

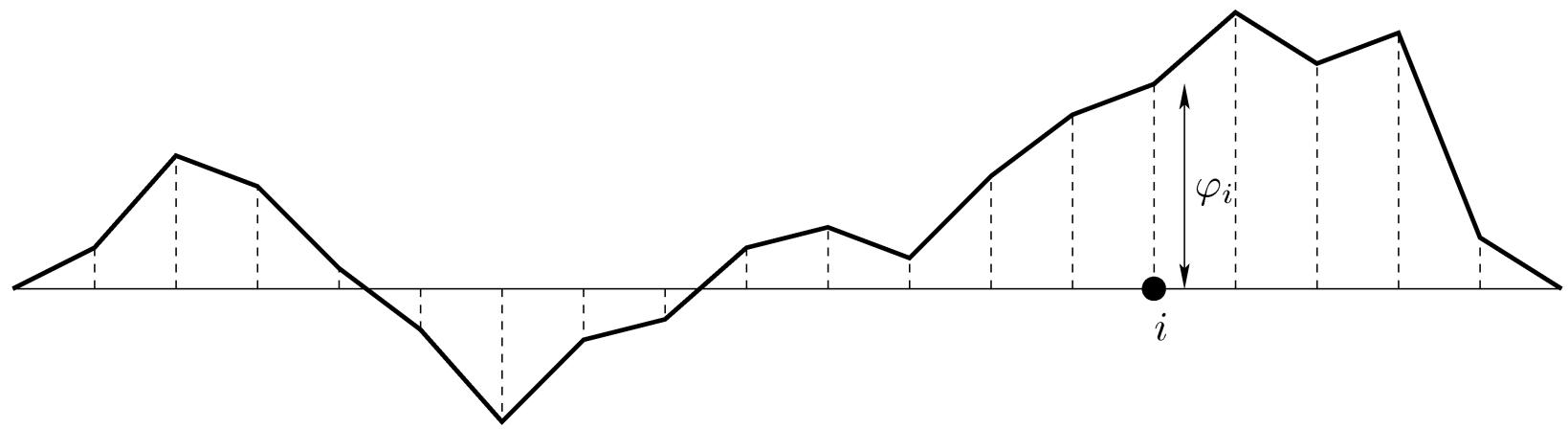
Uniqueness of random gradient states

Codina Cotar

(Based on joint work with Christof Külske)

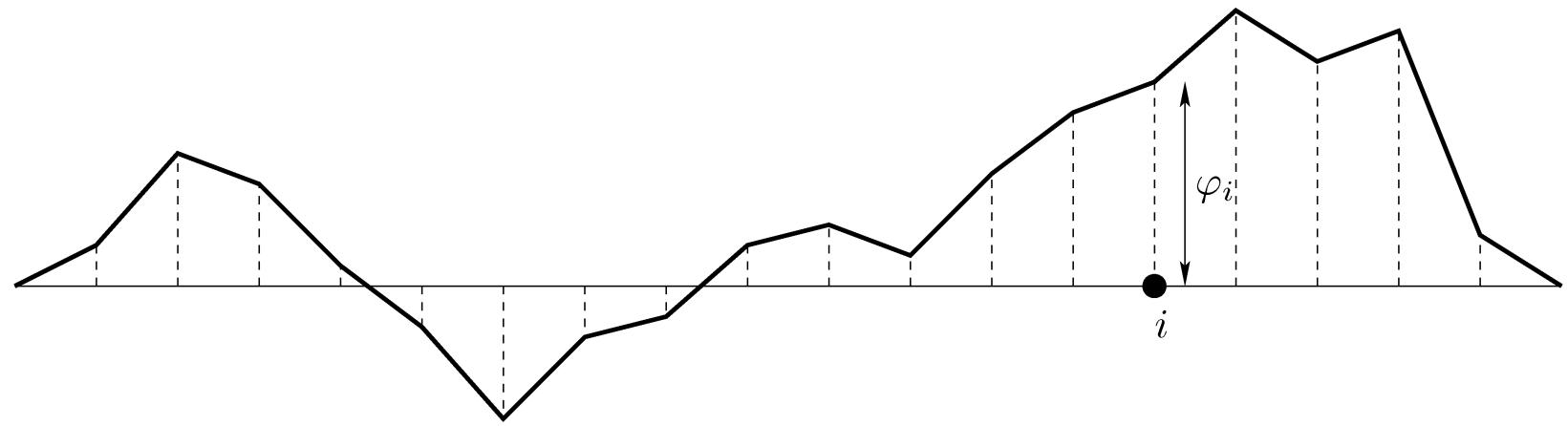
Model

- Interface — transition region that separates different phases



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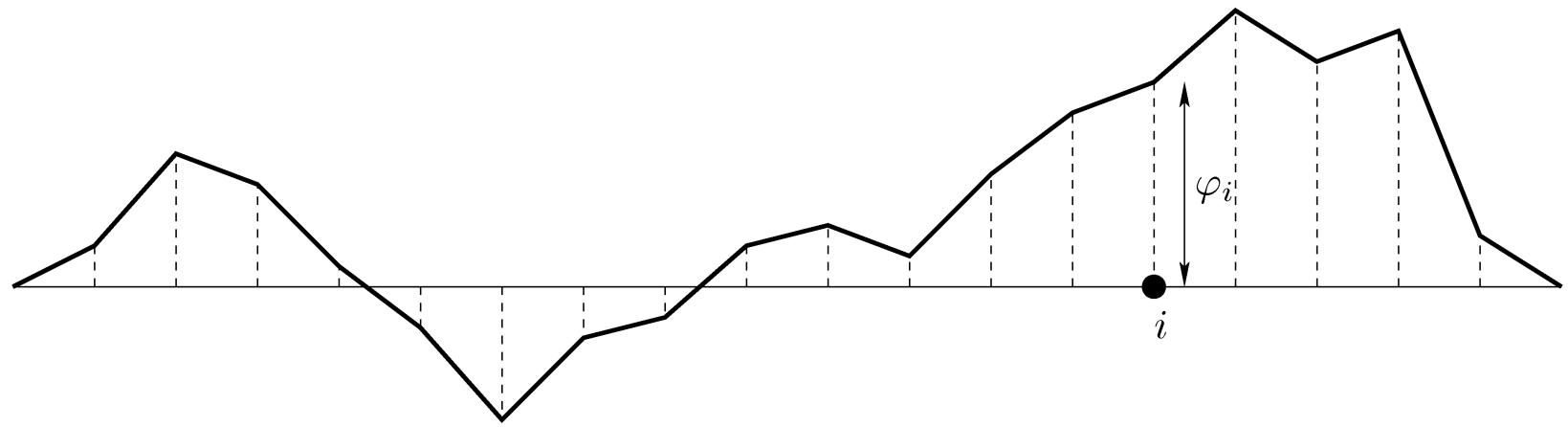
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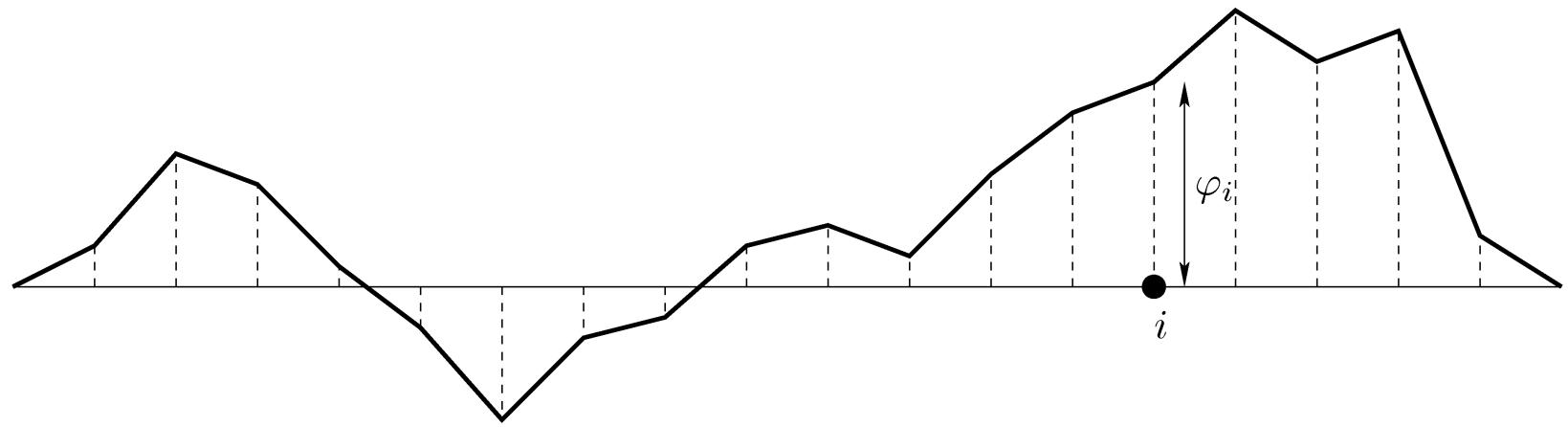
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- Finite set $\Lambda \subset \mathbb{Z}^d$
- Height Variables $\phi: \Lambda \rightarrow \mathbb{R}$
- Boundary $\partial\Lambda$ with boundary condition ψ , such that

$$\phi_x = \psi_x, \text{ when } x \in \partial\Lambda.$$

- **Potentials** $V : \mathbb{R} \rightarrow \mathbb{R}$, even, diverges at $\pm\infty$
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- **Hamiltonian**

$$H_{\Lambda}^{\psi}(\phi) = \sum_{\substack{x \in \Lambda, y \in \Lambda, \\ |x-y|=1}} V(\phi_x - \phi_y) + 2 \sum_{\substack{x \in \Lambda, y \in \partial\Lambda, \\ |x-y|=1}} V(\phi_x - \phi_y)$$

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- $\beta = \frac{1}{T} > 0$, where T is the temperature
- Finite volume **surface tension** $\sigma_{\Lambda}(u)$: macroscopic energy of a surface with tilt $u = (u_1, \dots, u_d) \in \mathbb{R}^d$.

ϕ -Gibbs Measure

- The **finite volume Gibbs measure** in Λ

$$\nu_{\Lambda}^{\psi}(d\phi) = \frac{1}{Z_{\Lambda}^{\psi}} \exp \left\{ -\beta H_{\Lambda}^{\psi}(\phi) \right\} \prod_{x \in \Lambda} d\phi_x.$$

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- The probability measure $\nu \in P(\mathbb{R}^{\mathbb{Z}^d})$ is called a **Gibbs measure** for the ϕ -field if

$$\nu(\cdot | \mathcal{F}_{\Lambda^c})(\psi) = \nu_{\Lambda}^{\psi}(\cdot), \quad \nu - \text{a.e. } \psi,$$

for every $\Lambda \subset \mathbb{Z}^d$, where \mathcal{F}_{Λ^c} is the σ -field of $\mathbb{R}^{\mathbb{Z}^d}$ generated by $\{\phi(x) : x \notin \Lambda\}$.

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- $\nabla\phi$ -Gibbs measures exist for $d \geq 1$.

Questions

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6. Large deviations (LDP) results.

Results: Strictly Convex Potentials

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3. **Funaki-Spohn (CMP, 1997):** Assume $C_1 \leq V'' \leq C_2$. For every $u = (u_1, \dots, u_d) \in \mathbb{R}^d$ there exists a unique shift-invariant $\nabla\phi$ -Gibbs measure μ with $E_\mu[\phi(e_i) - \phi(0)] = u_i$, for all $i = 1, \dots, d$.

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$$\epsilon^{\frac{d}{2}} \sum_{x \in \mathbb{Z}^d} f(\epsilon x) [\phi_{x+e_i} - \phi_x - u_i] \xrightarrow{\epsilon \rightarrow 0} N(0, \sigma_u^2(f)).$$

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Techniques

- Brascamp-Lieb Inequality: for all $x \in \Lambda$

$$\text{var}_{\nu_\Lambda^\psi}(\phi_x) \leq \text{var}_{\tilde{\nu}_\Lambda^\psi}(\phi_x),$$

$\tilde{\nu}_\Lambda^\psi$ is the Gibbs measure with potential
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- Random Walk Representation: Representation of the Covariance Matrix in terms of the Green function of a particular random walk.

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$$\begin{aligned} V &= V_0 + g, \quad C_1 \leq V_0'' \leq C_2, \\ -C_0 &\leq g'' < 0, \quad \sqrt{\beta} \|g''\|_{L^1(\mathbb{R})} \text{ small } (C_0, C_1, C_2). \end{aligned}$$

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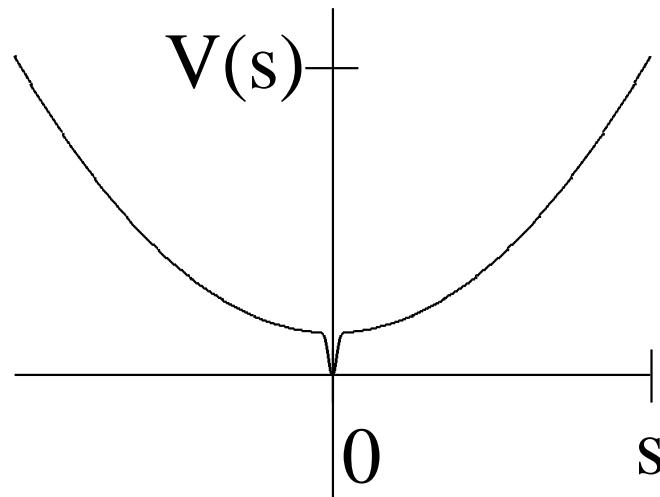
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- **Adams-Kotecký-Müller (preprint, 2011):** Strict convexity of the surface tension for small tilts u and large β .

Non-Convex potentials

- For the potential

$$e^{-V(s)} = pe^{-k_1 \frac{s^2}{2}} + (1-p)e^{-k_2 \frac{s^2}{2}}$$



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- **Biskup-Spohn (AOP, 2011):** CLT for every ergodic $\nabla\phi$ -Gibbs measure μ with $E_\mu[\phi(e_i) - \phi(0)] = 0$.

Non-Convex potentials

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Interfaces with Disorder

Two ways to add randomness

(a) Add to the Hamiltonian the random part

$$\sum_{x \in \Lambda} \xi_x \phi_x$$

$(\xi_x)_{x \in \mathbb{Z}^d}$ i.i.d with $E_{\mathbb{P}}(\xi_x^2) < \infty$.

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Or

(b) Make the potentials random: $V_{\langle x,y \rangle}$ i.i.d with
 $A s^2 - B_{(x,y)} \leq V_{(x,y)}(s) \leq C_2 s^2$ and
 $\mathbb{E}|B_{\langle x,y \rangle}| < \infty$.

Results

- Translation-covariance:

$$\int \mu[\tau_v(\xi)](d\eta) F(\eta) = \int \mu[\xi](d\eta) F(\tau_v \eta),$$

for $(\tau_v \eta)(b) := \eta(b - v)$, for all bonds b and all $v \in \mathbb{Z}^d$.

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