

## Curves on Surfaces.

Let  $S(u, v)$  be a (nonsingular) surface patch (i.e.  $\vec{S}_u \times \vec{S}_v \neq \vec{0}$ ). A curve  $(u(t), v(t))$  in the  $(u, v)$  domain of definition of  $S$  gives rise to a curve  $S(u(t), v(t))$  "in" the surface  $S$ , i.e. the image of the curve lies in the image of  $S$  (obviously) and if  $S$  is one-to-one, all curves in  $S$  arises in this way. We are interested in the curvature behavior of such curves

$\gamma(t) = S(u(t), v(t))$ . As usual, we compute the tangent vector:  $\gamma'(t) = \frac{du(t)}{dt} \vec{S}_u + \frac{dv(t)}{dt} \vec{S}_v$  so that  $\gamma'(t)$  belongs to the tangent plane at  $\gamma(t)$  (which is in  $S$ ). We restrict our attention to arc length parameter curves, i.e.  $\|\gamma'(t)\| = 1$  for all  $t$ , and we write  $\gamma(s)$  to emphasize this condition. Then  $T = T(s) = \frac{du(s)}{ds} \vec{S}_u + \frac{dv(s)}{ds} \vec{S}_v$

and (because  $\|T\| = 1$ )

$$1 = \left( \frac{du}{ds} \right)^2 E + 2F \frac{du}{ds} \frac{dv}{ds} + \left( \frac{dv}{ds} \right)^2 G$$

where  $E = \langle \vec{S}_u, \vec{S}_u \rangle$ ,  $F = \langle \vec{S}_u, \vec{S}_v \rangle$ , and  $G = \langle \vec{S}_v, \vec{S}_v \rangle$  as usual. Then a further differentiation yields

$$\frac{dT}{ds} (= k N_r) = \frac{d^2 u}{ds^2} \vec{S}_u + \frac{d^2 v}{ds^2} \vec{S}_v$$

$$+ \frac{du}{ds} \frac{d}{ds} \left( \vec{S}_u \Big|_{\gamma(s)} \right) + \frac{dv}{ds} \frac{d}{ds} \left( \vec{S}_v \Big|_{\gamma(s)} \right).$$

$N_r$  = normal of curve

The presence of the last two terms require a little thought. The point is that when  $t$  varies the vectors  $S_u$  and  $S_v$ , which are evaluated at  $\gamma(s)$ , also change. Indeed,

$$\frac{d}{ds}(S_u|_{\gamma(s)}) = \frac{du}{ds} S_{uu} + \frac{dv}{ds} S_{uv}$$

and  $\frac{d}{ds}(S_v|_{\gamma(s)}) = \frac{du}{ds} S_{vu} + \frac{dv}{ds} S_{vv}.$

Thus

$$\begin{aligned} \frac{dT}{ds} &= \frac{d^2 u}{ds^2} S_{uu} + \frac{d^2 v}{ds^2} S_{vv} \\ &\quad + \left(\frac{du}{ds}\right)^2 S_{uu} + 2\left(\frac{du}{ds}\right)\left(\frac{dv}{ds}\right) S_{uv} + \left(\frac{dv}{ds}\right)^2 S_{vv}. \end{aligned}$$

$= KN_\gamma$  (where  $K, N$  are as usual for space curves,  $N_\gamma$  = curve normal of  $\gamma$ ).

There is of course no reason why the curve normal  $N_\gamma$  has to be the surface normal  $N$  (Example: think about a nonequatorial parallel of latitude on the sphere). It is of interest to look, however, at the "normal component" of  $\frac{dT}{ds}$  here: namely,  $\langle N, \frac{dT}{ds} \rangle$ .

Thus it satisfies

$$\begin{aligned} -\langle N, \frac{dT}{ds} \rangle &= \left(\frac{du}{ds}\right)^2 L_{11} + 2\left(\frac{du}{ds}\right)\left(\frac{dv}{ds}\right) L_{12} \\ &\quad + \left(\frac{dv}{ds}\right)^2 L_{22}. \end{aligned}$$

(The - sign arises from  $L_{11} = -\langle S_{uu}, N \rangle$ ).

The tangent part of  $\frac{dT}{ds}$  is given by

$$\begin{aligned} \frac{d^2u}{ds^2} S_u + \frac{d^2v}{ds^2} S_v + \left(\frac{du}{ds}\right)^2 T(S_{uu}) + 2 \left(\frac{du}{ds} \frac{dv}{ds}\right) T(S_{uv}) \\ + \left(\frac{dv}{ds}\right)^2 T(S_{vv}) \end{aligned}$$

where  $T(\cdot)$  = tangent part as earlier.

The normal part of the  $\frac{dT}{ds}$  vector is determined entirely by  $T$  itself, that is, by  $\frac{du}{ds}$  &  $\frac{dv}{ds}$ . But the tangent part depends also on the second derivatives of  $\gamma$ ,  $\frac{d^2u}{ds^2}$  and  $\frac{d^2v}{ds^2}$ .

Note that for a given tangent vector (i.e. given  $\frac{du}{dt}, \frac{dv}{dt}$ ), the minimum possible value for  $\|\frac{dT}{ds}\|$  is obtained if the tangent part of  $\frac{dT}{ds}$  is 0, in which case the value of  $\|\frac{dT}{ds}\|$  is

$$\left| \left(\frac{du}{ds}\right)^2 L_{11} + 2 \frac{du}{ds} \frac{dv}{ds} L_{12} + \left(\frac{dv}{ds}\right)^2 L_{22} \right|.$$

(We could always realize this minimum, too, by taking  $\gamma$  to be a "sector" of  $S$  by the plane contain  $N$  and the specified tangent vector).

We set  $K_g =$  the length of the tangent ( $\gamma$ ) part of  $\frac{dT}{ds}$  and  $|K_n|$  (or normal) = the length of the normal part. Then the curvature  $K$  of  $\gamma$  as a space curve satisfies

$$K^2 = K_g^2 + K_n^2$$

Moreover,  $K |\cos \theta| = |K_n|$  where  $\theta =$  the angle between the surface normal  $N$  and the curve normal  $N_\gamma$ . The <sup>seemingly</sup> strange notation  $|K_n|$  suggests that there is a  $K_n$  with a  $\pm$  sign: this is just  $K_n = \langle (\text{normal part of } \frac{dT}{ds}), N \rangle$  where  $N =$  the surface normal!  $-K_n = \left( \frac{du}{dt} \right)^2 L_{11} + 2 \left( \frac{du}{dt} \right) \left( \frac{dv}{dt} \right) L_{12} + \left( \frac{dv}{dt} \right)^2 L_{22}$ . The standard theory of symmetric quadratic forms gives that the Gauss curvature at the point = the product of the maximum and minimum values of the "normal curvature"  $K_n$  of arclength parameter curves through the point.

Recalling that the tangent parts of  $S_{uu}$ ,  $S_{uv}$ , and  $S_{vv}$  are intrinsic, we see also that the tangent part of  $\frac{dT}{ds}$  is also intrinsic.

Definition: A curve  $\gamma(s)$  (with arclength parameter) is an arclength parameter geodesic in  $S$  if the tangent part of  $\frac{dT}{ds}$  is 0 for all  $s$  values (for which the curve is defined).

It is not clear *a priori* that geodesics exist! All we have is an equation that they have to satisfy everywhere, i.e., for all  $s$  values.

Let us write this equation a little more neatly.

$$\text{First write } T(S_{uu}) = \Gamma_{11}^1 S_u + \Gamma_{11}^2 S_v$$

$$T(S_{uv}) = \Gamma_{12}^1 S_u + \Gamma_{12}^2 S_v$$

$$T(S_{vu}) = \Gamma_{21}^1 S_u + \Gamma_{21}^2 S_v$$

$$T(S_{vv}) = \Gamma_{22}^1 S_u + \Gamma_{22}^2 S_v.$$

Since  $S_{uv} = S_{vu}$ , we see that  $\Gamma_{12}^1 = \Gamma_{21}^1$  and  $\Gamma_{12}^2 = \Gamma_{21}^2$ . The  $\Gamma$ 's are function on the surface: they can be and usually are different at different points of the patch  $S$ .

In this notation, separating \$S\_u\$ and \$S\_v\$ components:

$$T\left(\frac{dT}{ds}\right) = \left( \frac{d^2u}{ds^2} + \left(\frac{du}{ds}\right)^2 \Gamma_{11}^1 + 2 \frac{du}{ds} \frac{dv}{ds} \Gamma_{12}^1 + \left(\frac{dv}{ds}\right)^2 \Gamma_{22}^1 \right) S_u$$

$$+ \left( \frac{d^2v}{ds^2} + \left(\frac{du}{ds}\right)^2 \Gamma_{11}^2 + 2 \frac{du}{ds} \frac{dv}{ds} \Gamma_{12}^2 + \left(\frac{dv}{ds}\right)^2 \Gamma_{22}^2 \right) S_v.$$

In particular, a curve  $\gamma(s) = (u(s), v(s))$  is a geodesic if and only if the  $u(s)$  and  $v(s)$  satisfy two ("coupled") ordinary differential equations

$$0 = \frac{d^2u}{ds^2} + \left(\frac{du}{ds}\right)^2 \Gamma_{11}^1 \Big|_{(u(s), v(s))} + 2 \frac{du}{ds} \frac{dv}{ds} \Gamma_{12}^1 + \left(\frac{dv}{ds}\right)^2 \Gamma_{22}^1$$

$$0 = \frac{d^2v}{ds^2} + \left(\frac{du}{ds}\right)^2 \Gamma_{11}^2 + 2 \frac{du}{ds} \frac{dv}{ds} \Gamma_{12}^2 + \left(\frac{dv}{ds}\right)^2 \Gamma_{22}^2$$

This looks nicer if we start writing  $u_1, u_2$  for  $u$  and  $v$ . Namely, the equations become

$$0 = \frac{d^2u_1}{ds^2} + \sum_{i,j} \Gamma_{ij}^1 \frac{du_i}{ds} \frac{du_j}{ds}$$

$$0 = \frac{d^2u_2}{ds^2} + \sum_{i,j} \Gamma_{ij}^2 \frac{du_i}{ds} \frac{du_j}{ds}.$$

It is worth noting that we did not really make much use of arclength parameter here: it was only relevant in talking about the curvature of the curve as a space curve and so on. In particular,

7

we could talk about geodesics as being  
 curves that satisfied these equations  
 whether or not  $s$  was an length parameter.  
 It turns out this does not matter too much:  
 a geodesic in this extended sense always  
 has  $\|\frac{d\gamma}{ds}\|$  constant, though not necessarily 1.  
 But it is convenient to allow this great  
 freedom, so from now on we call any  
 curve  $\gamma$  satisfying the two equations a  
 geodesic and call it an "arc length  
 parameter geodesic" explicitly if it does  
 have arc length parameter.

It is not hard to see why geodesics  
 in this more general sense are parameterized  
 "proportional to arc length": For any curve  
 $\gamma(t)$  at all (in a surface  $S$ ),  $\frac{d}{dt} \langle \gamma', \gamma' \rangle$   
 $= 2 \langle \gamma'', \gamma' \rangle = 2 \langle \text{tangent part of } \gamma'', \gamma' \rangle$   
 since  $\gamma'$  is tangent to  $S$  (lies in the tangent plane)  
 [Here  $' = t\text{-derivative}$ ]. So if tangent part  
 of  $\gamma''$  ( $=$  tangent part of  $\frac{d^2\gamma}{dt^2}$ ) = 0, then  $\langle \gamma', \gamma' \rangle$   
 is constant, which is the same as  $\|\frac{d\gamma}{dt}\|$  is  
 constant, which is itself the same as proportional  
 to arc length parameterization.

Geodesics (with arclength parameter) are important because they are the analogues for surfaces in general of straight lines in the plane or great circles on the sphere.

Namely, if  $\gamma(s)$  is an (arclength parameter) geodesic then for a fixed  $s_0$  and all  $s$  close enough to  $s_0$ , the length of any curve in  $S$  connecting  $\gamma(s_0)$  to  $\gamma(s)$  is at least  $|s-s_0|$ , with equality if and only if  $\sigma$  is  $\gamma$  up to reparametrization. In summary form,  $\gamma$  is locally the unique shortest connection between points on it. (Think about great circles to see why we need "locally" and cannot just say that  $\gamma$  is the shortest connection).

Proving this will take a little work! but intuitively, it seems likely: a geodesic does not "turn" in terms of the tangential component of its acceleration, so one expects it is heading straight for where it goes.