

Coordinate Invariance of d on 1-forms

We define $d\omega$, $\omega = \sum f_i dx_i$ in local coordinates (x_1, \dots, x_n) by

$$d\omega = \sum_{i < j} \left(-\frac{\partial f_i}{\partial x_j} + \frac{\partial f_j}{\partial x_i} \right) dx_i \wedge dx_j$$

(this is motivated by $d(f dx_i)$ should be $df \wedge dx_i$). For emphasis, let us call this $d_{(x_1, \dots, x_n)} \omega$ since in principle it depends on the choice of local coordinates (x_1, \dots, x_n) .

We want to show that $d_{(x_1, \dots, x_n)} \omega = d_{(\hat{x}_1, \dots, \hat{x}_n)} \omega$ in the sense that for the right-hand side we write ω in $(\hat{x}_1, \dots, \hat{x}_n)$ coordinates and compute d with the same formula as in the third line above - except with \hat{x} coordinates. Now

$$dx_i = \sum_l \frac{\partial x_i}{\partial \hat{x}_l} d\hat{x}_l \quad \text{so} \quad \omega = \sum_{i, l} f_i \frac{\partial x_i}{\partial \hat{x}_l} d\hat{x}_l \\ = \sum_i \left(\sum_l f_l \frac{\partial x_i}{\partial \hat{x}_l} \right) d\hat{x}_i$$

$$\text{So } \omega = \sum \hat{f}_i d\hat{x}_i \text{ where } \hat{f}_i = \sum_l f_l \frac{\partial x_i}{\partial \hat{x}_l}$$

$$\text{Thus } d_{(\hat{x}_1, \dots, \hat{x}_n)} \omega = \sum_{i < j} \left(-\frac{\partial \hat{f}_i}{\partial \hat{x}_j} + \frac{\partial \hat{f}_j}{\partial \hat{x}_i} \right) d\hat{x}_i \wedge d\hat{x}_j$$

$$\text{and } \frac{\partial \hat{f}_i}{\partial \hat{x}_j} = \frac{\partial}{\partial \hat{x}_j} \left(\sum_l f_l \frac{\partial x_i}{\partial \hat{x}_l} \right) = \sum_l \left(\frac{\partial f_l}{\partial \hat{x}_j} \frac{\partial x_i}{\partial \hat{x}_l} + f_l \frac{\partial^2 x_i}{\partial \hat{x}_j \partial \hat{x}_l} \right)$$

While $\hat{\frac{\partial f_i}{\partial \hat{x}_j}} = \sum_l \frac{\partial f_l}{\partial \hat{x}_i} \frac{\partial \hat{x}_l}{\partial \hat{x}_j} + f_l \frac{\partial^2 f_l}{\partial \hat{x}_i \partial \hat{x}_j}$ (by ② interchanging i, j in previous)

$$\text{So } \hat{\frac{\partial f_i}{\partial \hat{x}_j}} - \hat{\frac{\partial f_i}{\partial \hat{x}_j}} = \sum_l \left(-\frac{\partial f_l}{\partial \hat{x}_i} \frac{\partial \hat{x}_l}{\partial \hat{x}_j} + \frac{\partial f_l}{\partial \hat{x}_i} \frac{\partial \hat{x}_l}{\partial \hat{x}_j} \right) + 0$$

↑
2nd derivative terms cancel

Hence

$$\begin{aligned}
 d\omega_{(x_1, \dots, x_n)} &= \sum_{i < j} \left(\sum_l \left(-\frac{\partial f_l}{\partial \hat{x}_j} \frac{\partial \hat{x}_l}{\partial \hat{x}_i} + \frac{\partial f_l}{\partial \hat{x}_i} \frac{\partial \hat{x}_l}{\partial \hat{x}_j} \right) d\hat{x}_i \wedge d\hat{x}_j \right) \\
 &= \sum_{i < j} \left(\sum_{k,l} \left(-\frac{\partial f_k}{\partial \hat{x}_p} \frac{\partial \hat{x}_k}{\partial \hat{x}_j} \frac{\partial \hat{x}_l}{\partial \hat{x}_i} + \frac{\partial f_k}{\partial \hat{x}_p} \frac{\partial \hat{x}_k}{\partial \hat{x}_i} \frac{\partial \hat{x}_l}{\partial \hat{x}_j} \right) d\hat{x}_i \wedge d\hat{x}_j \right) \\
 &= \frac{1}{2} \sum_{l,k} \left[\sum_{i,j} \left(-\frac{\partial f_k}{\partial \hat{x}_k} \frac{\partial \hat{x}_k}{\partial \hat{x}_j} \frac{\partial \hat{x}_l}{\partial \hat{x}_i} + \frac{\partial f_k}{\partial \hat{x}_k} \frac{\partial \hat{x}_k}{\partial \hat{x}_i} \frac{\partial \hat{x}_l}{\partial \hat{x}_j} \right) d\hat{x}_i \wedge d\hat{x}_j \right] \\
 &\quad \text{= } dx_k \text{ when } i \text{ is summed out} \\
 &= \frac{1}{2} \sum_{l,k} -\frac{\partial f_k}{\partial \hat{x}_k} \left(\sum_{i,j} \left(\frac{\partial \hat{x}_k}{\partial \hat{x}_i} d\hat{x}_i \wedge \frac{\partial \hat{x}_k}{\partial \hat{x}_j} d\hat{x}_j \right) \right) \\
 &\quad \text{= } dx_k \text{ when } j \text{ is summed out} \\
 &+ \frac{1}{2} \sum_{l,k} \frac{\partial f_k}{\partial \hat{x}_k} \left(\sum_{i,j} \left(\frac{\partial \hat{x}_k}{\partial \hat{x}_i} d\hat{x}_i \wedge \frac{\partial \hat{x}_k}{\partial \hat{x}_j} d\hat{x}_j \right) \right) \\
 &\quad \text{= } dx_k \text{ when } i \text{ is summed out} \quad \text{= } dx_k \text{ when } j \text{ is summed out} \\
 &= \frac{1}{2} \sum_{l,k} \left(-\frac{\partial f_k}{\partial \hat{x}_k} dx_k \wedge dx_k + \frac{\partial f_k}{\partial \hat{x}_k} dx_k \wedge dx_k \right) \\
 &= \sum_{k < l} \left(-\frac{\partial f_k}{\partial \hat{x}_k} + \frac{\partial f_k}{\partial \hat{x}_k} \right) (dx_k \wedge dx_k) = d_{(x_1, \dots, x_n)} \omega.
 \end{aligned}$$

(3)

This is a bit messy, but it does work.

The messiness justifies the process of isolating (coordinate dependent) characterizing properties of ω and then verifying that ω in local coordinate has the characterizing properties and hence is coordinate independent.

Alternatively, one can show that at $p \in M$, $\omega = \sum f_i dx^i$

$$d\omega(X, Y) = X\omega(Y) - Y\omega(X) - \omega([X, Y])$$

X, Y vector fields in nbhd of p

Then (x_1, \dots, x_n) coordinates around p

Since RHS is coordinate independent, so is left hand side. To check this formula, note:

(1) It holds if $X = \frac{\partial}{\partial x_i}$, $Y = \frac{\partial}{\partial x_j}$; left hand side

$$= -\frac{\partial f_i}{\partial x_j} + \frac{\partial f_j}{\partial x_i} \quad \text{while} \quad \omega(Y) = f_j \quad \omega(X) = f_i$$

$$\text{so } X\omega(Y) - Y\omega(X) - \omega([X, Y]) = \frac{\partial f_i}{\partial x_j} - \frac{\partial f_i}{\partial x_j} \quad \checkmark$$

(2) $RHS(fX, Y) = f RHS(X, Y)$ and similarly

$$RHS(X, fY) = f RHS(X, Y).$$

Proof: $(fX)\omega(Y) = f(X\omega(Y))$

$$-Y\omega(fX) = -Y(f\omega(X)) = -Yf\omega(X) \\ -fY\omega(X)$$

$$-\omega([fX, Y]) = -\omega((-Yf)X + f[X, Y]) \\ = +(-Yf)\omega(X) - f\omega([X, Y]).$$

So Yf terms cancel when RHS is added up, f factors out

(4)

as required.
Of course $LHS(fX, Y) = f LHS(X, Y)$ and similarly for XfY .
Points (1) and (2) (and additive linearity on X and Y) of both sides thus imply that

$LHS = RHS$ since both give same answers
on $X = \sum f_i \frac{\partial}{\partial x_i}$ and $Y = \sum g_j \frac{\partial}{\partial x_j}$, namely

$$\text{both sides} = \sum_{i,j} f_i g_j LHS\left(\frac{\partial}{\partial x_i}, \frac{\partial}{\partial x_j}\right)$$

or
 RHS

(which at p depends only on $f_i(p)$ and $g_j(p)$).

This may seem a bit obscure at first!

Think it over. But it does show that

$d_{(x_1, \dots, x_n)} w$ at p is the same as dw

calculated at p in any other local coordinate system.

This is all a lot of trouble but d is really important so the trouble is worth going through! It also has a reasonably straightforward generalization to

k -forms, $k > 2$: $d(f dx_1 \wedge \dots \wedge dx_k) = df \wedge dx_1 \wedge \dots \wedge dx_k$
as definition and

$$dw(X_0, X_1, \dots, X_k) = \sum (-1)^i X_i w(X_0, X_1, \overset{i+1}{\underset{\text{omitted}}{X_2}}, \dots, X_k)$$

It is worthwhile to work out some specific instances. (5)
 For example, suppose $\omega = \frac{1}{2} r^2 d\theta$ in polar coordinates
 on $\mathbb{R}^2 - \{(0,0)\}$. Then $d_{(r,\theta)} \omega = r dr \wedge d\theta$.

$$\text{In } (x,y) \text{ coordinates, } dr = \frac{\partial r}{\partial x} dx + \frac{\partial r}{\partial y} dy = \frac{x}{\sqrt{x^2+y^2}} dx + \frac{y}{\sqrt{x^2+y^2}} dy$$

$$\text{and } d\theta = \frac{\partial \theta}{\partial x} dx + \frac{\partial \theta}{\partial y} dy = -\frac{y}{x^2+y^2} dx + \frac{x}{x^2+y^2} dy.$$

$$\text{So } \omega = \frac{1}{2} ((x^2+y^2)^{\frac{1}{2}})^2 \left(-\frac{y}{x^2+y^2} dx + \frac{x}{x^2+y^2} dy \right) = -\frac{y}{2} dx + \frac{x}{2} dy$$

$$\text{So } d_{(x,y)} \omega = -\frac{1}{2} dy \wedge dx + \frac{1}{2} dx \wedge dy = dx \wedge dy$$

$$\text{while } r dr \wedge d\theta = \sqrt{x^2+y^2} \left(\frac{x}{\sqrt{x^2+y^2}} dx + \frac{y}{\sqrt{x^2+y^2}} dy \right) \wedge \left(-\frac{y}{(x^2+y^2)^{3/2}} dx + \frac{x}{(x^2+y^2)^{3/2}} dy \right)$$

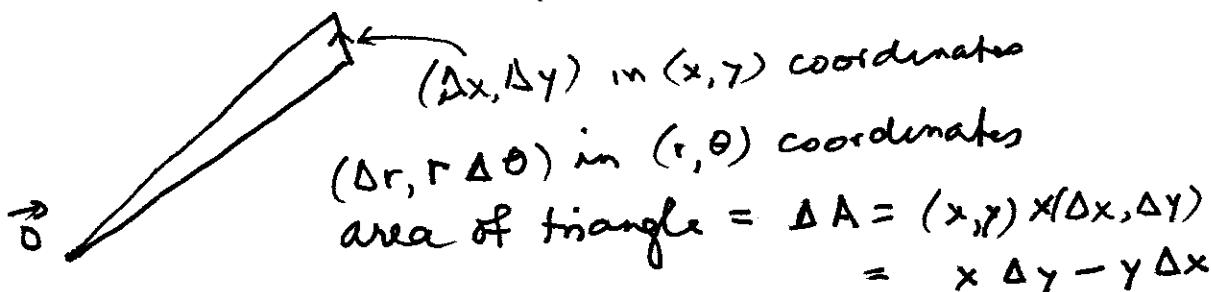
$$= \sqrt{x^2+y^2} \left(\frac{x^2}{(x^2+y^2)^{3/2}} dx \wedge dy - \frac{y^2}{(x^2+y^2)^{3/2}} dy \wedge dx \right)$$

$$= \frac{(x^2+y^2)\sqrt{x^2+y^2}}{(x^2+y^2)^{3/2}} dx \wedge dy = dx \wedge dy.$$

This confirms the coordinate invariance of d in
 this instance.

The geometric interpretation is of interest: By Green's/Stokes
 Theorem, $\oint_C \omega$ around a ^{simple} closed curve $C = \int_{\text{interior of } C} dx \wedge dy$

if C is oriented counterclockwise



also area of triangle = $\frac{1}{2} r \cdot (r\Delta\theta)$ part ⊥ radius,
 $r\Delta\theta$ does not
 count

 $= \frac{1}{2} r^2 \Delta\theta$

So $\omega(C) = \frac{dA}{dt}$ in both coordinate systems,
 if $A = \text{area swept out}$.