

Amplituhedra and origami

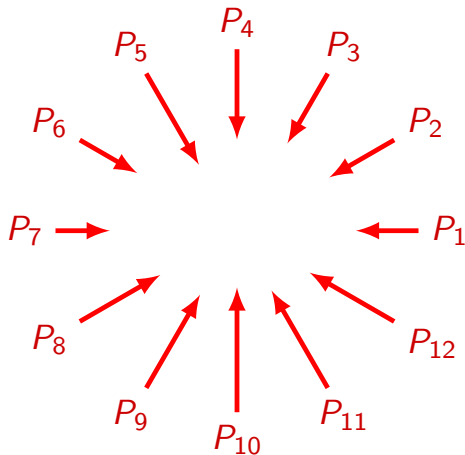
Pavel Galashin (UCLA)

Cornell Discrete Geometry and Combinatorics Seminar
December 9, 2024

[arXiv:2410.09574](https://arxiv.org/abs/2410.09574)

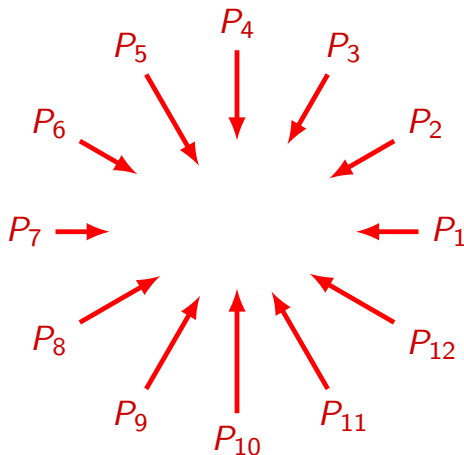
Scattering amplitudes

- Consider incoming particles with momenta $P_1, P_2, \dots, P_n$.



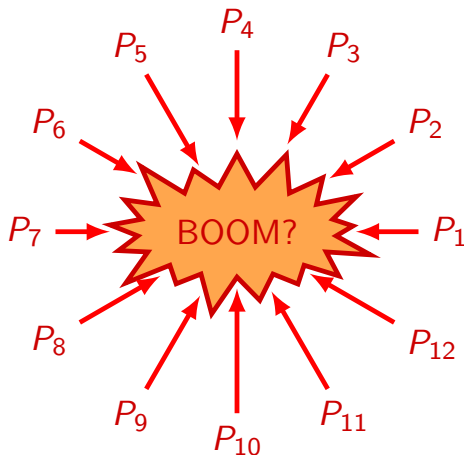
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Computing scattering amplitudes:

Efficient? Conceptual?

Old-fashioned perturbation theory [Dirac, 1926] 1/10 2/10

$$\begin{aligned}
 E_n^{(1)} &= V_{nn} \\
 E_n^{(2)} &= \frac{|V_{nnk_2}|^2}{E_{nk_2}} \\
 E_n^{(3)} &= \frac{V_{nk_3} V_{k_3k_2} V_{k_2n}}{E_{nk_2} E_{nk_3}} - V_{nn} \frac{|V_{nk_3}|^2}{E_{nk_3}^2} \\
 E_n^{(4)} &= \frac{V_{nk_4} V_{k_4k_3} V_{k_3k_2} V_{k_2n}}{E_{nk_2} E_{nk_3} E_{nk_4}} - \frac{|V_{nk_4}|^2 |V_{nk_2}|^2}{E_{nk_4}^2 E_{nk_2}} - V_{nn} \frac{V_{nk_4} V_{k_4k_3} V_{k_3n}}{E_{nk_3}^2 E_{nk_4}} - V_{nn} \frac{V_{nk_4} V_{k_4k_2} V_{k_2n}}{E_{nk_2} E_{nk_4}^2} + V_{nn}^2 \frac{|V_{nk_4}|^2}{E_{nk_4}^3} \\
 &= \frac{V_{nk_4} V_{k_4k_3} V_{k_3k_2} V_{k_2n}}{E_{nk_2} E_{nk_3} E_{nk_4}} - E_n^{(2)} \frac{|V_{nk_4}|^2}{E_{nk_4}^2} - 2V_{nn} \frac{V_{nk_4} V_{k_4k_3} V_{k_3n}}{E_{nk_3}^2 E_{nk_4}} + V_{nn}^2 \frac{|V_{nk_4}|^2}{E_{nk_4}^3} \\
 E_n^{(5)} &= \frac{V_{nk_5} V_{k_5k_4} V_{k_4k_3} V_{k_3k_2} V_{k_2n}}{E_{nk_2} E_{nk_3} E_{nk_4} E_{nk_5}} - \frac{V_{nk_5} V_{k_5k_4} V_{k_4n}}{E_{nk_4}^2 E_{nk_5}} \frac{|V_{nk_2}|^2}{E_{nk_2}} - \frac{V_{nk_5} V_{k_5k_2} V_{k_2n}}{E_{nk_2} E_{nk_5}^2} \frac{|V_{nk_2}|^2}{E_{nk_2}} - \frac{|V_{nk_5}|^2 V_{nk_3} V_{k_3k_2} V_{k_2n}}{E_{nk_3}^2 E_{nk_5}} \frac{E_{nk_2} E_{nk_3}}{E_{nk_2} E_{nk_3}} \\
 &\quad - V_{nn} \frac{V_{nk_5} V_{k_5k_4} V_{k_4k_3} V_{k_3n}}{E_{nk_3}^2 E_{nk_4} E_{nk_5}} - V_{nn} \frac{V_{nk_5} V_{k_5k_4} V_{k_4k_2} V_{k_2n}}{E_{nk_2} E_{nk_4}^2 E_{nk_5}} - V_{nn} \frac{V_{nk_5} V_{k_5k_3} V_{k_3k_2} V_{k_2n}}{E_{nk_2} E_{nk_3} E_{nk_5}^2} + V_{nn} \frac{|V_{nk_5}|^2 |V_{nk_3}|^2}{E_{nk_5}^2 E_{nk_3}^2} + 2V_{nn} \frac{|V_{nk_5}|^2 |V_{nk_2}|^2}{E_{nk_5}^3 E_{nk_2}} \\
 &\quad + V_{nn}^2 \frac{V_{nk_5} V_{k_5k_4} V_{k_4n}}{E_{nk_4}^3 E_{nk_5}} + V_{nn}^2 \frac{V_{nk_5} V_{k_5k_3} V_{k_3n}}{E_{nk_3}^3 E_{nk_5}} + V_{nn}^2 \frac{V_{nk_5} V_{k_5k_2} V_{k_2n}}{E_{nk_2}^3 E_{nk_5}} - V_{nn}^3 \frac{|V_{nk_5}|^2}{E_{nk_5}^4} \\
 &= \frac{V_{nk_5} V_{k_5k_4} V_{k_4k_3} V_{k_3k_2} V_{k_2n}}{E_{nk_2} E_{nk_3} E_{nk_4} E_{nk_5}} - 2E_n^{(2)} \frac{V_{nk_5} V_{k_5k_4} V_{k_4n}}{E_{nk_4}^2 E_{nk_5}} - \frac{|V_{nk_5}|^2 V_{nk_3} V_{k_3k_2} V_{k_2n}}{E_{nk_3}^2 E_{nk_5}} \frac{E_{nk_2} E_{nk_3}}{E_{nk_2} E_{nk_3}} \\
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 &\quad + V_{nn}^2 \left(2 \frac{V_{nk_5} V_{k_5k_4} V_{k_4n}}{E_{nk_4}^3 E_{nk_5}} + \frac{V_{nk_5} V_{k_5k_3} V_{k_3n}}{E_{nk_3}^2 E_{nk_5}^2} \right) - V_{nn}^3 \frac{|V_{nk_5}|^2}{E_{nk_5}^4}
 \end{aligned}$$



Scattering amplitudes

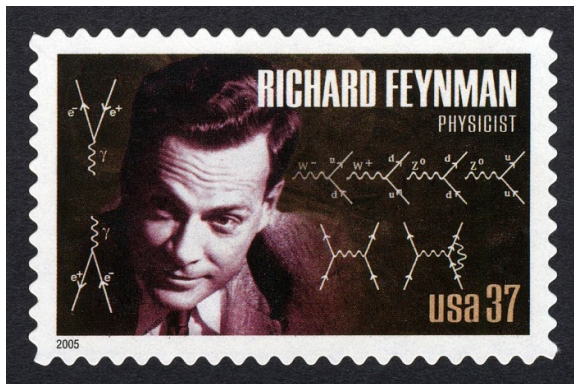
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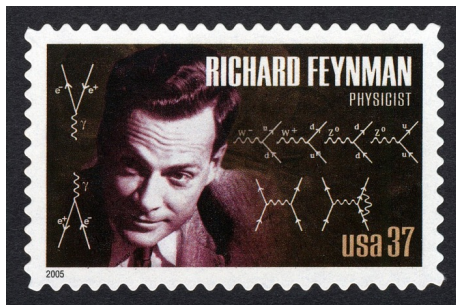
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Each Feynman diagram is the sum of exponentially many old-fashioned terms.

— [Wikipedia]

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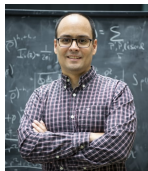
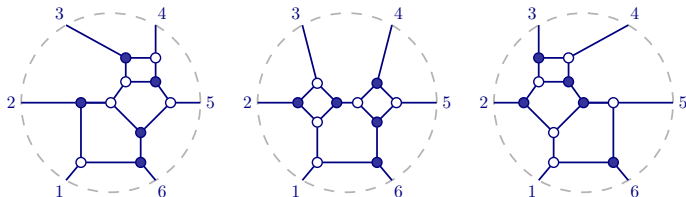
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Loops	# Feynman (super)graphs	# BCFW Cells
1	940	1
2	47,380	10
3	4,448,500	146
4	672,315,700	2,684
5	148,251,680,500	56,914
6	44,838,422,282,500	1,329,324
7	17,796,990,083,372,500	33,291,164
8	8,968,512,580,259,732,500	878,836,728
9	5,592,013,331,255,143,292,500	24,175,924,094
10	4,225,692,640,945,498,084,862,500	687,444,432,396
11	3,804,754,710,505,713,091,940,312,500	20,086,271,785,340

credit: J. Bourjaily

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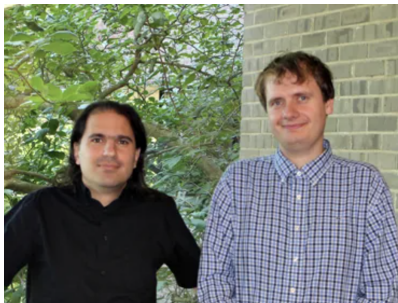
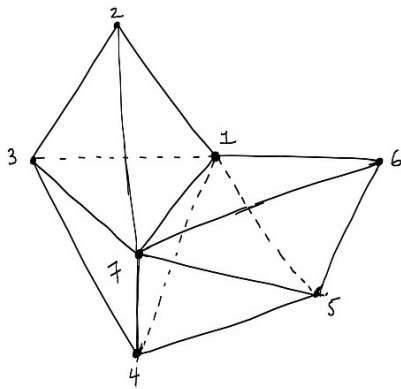
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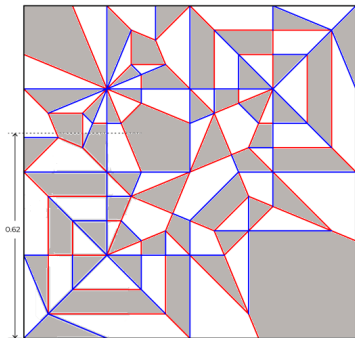
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Today: origami [G., 2024]	?/10	?/10



Classic Snowman v1

Model by Michelle Fung
Designed: 12/2020
Crease Pattern: 12/2020
www.michellefung.net



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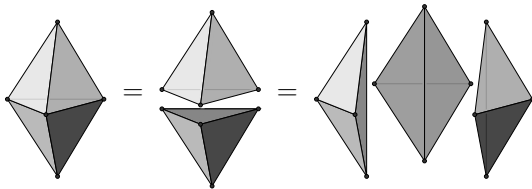
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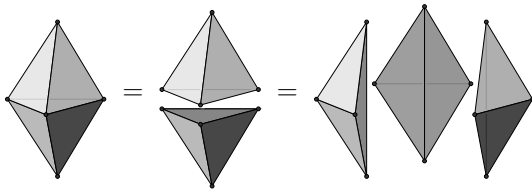


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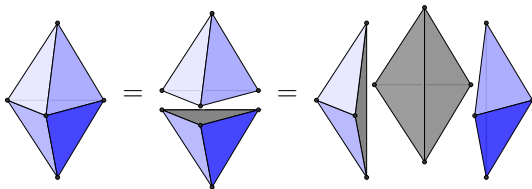


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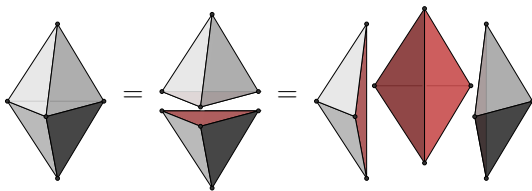


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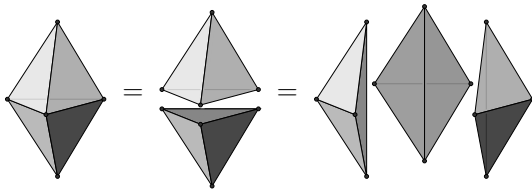


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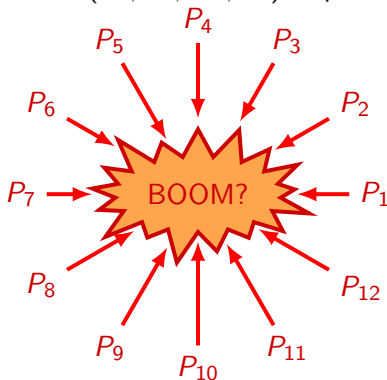


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- Spurious singularities $\longleftrightarrow$ boundaries between pieces
- $\mathcal{A}(P_1, P_2, \dots, P_n)$ $\longleftrightarrow$ “volume” of the amplituhedron

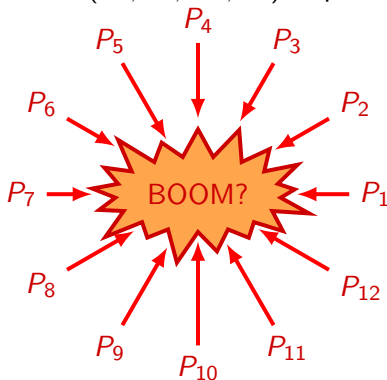
Particle momenta

- Consider incoming particles with momenta $P_1, P_2, \dots, P_n$.
- Scattering amplitude $\mathcal{A}(P_1, P_2, \dots, P_n) = \text{probability of 'scattering'}$.



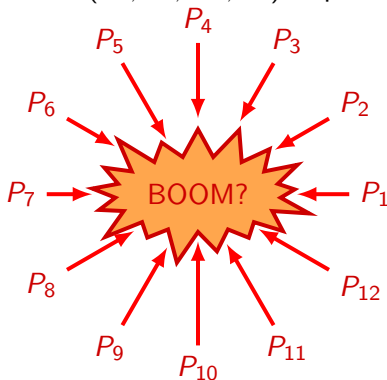
Particle momenta

- Consider incoming particles with momenta $P_1, P_2, \dots, P_n \in \mathbb{R}^{3,1}$ which are light-like ($P_i^2 = 0$) and satisfy $P_1 + P_2 + \dots + P_n = 0$.
- Scattering amplitude $\mathcal{A}(P_1, P_2, \dots, P_n) =$ probability of 'scattering'.



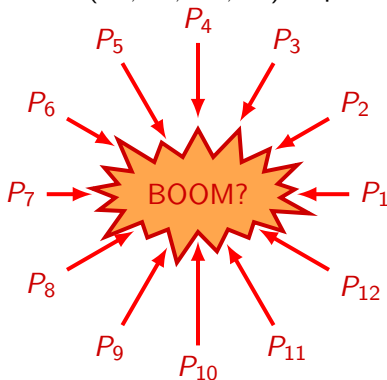
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- Here $P = (p_0, p_1, p_2, p_3)$ and $P^2 = p_0^2 + p_1^2 + p_2^2 - p_3^2$.
- Scattering amplitude $\mathcal{A}(P_1, P_2, \dots, P_n) =$ probability of 'scattering'.



Particle momenta

- Consider incoming particles with momenta $P_1, P_2, \dots, P_n \in \mathbb{R}^{2,2}$ which are light-like ($P_i^2 = 0$) and satisfy $P_1 + P_2 + \dots + P_n = 0$.
- Here $P = (p_0, p_1, p_2, p_3)$ and $P^2 = p_0^2 + p_1^2 - p_2^2 - p_3^2$.
- Scattering amplitude $\mathcal{A}(P_1, P_2, \dots, P_n) = \text{probability of 'scattering'}$.



Particle momenta

- Consider incoming particles with momenta $P_1, P_2, \dots, P_n \in \mathbb{R}^{2,2}$ which are light-like ($P_i^2 = 0$) and satisfy $P_1 + P_2 + \dots + P_n = 0$.
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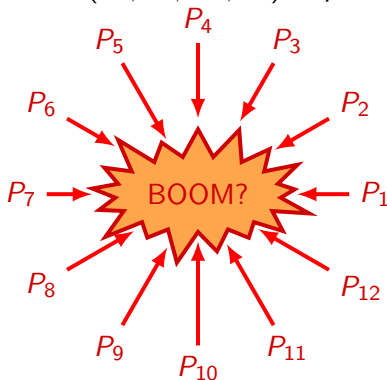
On the contrary, for *complex-valued momenta* p^μ , the angle and square spinors are independent.¹ It may not seem physical to take p^μ complex, but it is a very very very useful strategy. We will see this repeatedly.

¹ One can keep p^μ real and change the spacetime signature to $(-, +, -, +)$; in that case, the angle and square spinors are real and independent.

[Evang, Huang. *Scattering amplitudes in gauge theory and gravity*. Cambridge University Press, Cambridge, 2015.]

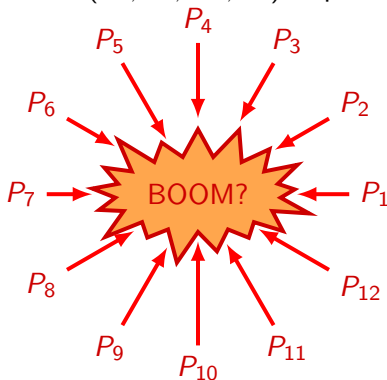
Particle momenta

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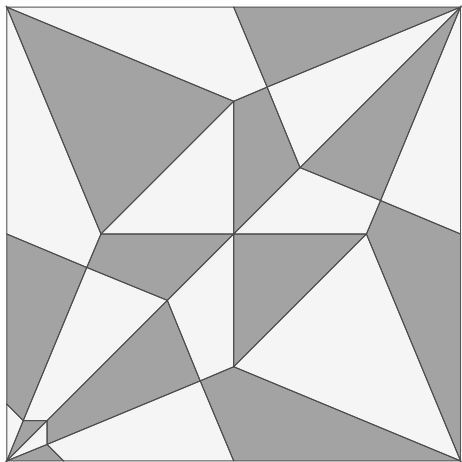
Particle momenta

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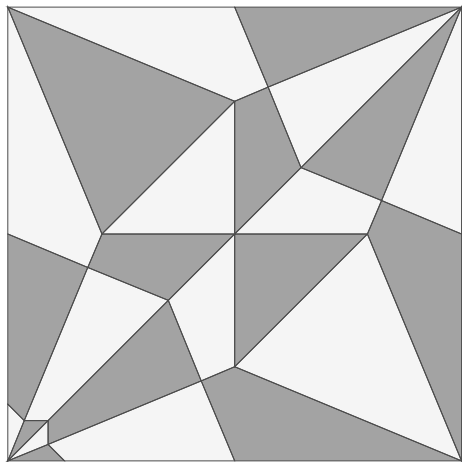


- Think of a light-like momentum vector $P \in \mathbb{R}^{2,2}$ as a pair $(P^T, P^O) \in \mathbb{C}^2$ of complex numbers satisfying $|P^T| = |P^O|$.

Origami crease patterns

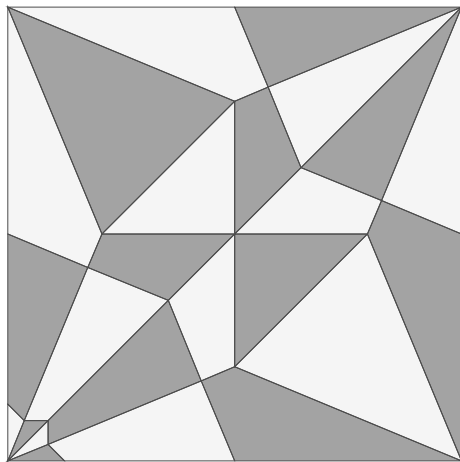


Origami crease patterns



Faces: convex polygons
colored black and white;

Origami crease patterns



Faces: convex polygons
colored black and white;

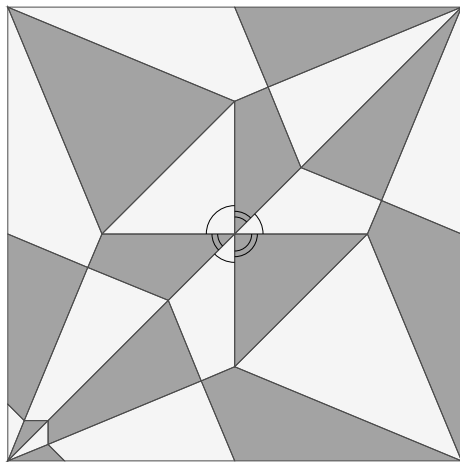
Angle condition:

$\text{sum}(\text{white angles}) = \pi$,

$\text{sum}(\text{black angles}) = \pi$,

around each interior vertex.

Origami crease patterns



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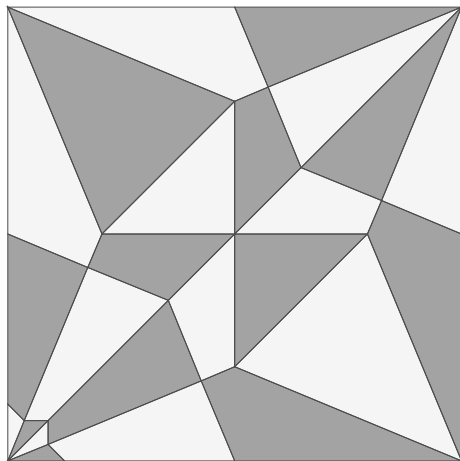
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around each interior vertex.

Origami crease patterns

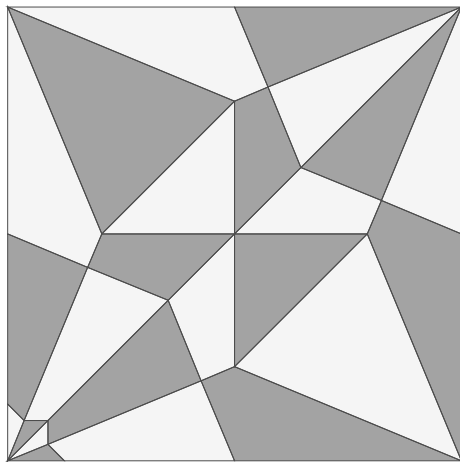


Faces: convex polygons
colored black and white;

Angle condition:
 $\text{sum}(\text{white angles}) = \pi$,
 $\text{sum}(\text{black angles}) = \pi$,
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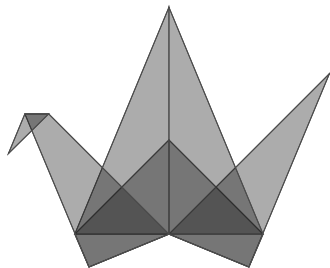
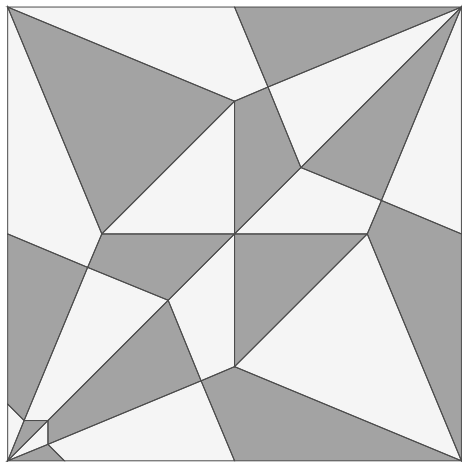
Origami map $\mathcal{O}$:
isometry on each face
preserving/reversing
the orientations of
white/black faces.

Origami crease patterns



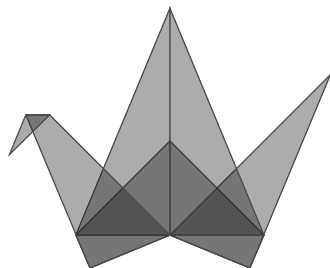
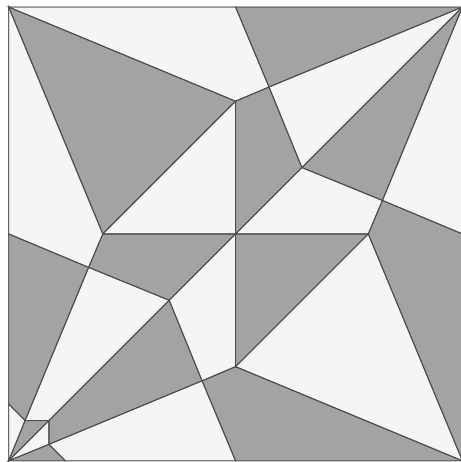
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Origami crease patterns



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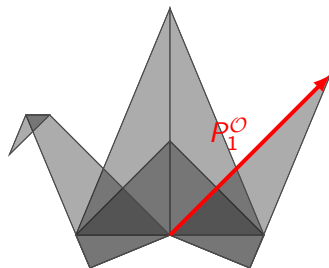
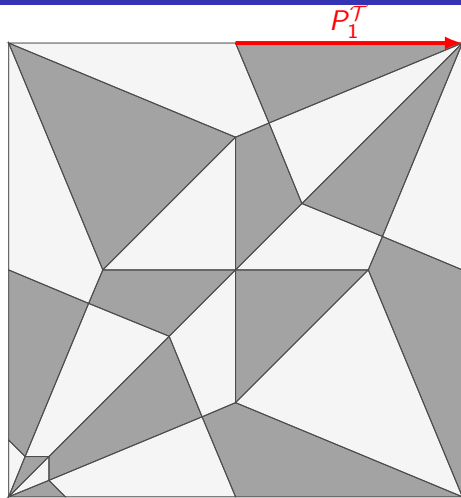
Origami crease patterns



Origami map $\mathcal{O}$:
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the orientations of
white/black faces.

Boundary vectors P_i^T and their images $P_i^{\mathcal{O}}$ under $\mathcal{O}$ satisfy $|P_i^T| = |P_i^{\mathcal{O}}|$!

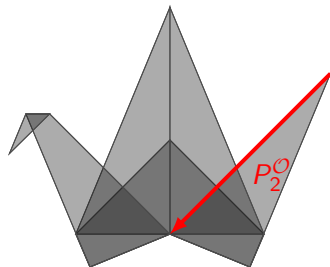
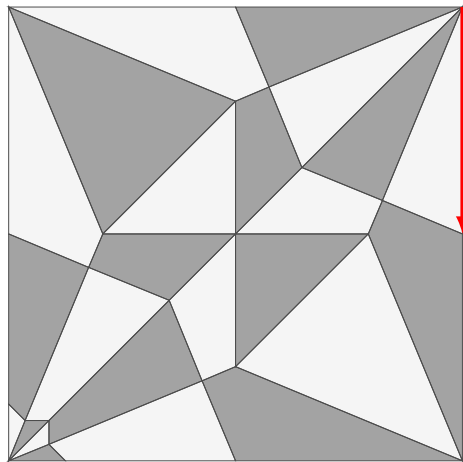
Origami crease patterns



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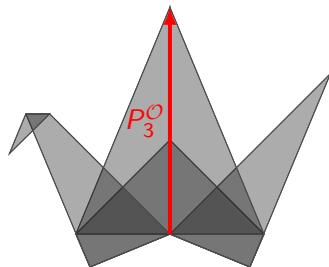
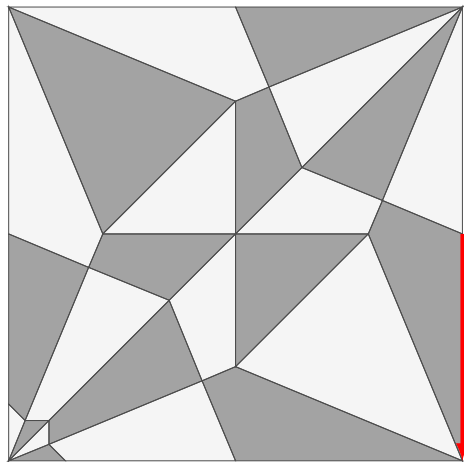
Origami crease patterns



Origami map $\mathcal{O}$:
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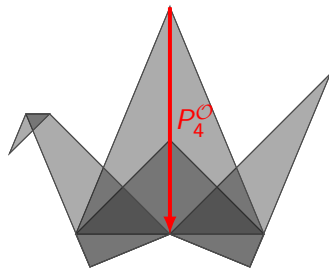
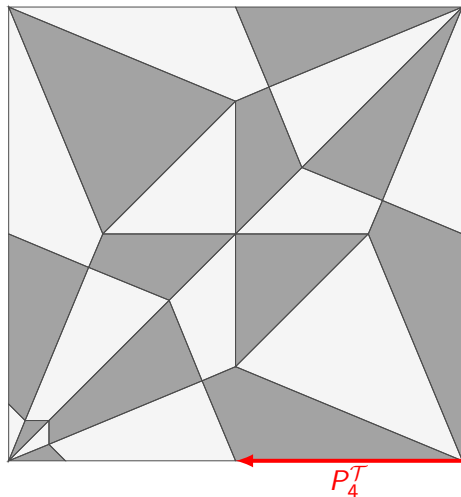
Origami crease patterns



Origami map $\mathcal{O}$:
 P_3^T isometry on each face
preserving/reversing
the orientations of
white/black faces.

Boundary vectors P_i^T and their images $P_i^{\mathcal{O}}$ under $\mathcal{O}$ satisfy $|P_i^T| = |P_i^{\mathcal{O}}|$!

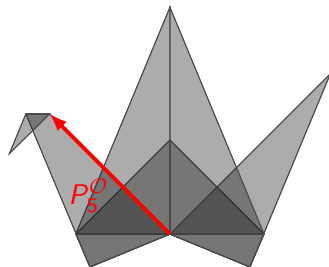
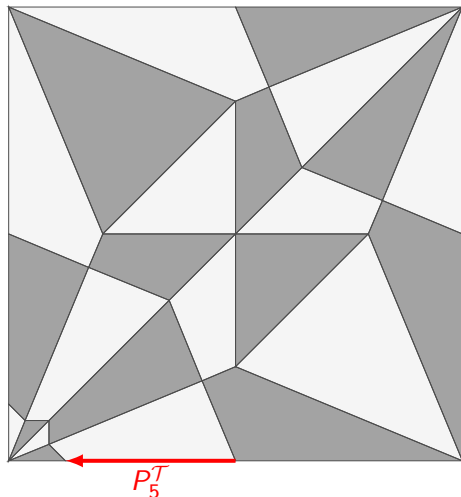
Origami crease patterns



Origami map $\mathcal{O}$:
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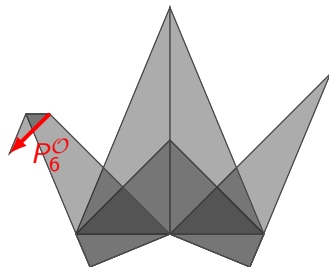
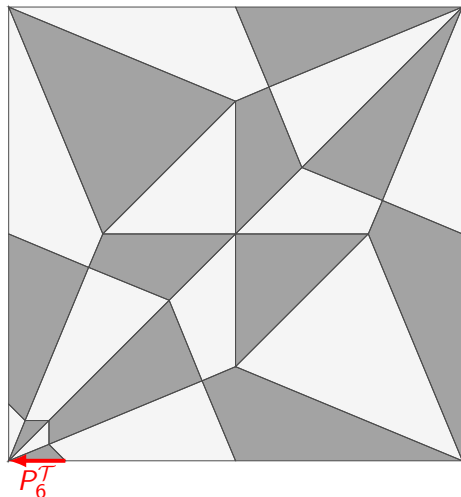
Origami crease patterns



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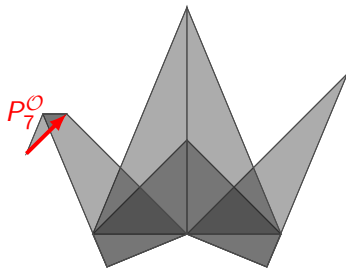
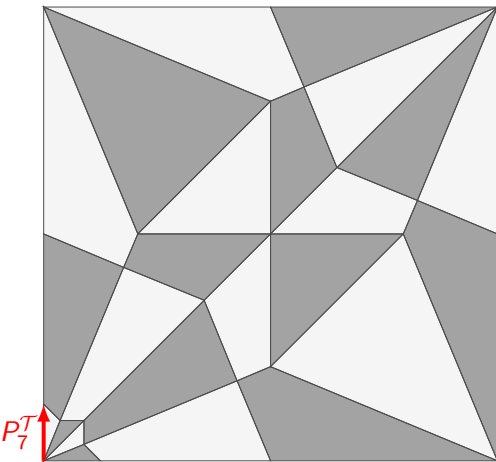
Origami crease patterns



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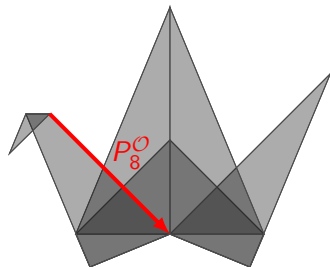
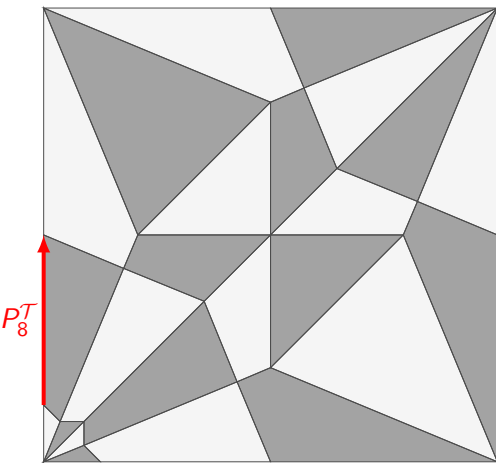
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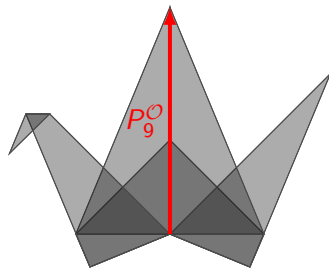
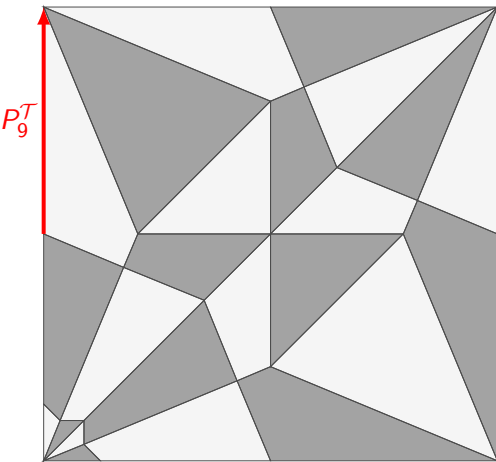
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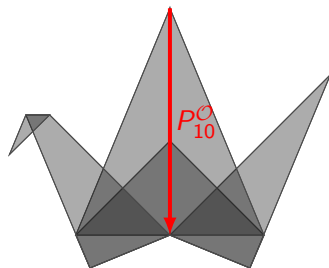
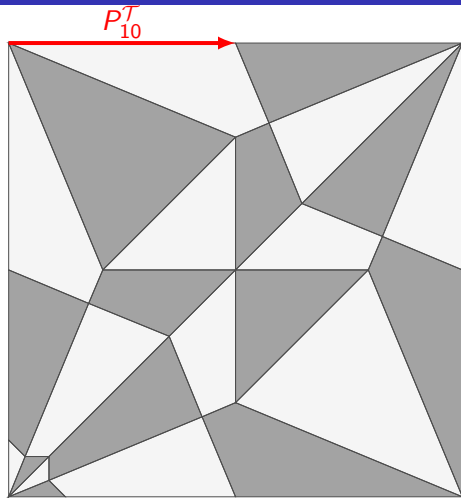
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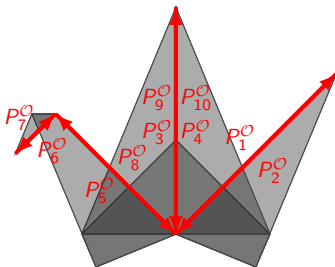
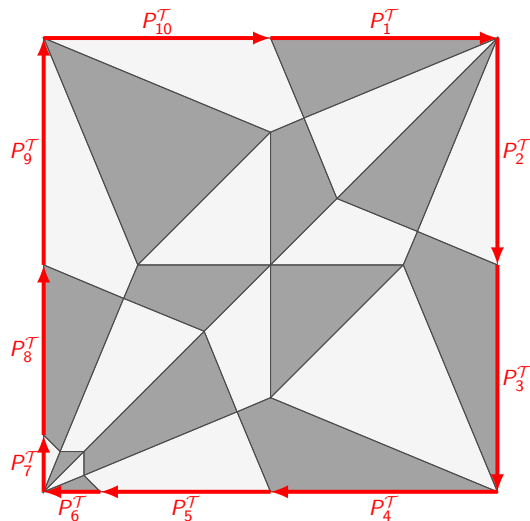
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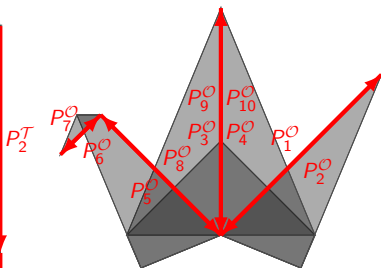
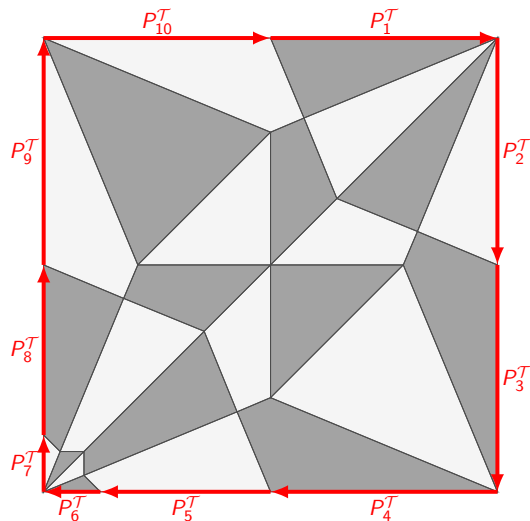
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Origami crease patterns



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Main result (preview):

$\mathcal{A}(P_1, \dots, P_n) = \text{integral over origami crease patterns with boundary } P_1, \dots, P_n.$

Grassmannian

$k \times n$ matrix C

$$C = \begin{bmatrix} 1 & a & 0 & -b \\ 0 & c & 1 & d \end{bmatrix}$$

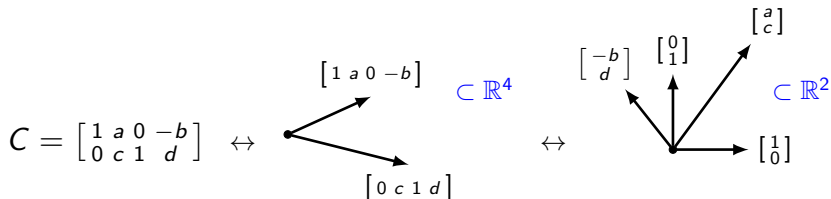
Grassmannian

$k \times n$ matrix $C \iff k$ row vectors in $\mathbb{R}^n$

$$C = \begin{bmatrix} 1 & a & 0 & -b \\ 0 & c & 1 & d \end{bmatrix} \iff \begin{array}{c} \nearrow [1 \ a \ 0 \ -b] \\ \searrow [0 \ c \ 1 \ d] \end{array} \subset \mathbb{R}^4$$

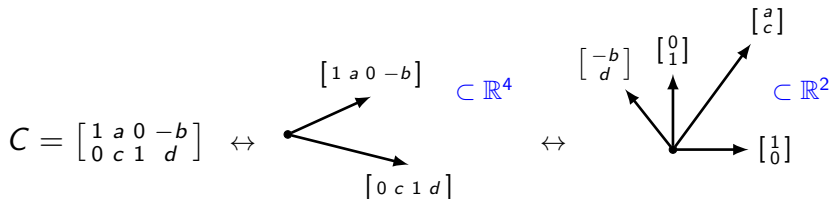
Grassmannian

$k \times n$ matrix $C \iff k$ row vectors in $\mathbb{R}^n \iff n$ column vectors in $\mathbb{R}^k$



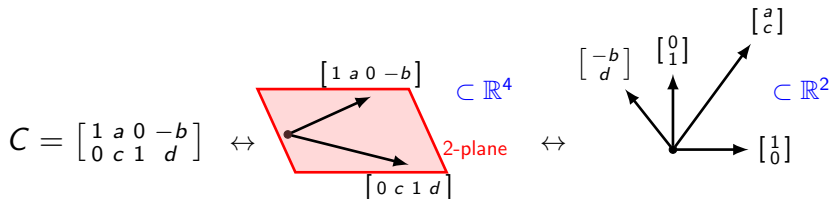
Grassmannian

$k \times n$ matrix $C \iff k$ row vectors in $\mathbb{R}^n \iff n$ column vectors in $\mathbb{R}^k$
mod row operations



Grassmannian

$k \times n$ matrix C mod row operations $\leftrightarrow$ ~~k row vectors in $\mathbb{R}^n$~~
 k -plane inside $\mathbb{R}^n$ $\leftrightarrow$ n column vectors in $\mathbb{R}^k$

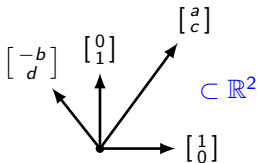
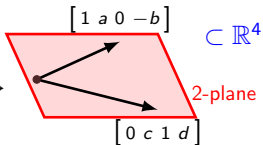


Grassmannian

 $k \times n$ matrix C
mod row operations
$$\Leftrightarrow \begin{array}{l} \text{~~k row vectors in } \mathbb{R}^n \\ \text{k-plane inside } \mathbb{R}^n \end{array}~~$$
$$\leftrightarrow n \text{ column vectors in } \mathbb{R}^k$$

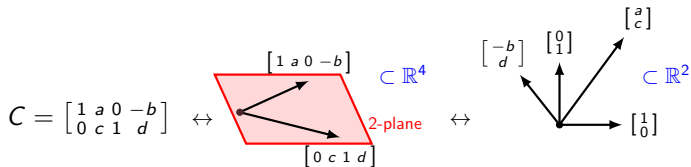
mod $GL_k(\mathbb{R})$ -action

$$C = \begin{bmatrix} 1 & a & 0 & -b \\ 0 & c & 1 & d \end{bmatrix}$$

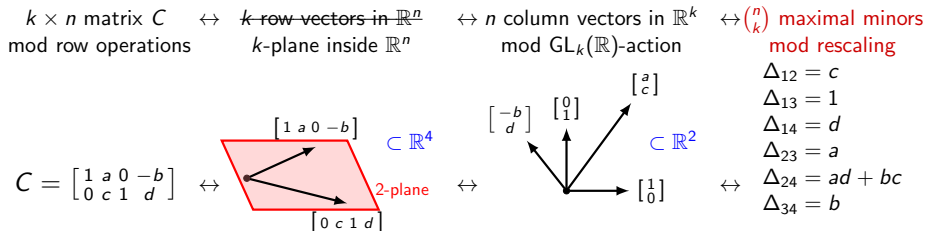


Grassmannian

$k \times n$ matrix C $\leftrightarrow$ ~~k row vectors in $\mathbb{R}^n$~~
 mod row operations $\quad$ k -plane inside $\mathbb{R}^n$ $\leftrightarrow$ n column vectors in $\mathbb{R}^k$
 mod $\text{GL}_k(\mathbb{R})$ -action

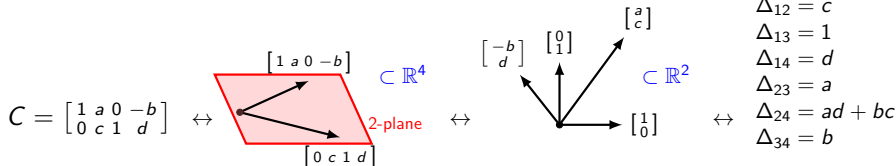


Grassmannian



Grassmannian

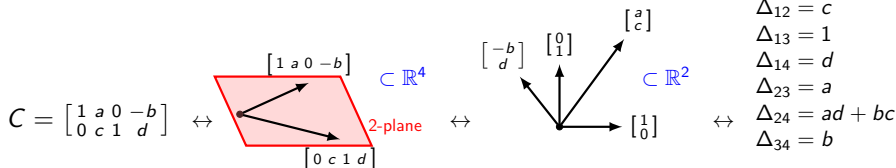
$k \times n$ matrix C mod row operations $\leftrightarrow$ ~~k row vectors in $\mathbb{R}^n$~~
 k -plane inside $\mathbb{R}^n$ $\leftrightarrow$ n column vectors in $\mathbb{R}^k$ mod $\text{GL}_k(\mathbb{R})$ -action $\leftrightarrow$ $\binom{n}{k}$ maximal minors mod rescaling



$$\text{Gr}(k, n) := \{C \in \text{Mat}(k, n)\} / (\text{row operations}) = \{k\text{-planes inside } \mathbb{R}^n\};$$

Positive Grassmannian

$k \times n$ matrix C mod row operations $\leftrightarrow$ ~~k row vectors in $\mathbb{R}^n$~~
 k -plane inside $\mathbb{R}^n$ $\leftrightarrow$ n column vectors in $\mathbb{R}^k$ mod $\mathrm{GL}_k(\mathbb{R})$ -action $\leftrightarrow$ $\binom{n}{k}$ maximal minors mod rescaling



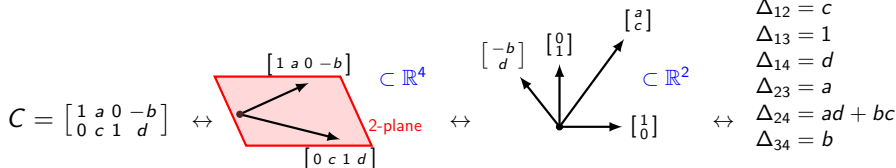
$\mathrm{Gr}(k, n) := \{C \in \mathrm{Mat}(k, n)\} / (\text{row operations}) = \{k\text{-planes inside } \mathbb{R}^n\};$

$\mathrm{Gr}_{\geq 0}(k, n) := \{C \in \mathrm{Mat}(k, n) \mid \Delta_I(C) \geq 0 \text{ for all } I \text{ of size } k\} / (\text{row operations}).$

[\[Lusztig '94\]](#), [\[Postnikov '06\]](#).

Positive Grassmannian

$k \times n$ matrix C mod row operations $\leftrightarrow$ ~~k row vectors in $\mathbb{R}^n$~~
 k -plane inside $\mathbb{R}^n$ $\leftrightarrow$ n column vectors in $\mathbb{R}^k$ mod $\text{GL}_k(\mathbb{R})$ -action $\leftrightarrow$ $\binom{n}{k}$ maximal minors mod rescaling



$\text{Gr}(k, n) := \{C \in \text{Mat}(k, n)\} / (\text{row operations}) = \{k\text{-planes inside } \mathbb{R}^n\};$

$\text{Gr}_{\geq 0}(k, n) := \{C \in \text{Mat}(k, n) \mid \Delta_I(C) \geq 0 \text{ for all } I \text{ of size } k\} / (\text{row operations}).$

[\[Lusztig '94\]](#), [\[Postnikov '06\]](#).

Cyclic symmetry: $[C_1 | C_2 | \cdots | C_n] \mapsto [C_2 | \cdots | C_n | (-1)^{k-1} C_1]$ preserves $\text{Gr}_{\geq 0}(k, n)$.

Positive kinematic space

- **Spinor-helicity formalism:** Since $P_i = (P_i^{\mathcal{T}}, P_i^{\mathcal{O}}) \in \mathbb{C}^2$ with $|P_i^{\mathcal{T}}| = |P_i^{\mathcal{O}}|$, can choose $\lambda_i, \tilde{\lambda}_i \in \mathbb{C}$ such that $P_i^{\mathcal{T}} = \lambda_i \tilde{\lambda}_i$ and $P_i^{\mathcal{O}} = \bar{\lambda}_i \tilde{\lambda}_i$.

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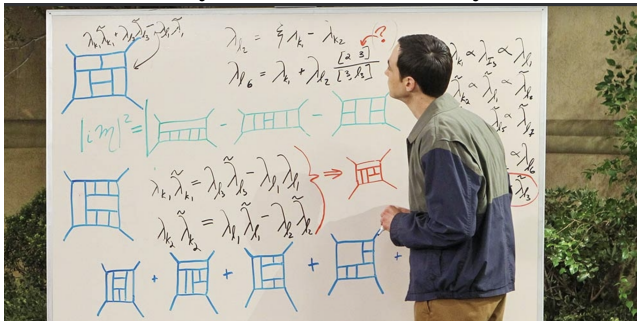
$$\left(\begin{array}{c} \frac{[4|5+6|1]^3}{[34][23]\langle 56 \rangle \langle 61 \rangle [2|3+4|5] S_{234}} \\ + \frac{[6|1+2|3]^3}{[61][12]\langle 34 \rangle \langle 45 \rangle [2|3+4|5] S_{612}} \end{array} \right) = \left(\begin{array}{c} \frac{(S_{123})^3}{[12][23]\langle 45 \rangle \langle 56 \rangle [1|2+3|4][3|4+5|6]} \\ + \frac{\langle 12 \rangle^3 [45]^3}{\langle 16 \rangle [34][3|4+5|6][5|6+1|2] S_{612}} \\ + \frac{\langle 23 \rangle^3 [56]^3}{\langle 34 \rangle [16][1|2+3|4][5|6+1|2] S_{234}} \end{array} \right)$$

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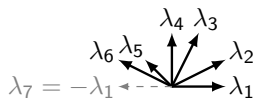
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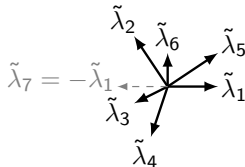
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$$\text{wind}(\lambda) = \pi$$



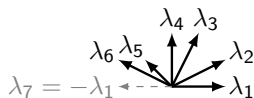
$$\text{wind}(\tilde{\lambda}) = 3\pi$$

Positive kinematic space

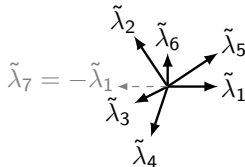
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Origami crease patterns are in natural bijection with triples $\lambda \subset C \subset \tilde{\lambda}^\perp$ such that $(\lambda, \tilde{\lambda}) \in \mathcal{K}_{k,n}^+$ and $C \in \text{Gr}_{\geq 0}(k, n).$*

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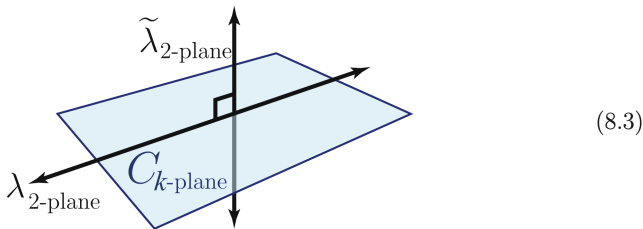
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- $(\lambda, \tilde{\lambda})$ determine the (4-dimensional) boundary of the origami crease pattern.
- $\mathcal{A}(P_1, \dots, P_n) = \text{integral over } \{C \in \text{Gr}_{\geq 0}(k, n) \mid \lambda \subset C \subset \tilde{\lambda}^\perp\}$ [ABCGPT '16].

As we saw in section 7, this can also be written as a residue of the top-form,

$$f_{\sigma}^{(k)} = \oint_{C \subset \Gamma_{\sigma}} \frac{d^{k \times n} C}{\text{vol}(GL(k))} \frac{\delta^{k \times 4}(C \cdot \tilde{\eta})}{(1 \cdots k) \cdots (n \cdots k-1)} \delta^{k \times 2}(C \cdot \tilde{\lambda}) \delta^{2 \times (n-k)}(\lambda \cdot C^{\perp}). \quad (8.2)$$

Recall from section 4, the (ordinary) δ -functions in (8.2) have the geometric interpretation of constraining the k -plane C to be *orthogonal* to the 2-plane $\tilde{\lambda}$ and to *contain* the 2-plane λ , [14]:



[Arkani-Hamed, Bourjaily, Cachazo, Goncharov, Postnikov, Trnka. *Grassmannian Geometry of Scattering Amplitudes*. Cambridge University Press, Cambridge, 2016.]

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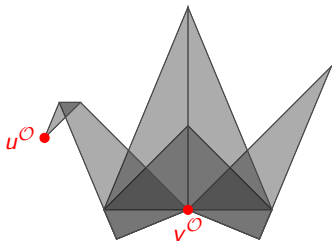
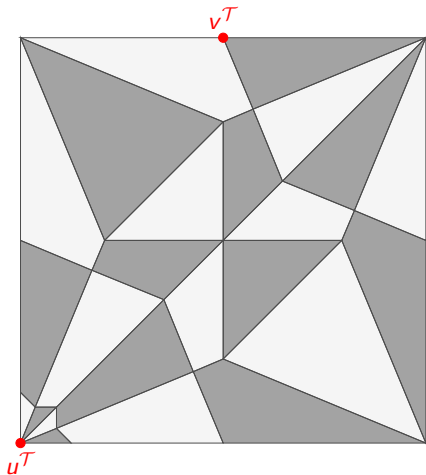
Theorem (G. (2024), “Main bijection”)

Origami crease patterns are in natural bijection with triples $\lambda \subset C \subset \tilde{\lambda}^\perp$ such that $(\lambda, \tilde{\lambda}) \in \mathcal{K}_{k,n}^+$ and $C \in \text{Gr}_{\geq 0}(k, n)$.*

*modulo Lorentz transformations, etc.

- $(\lambda, \tilde{\lambda})$ determine the (4-dimensional) boundary of the origami crease pattern.
- $\mathcal{A}(P_1, \dots, P_n) = \text{integral over } \{C \in \text{Gr}_{\geq 0}(k, n) \mid \lambda \subset C \subset \tilde{\lambda}^\perp\}$ [ABCGPT '16].
- **Corollary:** BCFW cells triangulate (Mandelstam-positive region of) $\mathcal{K}_{k,n}^+$.

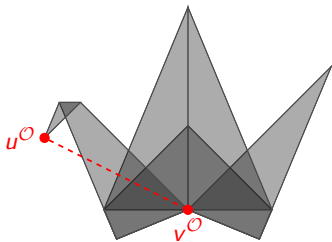
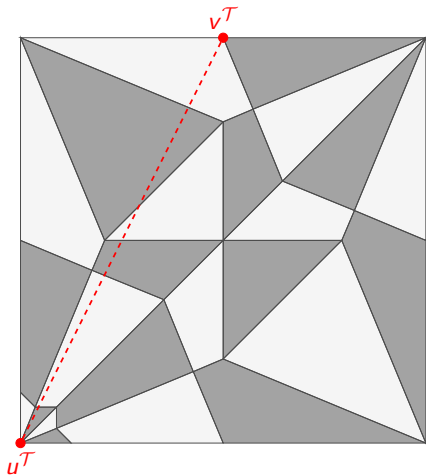
- Let u^T, v^T be two vertices of an origami crease pattern and u^O, v^O be their images under the origami map.



- Let u^T, v^T be two vertices of an origami crease pattern and u^O, v^O be their images under the origami map.

Question

True or False: we always have $|u^T - v^T| \geq |u^O - v^O|$?

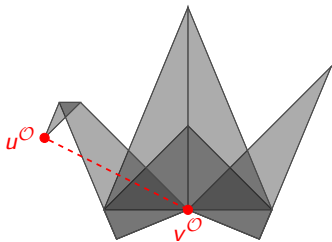
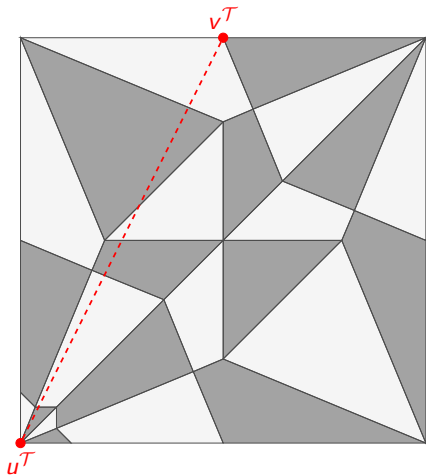


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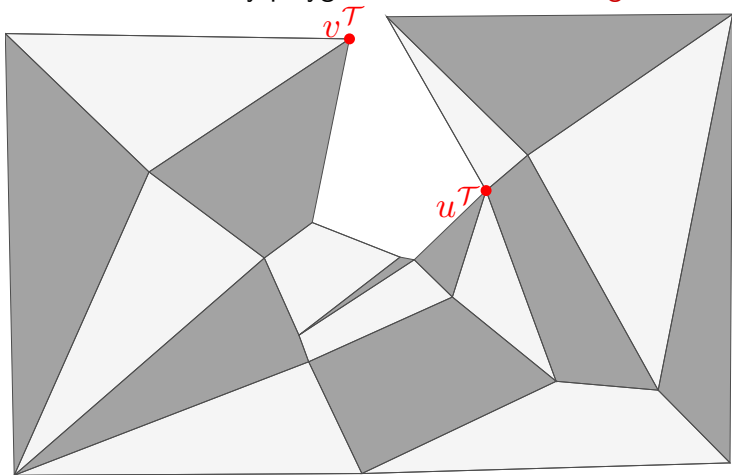


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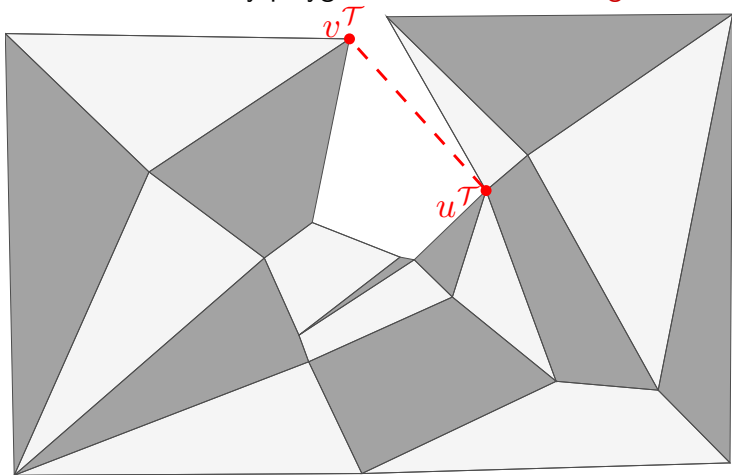


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- **Mandelstam-positive region:**

$$\mathcal{K}_{k,n}^{M+} := \left\{ \lambda \perp \tilde{\lambda} \left| \begin{array}{l} \langle i i + 1 \rangle > 0, [i i + 1] > 0 \text{ for } i = 1, \dots, n, \\ \text{wind}(\lambda) = (k - 1)\pi, \text{wind}(\tilde{\lambda}) = (k + 1)\pi, \text{ and} \\ S_{(i,j]} > 0 \text{ for all } i, j \end{array} \right. \right\}.$$

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$$\left(\begin{array}{c} \frac{[4|5+6|1]^3}{[34][23]\langle 56 \rangle \langle 61 \rangle [2|3+4|5] S_{234}} \\ + \frac{[6|1+2|3]^3}{[61][12]\langle 34 \rangle \langle 45 \rangle [2|3+4|5] S_{612}} \end{array} \right) = \left(\begin{array}{c} \frac{(S_{123})^3}{[12][23]\langle 45 \rangle \langle 56 \rangle [1|2+3|4][3|4+5|6] \langle 12 \rangle^3 [45]^3} \\ + \frac{\langle 16 \rangle [34][3|4+5|6][5|6+1|2] S_{612}}{\langle 23 \rangle^3 [56]^3} \\ + \frac{\langle 34 \rangle [16][1|2+3|4][5|6+1|2] S_{234}}{\langle 23 \rangle^3 [56]^3} \end{array} \right)$$

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Theorem (G. (2024))

BCFW cells triangulate $\mathcal{K}_{k,n}^{M^+}$.

See also: [Arkani-Hamed–Trnka '14], [Even-Zohar–Lakrec–Tessler '21],
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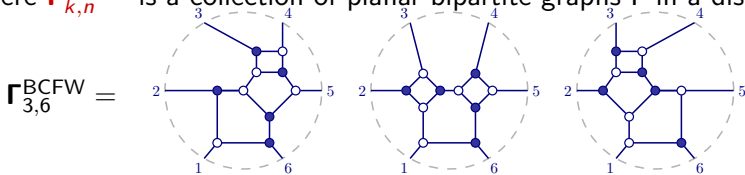
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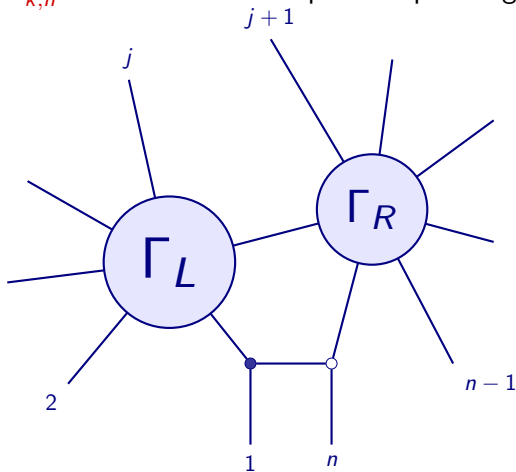


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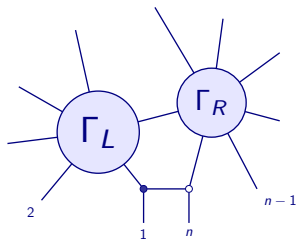
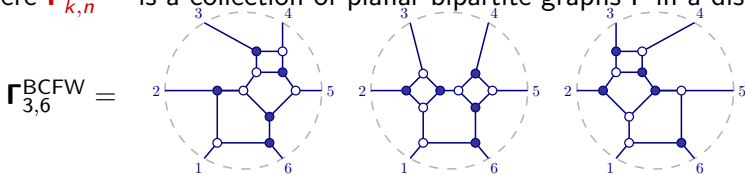


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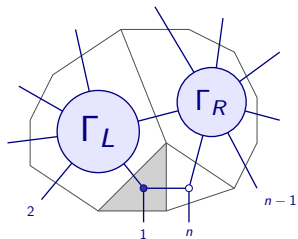
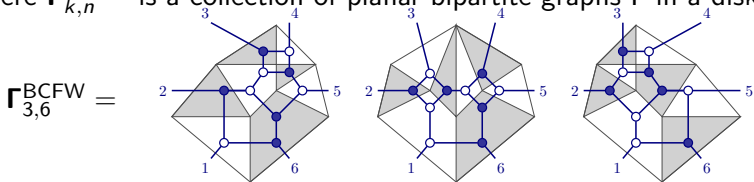


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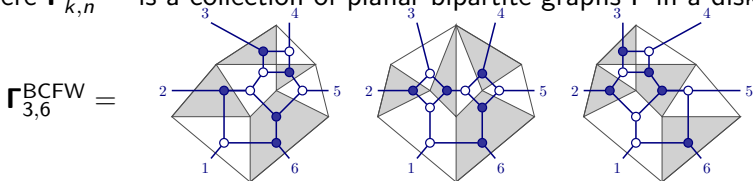


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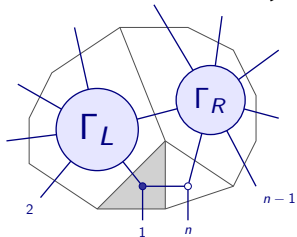
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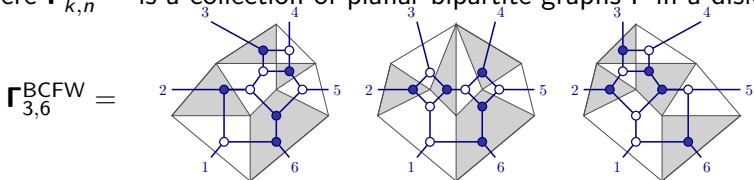


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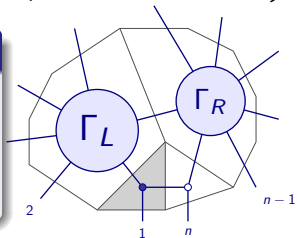
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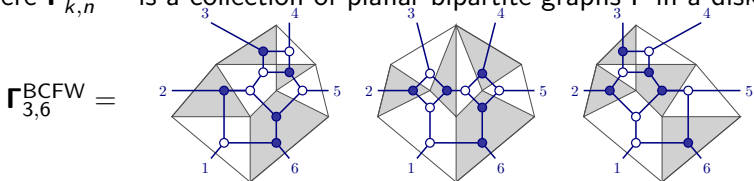


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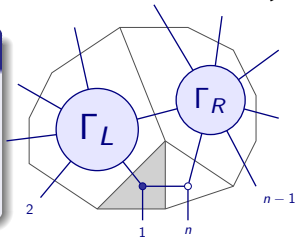
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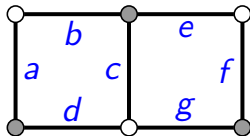
Theorem (G. (2024))

- 1 $\mathcal{K}_{k,n}^{M^+} = \bigsqcup_{\Gamma \in \Gamma_{k,n}^{\text{BCFW}}} \mathcal{K}_{\Gamma}^{M^+}$.
- 2 Each $(\lambda, \tilde{\lambda}) \in \mathcal{K}_{\Gamma}^{M^+}$ is the boundary of a **unique** origami crease pattern planar dual to Γ .



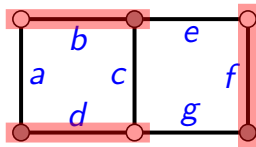
Dimer model

- **Dimer model**: study of perfect matchings on weighted bipartite graphs



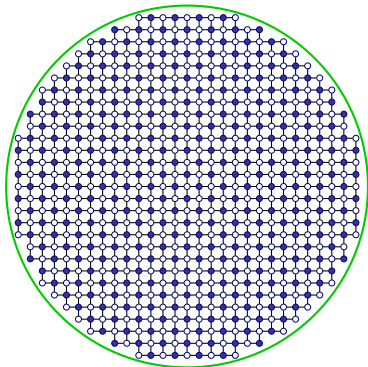
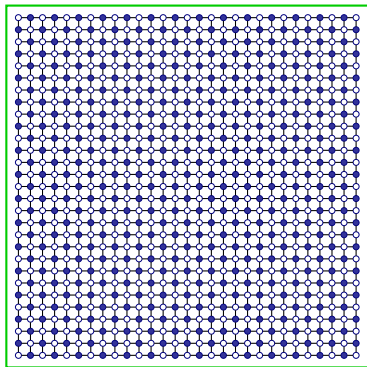
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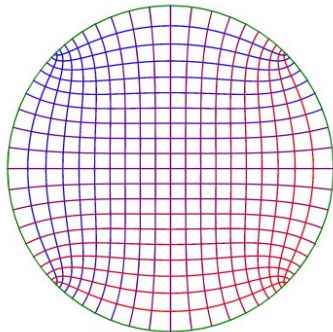
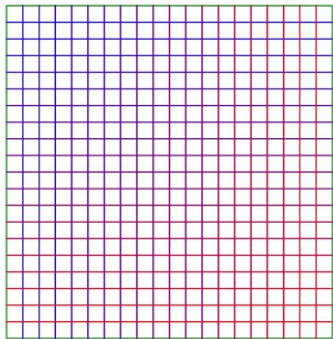
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- Conformal invariance for the **dimer model** on $\mathbb{Z}^2$ [Kenyon '00]



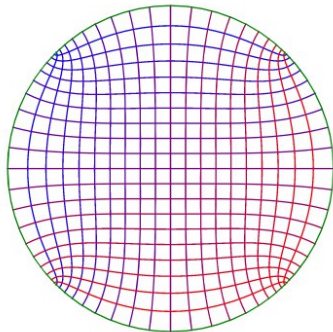
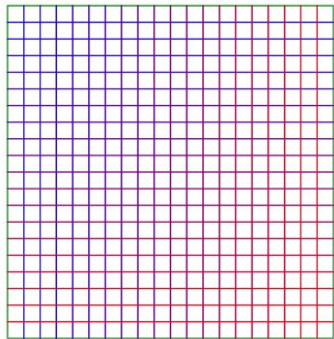
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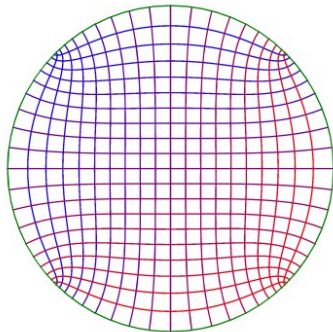
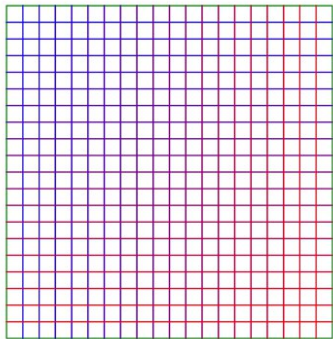
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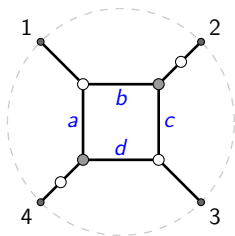
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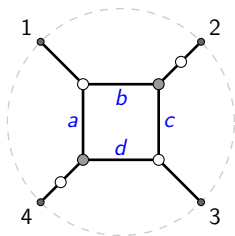
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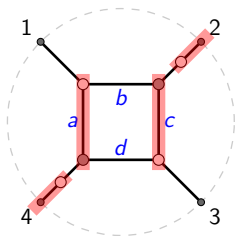
$$\Delta_{12}(C) = c \quad \Delta_{23}(C) = b \quad \Delta_{34}(C) = a \quad \Delta_{14}(C) = d$$

$$\Delta_{13}(C) = 1$$

$$\Delta_{24}(C) = ac + bd$$

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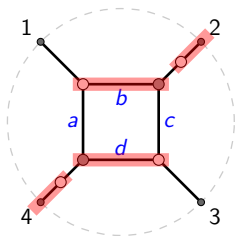
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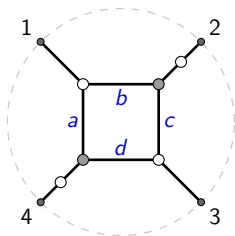
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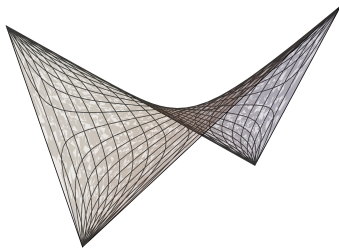
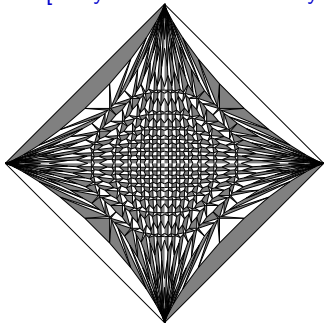
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Theorem (G. (2024))

Part 1 is true, part 2 is false.

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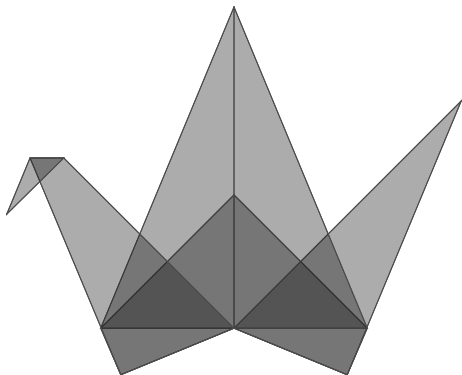
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 - **Hard direction:** if $(\lambda, \tilde{\lambda}) \in \mathcal{K}_{k,n}^+$ then $\int \lambda^\circ \tilde{\lambda}^\bullet dz$ gives a valid (embedded) origami crease pattern.



Thanks!