

Euler's method for the approximate solution of ODE's

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Euler's method for a single ordinary differential equation.

If one is given the ODE

$$\frac{dy}{dt} = f(y, t) \quad y(t_0) = y_0 \quad (1)$$

and asked to approximate the solution over a time interval $[t_0, T]$, then a common way to do this is to create a *difference* equation whose solution y_k approximates $y(t_k)$ at some finite set of times $t_k \subseteq [t_0, T]$ for $k = 0 \dots N$ (see Figure 1).

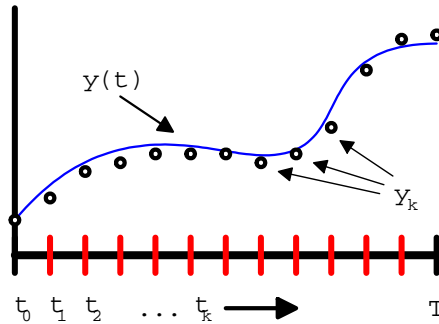


Figure 1

One of the simplest numerical methods based on this idea is Euler's method. In Euler's method the difference equation

$$y_{k+1} = y_k + \Delta t f(y_k, t_k) \quad y_0 = y(0) \quad (2)$$

is used to create the approximate solution values. Here $\Delta t = \frac{T - t_0}{N}$, where N is the number of timesteps to be taken.

Euler's method can be "derived" or "understood" in a number of different ways. One way to understand Euler's method is to recognize it as a repeated application of a first order Taylor series approximation to the solution. Specifically, if we know $y(t)$ at time $t = t_k$ and we want to approximate it at time $t = t_{k+1}$, then we can use a Taylor's series expansion with $\Delta t = t_{k+1} - t_k$:

$$y(t_{k+1}) = y(t_k) + \Delta t \left(\frac{dy}{dt} \right)_{t_k} + \frac{\Delta t^2}{2} \left(\frac{d^2y}{dt^2} \right)_{t_k} + \dots$$

Assuming that the timestep Δt is small enough, we neglect the higher order terms and use

$$y(t_{k+1}) \approx y(t_k) + \Delta t \left(\frac{dy}{dt} \right)_{t_k} \quad (3)$$

To use (3) requires knowing the derivative of the function at t_k , $\left(\frac{dy}{dt} \right)_{t_k}$. However, $y(t)$ satisfies the differential equation (1) so we have an explicit formula for its derivative at time t_k in terms of its value at time t_k , namely, $\left(\frac{dy}{dt} \right)_{t_k} = f(y(t_k), t_k)$. Substituting this relation into (3) gives us

$$y(t_{k+1}) \approx y(t_k) + \Delta t f(y(t_k), t_k). \quad (4)$$

Hence we see that Euler's method (2) is the repeated application of a first order Taylor series approximation (4). Now, one might worry that the errors in the repeated use of (4) accumulate in a catastrophic manner. In courses in numerical analysis, one proves that the errors tend to zero as the number of timesteps (over a fixed time interval) increases. In fact, under suitable assumptions on the differential equation, the difference between the computed and exact solution obeys the relation

$$\max_k |y(t_k) - y_k| \leq C \Delta t \quad k = 0 \dots N$$

for some constant C . Thus, as the number of steps $N \rightarrow \infty$, $\Delta t \rightarrow 0$, and so the approximate solution values converge to the true solution values.

Euler's method for systems.

Systems of ordinary differential equations have the form

$$\frac{d}{dt} \begin{pmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{pmatrix} = \begin{pmatrix} f_1(y_1, y_2, \dots, y_n, t) \\ f_2(y_1, y_2, \dots, y_n, t) \\ \vdots \\ f_n(y_1, y_2, \dots, y_n, t) \end{pmatrix} \quad \begin{pmatrix} y_1(t_0) \\ y_2(t_0) \\ \vdots \\ y_n(t_0) \end{pmatrix} = \begin{pmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{pmatrix}_0 \quad (5)$$

or in vector notation

$$\frac{d\vec{y}}{dt} = \vec{F}(\vec{y}, t) \quad \vec{y}(t_0) = \vec{y}_0$$

The solution one seeks is a vector valued function $\vec{y}(t)$ whose time derivative satisfies (5). Euler's method extends directly to systems of differential equations. Formally, Euler's method is obtained by just letting $y_k \rightarrow \vec{y}_k$, $y_{k+1} \rightarrow \vec{y}_{k+1}$, and $f \rightarrow \vec{F}$ in (2);

$$\vec{y}_{k+1} = \vec{y}_k + \Delta t \vec{F}(\vec{y}_k, t_k) \quad \vec{y}_0 = \vec{y}(t_0) \quad (6)$$

Here $\Delta t = \frac{T - t_0}{N}$, where N is the number of timesteps to be taken.

Just as in the single ODE case, one can "understand" Euler's method for systems as being the repeated application of a Taylor series approximation. Specifically, given $\vec{y}(t_k)$ one can think of approximating $\vec{y}(t_{k+1})$ using a Taylor Series approximation for each component,

$$\begin{aligned}
y_1(t_{k+1}) &= y_1(t_k) + \Delta t \left(\frac{dy_1}{dt} \right)_{t_k} + \frac{\Delta t^2}{2} \left(\frac{d^2 y_1}{dt^2} \right)_{t_k} + \dots \\
y_2(t_{k+1}) &= y_2(t_k) + \Delta t \left(\frac{dy_2}{dt} \right)_{t_k} + \frac{\Delta t^2}{2} \left(\frac{d^2 y_2}{dt^2} \right)_{t_k} + \dots \\
&\vdots \\
y_n(t_{k+1}) &= y_n(t_k) + \Delta t \left(\frac{dy_n}{dt} \right)_{t_k} + \frac{\Delta t^2}{2} \left(\frac{d^2 y_n}{dt^2} \right)_{t_k} + \dots
\end{aligned} \tag{7}$$

If we assume a sufficiently small timestep, then we use just the first two terms of each approximation

$$\begin{aligned}
y_1(t_{k+1}) &\approx y_1(t_k) + \Delta t \left(\frac{dy_1}{dt} \right)_{t_k} \\
y_2(t_{k+1}) &\approx y_2(t_k) + \Delta t \left(\frac{dy_2}{dt} \right)_{t_k} \\
&\vdots \\
y_n(t_{k+1}) &\approx y_n(t_k) + \Delta t \left(\frac{dy_n}{dt} \right)_{t_k}
\end{aligned}$$

or in vector notation,

$$\vec{y}(t_{k+1}) \approx \vec{y}(t_k) + \Delta t \left(\frac{d\vec{y}}{dt} \right)_{t_k} \tag{8}$$

As before, we need the derivatives of each component to evaluate this approximation. The differential equations provide us with a means of obtaining the derivatives, namely, $\left(\frac{d\vec{y}}{dt} \right)_{t_k} = \vec{F}(\vec{y}(t_k), t_k)$.

Thus, using this expressions in (8) we obtain

$$\vec{y}(t_{k+1}) \approx \vec{y}(t_k) + \Delta t \vec{F}(\vec{y}(t_k), t_k)$$

So Euler's method for a system is just repeated application of a Taylor series approximation to each component. As with single ordinary differential equations, one finds that under suitable assumptions on the differential equation the numerical solutions \vec{y}_k converge to $\vec{y}(t_k)$, in fact,

$$\max_k \|\vec{y}(t_k) - \vec{y}_k\| \leq C \Delta t \quad k = 0 \dots N$$

for some constant C . Thus, as the number of steps $N \rightarrow \infty$, $\Delta t \rightarrow 0$, and the approximate solution vector converges to the true solution vector at each t_k .