# CONVEX HYPERSURFACE THEORY IN CONTACT TOPOLOGY 

KO HONDA AND YANG HUANG


#### Abstract

We lay the foundations of convex hypersurface theory (CHT) in contact topology, extending the work of Giroux in dimension three. Specifically, we prove that any closed hypersurface in a contact manifold can be $C^{0}$ approximated by a convex one. We also prove that a $C^{0}$-generic family of mutually disjoint closed hypersurfaces parametrized by $t \in[0,1]$ is convex except at finitely many times $t_{1}, \ldots, t_{N}$, and that crossing each $t_{i}$ corresponds to a bypass attachment. As applications of CHT, we prove the existence of compatible (relative) open book decompositions for contact manifolds and an existence $h$-principle for codimension 2 contact submanifolds.


## Contents

1. Introduction ..... 2
1.1. Convex contact structures ..... 2
1.2. Main results ..... 4
1.3. Applications ..... 5
2. A convexity criterion ..... 8
2.1. Characteristic foliations ..... 8
2.2. Morse and 1-Morse vector fields ..... 9
2.3. A convexity criterion ..... 11
3. Construction of mushrooms in dimension 3 ..... 13
3.1. Construction of $Z_{\mathrm{PL}}$ ..... 13
3.2. Smoothing of $Z_{\text {PL }}$ ..... 15
4. Convex surface theory revisited ..... 18
4.1. 3-dimensional plugs ..... 18
4.2. Barricades ..... 19
4.3. Proof of Theorem 1.2.3 in dimension 3 ..... 20
4.4. Installing and uninstalling plugs ..... 20
4.5. Proof of Theorem 1.2 .5 in dimension 3 ..... 22
4.6. Further remarks ..... 24
5. Construction of mushrooms in dimension $>3$ ..... 25
5.1. Introduction ..... 25
5.2. Product hypersurface ..... 26
5.3. Dynamics of $Z_{\xi}^{\prime}$ ..... 27
5.4. Damping ..... 29
5.5. Description of the characteristic foliation of the mushroom ..... 31
6. Quantitative stabilization of an open book decomposition for $S^{2 n-1}$ ..... 32
6.1. Some definitions ..... 32
6.2. Quantitative stabilization for $S^{2 n-1}$ ..... 34
6.3. Verification of the Weinstein property ..... 35
7. Construction of the plug ..... 36
7.1. Definition of a plug ..... 36
7.2. A Peter-Paul contactomorphism ..... 37
7.3. A pre-plug ..... 38
7.4. $\epsilon$-short hypersurfaces ..... 41
7.5. Construction of the plug ..... 41
8. Bifurcations of characteristic foliations and bypass attachments ..... 44
8.1. Definitions and examples ..... 44
8.2. Normalization of contact structure near a folded Weinstein hypersurface ..... 46
8.3. Bypass attachment as a bifurcation ..... 48
8.4. Bypass attachment in terms of partial open book decompositions ..... 56
9. $C^{0}$-approximation by convex hypersurfaces ..... 58
10. The existence of (partial) open book decompositions ..... 59
11. Applications to contact submanifolds ..... 62
11.1. Models for pushing across mushrooms ..... 62
11.2. Existence $h$-principles ..... 67
Appendix A. Wrinkled and folded embeddings ..... 70
A.1. Wrinkled and cuspidal embeddings ..... 70
A.2. Cuspidal embeddings of a disk ..... 73
A.3. Folding hypersurfaces ..... 74
Appendix B. Quantitative stabilizations of open book decompositions in general ..... 75
B.1. Definitions ..... 75
B.2. Perturbing to a strongly adapted contact form ..... 76
B.3. Damped OBD ..... 78
B.4. Quantitative stabilization of OBD ..... 80
References ..... 90

## 1. Introduction

1.1. Convex contact structures. Morse theory is a topologist's favorite tool for exploring the structure of manifolds. The significance of Morse theory - here we mean the traditional finite-dimensional version, not Floer theory - in contact and symplectic topology was advocated by Eliashberg and Gromov in [EG91]. In particular, according to [EG91, Definition 3.5.A], a contact manifold $(M, \xi)$ is convex if there exists a Morse function, called a contact Morse function, which admits a gradient-like vector field whose flow preserves $\xi$. Just as a manifold can
be reconstructed from its Morse function by a sequence of handle attachments in traditional Morse theory, a contact manifold can be reconstructed from a contact Morse function by a sequence of contact handle attachments. The analogous theory in symplectic topology is known as the theory of Weinstein manifolds.

Eliashberg and Gromov asked in [EG91] whether there exist non-convex contact manifolds. Around 2000 Giroux gave a negative answer to the question by showing that every closed contact manifold is convex; see [Gir02]. This can also be formulated as his celebrated correspondence between contact structures and open book decompositions. This is in sharp contrast to the theory of Weinstein manifolds, where it is relatively easy to see that any compatible Morse function cannot have critical points of index greater than half of the dimension of the manifold.
Remark 1.1.1. It might be the case that "convexity" is one of the most abused terminologies in mathematics. We will not use the term "convex contact manifold" in the sense of Eliashberg and Gromov for the rest of the paper.

At this point, the question bifurcates into two:
Question 1.1.2. How do we establish Morse theory on contact manifolds?
Question 1.1.3. How do we use Morse theory to better understand contact manifolds?

Let us first address Question 1.1.2, which was first answered by Giroux in both dimension 3 and in higher dimensions. Giroux used two completely different sets of techniques to treat the 3 -dimensional and higher-dimensional cases.

We first discuss the 3 -dimensional case. In his thesis [Gir91], Giroux introduced what is now known as convex surface theory into 3 -dimensional contact topology. It is an extremely powerful and efficient way of studying embedded surfaces in contact 3 -manifolds, and can recover most of the pioneering results of Bennequin [Ben83] and Eliashberg [Eli92]. Using convex surface theory, Giroux showed that for closed contact 3 -manifolds, there is a one-to-one correspondence between isotopy classes of contact structures and compatible open book decompositions up to positive stabilization.

Before moving onto higher dimensions, let us recall the definition of a convex hypersurface following [Gir91]:
Definition 1.1.4. A hypersurface $\Sigma \subset(M, \xi)$ is convex if there exists a contact vector field $v$, i.e., a vector field whose flow preserves $\xi$, which is transverse to $\Sigma$ everywhere.

Observe that regular level sets of a contact Morse function are convex hypersurfaces.

The situation in dimensions $>3$ is quite different. Besides the fact that convex hypersurfaces can be defined in any dimension, until now there has been no systematic convex hypersurface theory. Giroux's proof [Gir02, Gir17] that every closed contact manifold is convex involves a completely different technology, i.e., Donaldson's [Don96] technique of approximately holomorphic sections, transplanted into contact topology independently by Ibort, Martínez-Torres, and Presas
[IMTP00] and by Mohsen [Moh, Moh19]. Donaldson used the approximate holomorphic technology to construct real codimension 2 symplectic hypersurfaces of a closed symplectic manifold as the zero locus of an approximately holomorphic section of a complex line bundle, while [IMTP00] and [Moh, Moh19] constructed certain codimension 2 contact submanifolds of a closed contact manifold. What Giroux realized is that [Don96], [IMTP00], and [Moh, Moh19] could be used to produce compatible open book decompositions. Roughly speaking, given a closed contact manifold $(M, \xi=\operatorname{ker} \alpha)$, one considers the trivial line bundle $\mathbb{C}$ on $M$ equipped with a suitable Hermitian connection determined by $\alpha$. Then there exists a section $s: M \rightarrow \mathbb{C}$ whose zero locus $B:=s^{-1}(0)$ is a closed codimension 2 contact submanifold called the binding, and

$$
\frac{s}{|s|}: M \backslash B \rightarrow S^{1}
$$

is a smooth fibration defining the compatible open book decomposition of $(M, \xi)$. As a consequence of using the approximate holomorphic technology, the higherdimensional Giroux correspondence (see Corollary 1.3.1) is a much weaker statement compared to its 3 -dimensional counterpart.
1.2. Main results. The main goal of this paper is to systematically generalize Giroux's convex surface theory to all dimensions. The main results of convex hypersurface theory (CHT) are Theorems 1.2.3 and 1.2.5. In fact, even in dimension 3, our method (see Section 4) differs somewhat from Giroux's original approach, is simpler, and is consistent with our more general approach in higher dimensions.

We first introduce some more terminology describing the anatomy of a convex hypersurface.

Definition 1.2.1. Let $\Sigma \subset(M, \xi=\operatorname{ker} \alpha)$ be a convex hypersurface with respect to a transverse contact vector field $v$. Define the dividing set $\Gamma(\Sigma):=\{\alpha(v)=0\}$ and $R_{ \pm}(\Sigma):=\{ \pm \alpha(v)>0\}$ as subsets of $\Sigma$.

It turns out that $\Gamma(\Sigma) \subset(M, \xi)$ is a codimension 2 contact submanifold, and $R_{ \pm}(\Sigma)$ are (complete) Liouville manifolds with Liouville form given by a suitable rescaling of $\left.\alpha\right|_{R_{ \pm}(\Sigma)}$, respectively. Moreover, the isotopy classes of $\Gamma(\Sigma), R_{ \pm}(\Sigma)$ are independent of the choices of $v$ and $\alpha$.

In dimensions $\geq 4$, there exist Liouville manifolds that are not Weinstein by McDuff [McD91], Geiges [Gei94, Gei95], Mitsumatsu [Mit95], and Massot, Niederkrüger, and Wendl [MNW13]. While these "exotic" Liouville manifolds are great for constructing (counter-)examples, there currently is no systematic understanding of such non-Weinstein Liouville manifolds, partially because of the lack of an appropriate Morse theory on such manifolds. This is slowly starting to change: For example it was recently shown in [EOY] that a stabilized Liouville manifold with the homotopy type of a half-dimensional CW-complex is symplectomorphic to a flexible Weinstein manifold and in $[\mathrm{BC}]$ that the stabilizations of Mitsumatsu's Liouville domains are Weinstein domains. ${ }^{1}$ This motivates the following definition:

[^0]Definition 1.2.2. A convex hypersurface $\Sigma$ is Weinstein (resp. 1-Weinstein) convex if $R_{+}(\Sigma)$ and $R_{-}(\Sigma)$ are Weinstein (resp. 1-Weinstein).

Here a Liouville domain that admits a Liouville vector field which is gradientlike with respect to a "1-Morse" function which also admits critical points of birthdeath type) will be called a 1-Weinstein domain (instead of a generalized Weinstein domain) in this paper.

Weinstein convex hypersurfaces admit a Morse-theoretic interpretation, given in Proposition 2.3.3.

Now we are ready to state the foundational theorems of CHT.
Theorem 1.2.3. Any closed oriented hypersurface in a contact manifold can be $C^{0}$-approximated by a Weinstein convex one.

Remark 1.2.4. The $C^{\infty}$-version of Theorem 1.2.3 remains open. Mori's candidate counterexample in [Mor] was shown by Breen [Bre21] to actually admit a $C^{\infty}$. small approximation by a convex one.
Theorem 1.2.5. Let $\xi$ be a contact structure on $\Sigma \times[0,1]$ such that the hypersurfaces $\Sigma \times\{0,1\}$ are Weinstein convex. Then, up to a boundary-relative contact isotopy, there exists a finite sequence $0<t_{1}<\cdots<t_{N}<1$ such that the following hold:

- $\Sigma \times\{t\}$ is 1 -Weinstein convex if $t \neq t_{i}$ for any $1 \leq i \leq N$.
- For each $i$, there exists a small $\epsilon>0$ such that $\xi$ restricted to $\Sigma \times\left[t_{i}-\right.$ $\left.\epsilon, t_{i}+\epsilon\right]$ is contactomorphic to a bypass attachment.

For an initial study of bypass attachments in higher dimensions the reader is referred to [HH].

Remark 1.2.6. Theorem 1.2.5 was conjectured by Paolo Ghiggini in the afternoon of April 10, 2015 in Paris.
1.3. Applications. As an immediate application of Theorems 1.2.3 and 1.2.5, we can extend Giroux's 3-dimensional approach to constructing compatible open book decompositions to higher dimensions. This is the content of the following three corollaries. Note, however, that we do not address the stabilization equivalence of the compatible open book decompositions in this paper. We plan to investigate this in future work.

Corollary 1.3.1 ([Gir02]). Any closed contact manifold admits a compatible open book decomposition, all of whose pages are 1-Weinstein.

Corollary 1.3.2. Any compact contact manifold with convex boundary admits a compatible partial open book decomposition, all of whose pages are 1-Weinstein domains or 1-Weinstein cobordisms.

Corollary 1.3.3. Given a possibly disconnected closed Legendrian submanifold $\Lambda$ in a closed contact manifold, there exists a compatible open book decomposition such that all the pages are 1 -Weinstein and $\Lambda$ is contained in a page.

Remark 1.3.4. Although Corollary 1.3.3 does not appear in the literature, asymptotically holomorphic techniques in contact geometry from [IMTP00] and [Moh, Moh19] - more specifically the combination of Giroux's existence theorem of compatible open book decompositions (see Presas [Pre14] for a published proof) and asymptotically holomorphic symplectic special position theorems for Lagrangians in symplectic manifolds [AMnP05] applied to the Lagrangianization of a Legendrian in the symplectization gives the result.

This completes our exploration of Question 1.1.2 for the time being.
Next we turn to Question 1.1.3, which is a much harder question. For example, we would like to obtain classification results for contact structures on higherdimensional manifolds (e.g., the spheres) besides the "flexible" ones due to Borman, Eliashberg, and Murphy [BEM15]. Unfortunately, our current understanding of contact Morse theory is not good enough for us to classify anything in higher dimensions. Instead, we will use the (mostly dynamical) techniques developed in this paper to address the existence problems of contact manifolds and submanifolds.

The existence problems of contact manifolds and submanifolds were first addressed by Gromov [Gro86] using his magnificent zoo of $h$-principles. In particular, he proved a full $h$-principle for contact structures on open manifolds (see [EM02, 10.3.2]) and an existence $h$-principle for isocontact embeddings $Y \subset$ $(M, \xi)$ under the assumptions that either $Y$ has $\operatorname{codim} Y \geq 4$ or is open with $\operatorname{codim} Y=2$.

The existence problem turned out to be much harder for closed manifolds. In dimension 3, an existence $h$-principle for contact structures was proved by Martinet [Mar71] and Lutz [Lut77]. For overtwisted contact 3-manifolds, a full $h$-principle was proved by Eliashberg [Eli89]. In dimension 5, there is a rich literature of partial results: the existence of contact structures on certain classes of 5-manifolds was established by Geiges [Gei91, Gei97], Geiges-Thomas [GT98, GT01], and Bourgeois [Bou02]. Afterwards, a complete existence $h$-principle for contact 5manifolds was established by Casals, Pancholi, and Presas [CPP15] (there is also the approach of of Etnyre [Etn], which currently has a gap). Finally, the existence $h$-principle for contact manifolds of any dimension, as well as the full $h$-principle for overtwisted contact manifolds of any dimension, was established by Borman, Eliashberg, and Murphy [BEM15].

So far the story has mostly been about the contact manifolds themselves. Now we turn to the existence problem of contact submanifolds or (iso-)contact embeddings. Results in this direction were surprisingly rare until recently. Besides the aforementioned $h$-principle of Gromov, there exist constructions of contact submanifolds by Ibort, Martínez-Torres, and Presas [IMTP00], mentioned earlier. In low dimensions, there also exist works by Kasuya [Kas16], Etnyre-Furukawa [EF17], and Etnyre-Lekili [EL] on embedding contact 3-manifolds into certain contact 5-manifolds.

In the rest of the introduction we will explain the existence $h$-principle for codimension 2 contact submanifolds. Since the case of open submanifolds has already
been settled by Gromov, we may assume that all the submanifolds involved are closed.

Definition 1.3.5. Let $(M, \xi)$ be a contact manifold. A submanifold $Y \subset M$ is an almost contact submanifold if there exists a homotopy $\left(\eta_{t}, \omega_{t}\right), t \in[0,1]$, where $\left.\eta_{t} \subset T M\right|_{Y}$ is a codimension 1 distribution of $T M$ along $Y$ and $\omega_{t}$ is a conformal symplectic structure on $\eta_{t}$, such that:
(1) $\eta_{0}=\left.\xi\right|_{Y}$ and $\omega_{0}$ is induced from $\left.\xi\right|_{Y}$; and
(2) $T Y \pitchfork \eta_{1}$ and the normal bundle $T_{Y} M \subset \eta_{1}$ is $\omega_{1}$-symplectic.

A straightforward calculation (see [BCS14, Lemma 2.17]) shows that any evencodimensional submanifold with trivial normal bundle is almost contact.

Corollary 1.3.6. Any almost contact submanifold can be $C^{0}$-approximated by a genuine contact submanifold.
Corollary 1.3.7. Any (coorientable) contact submanifold can be $C^{0}$-approximated by another contact submanifold with the opposite orientation.

Remark 1.3.8. Around the same time as our paper, Casals-Pancholi-Presas [CPP] proved the existence of iso-contact embeddings in codimension 2. Their work and ours are equivalent via the work of Pancholi-Pandit [PP] on iso-contact embeddings.

Wrapping up the introduction, by combining Corollary 1.3.6 with Gromov's $h$-principle for contact structures on open manifolds, one can deduce the following one-half of the groundbreaking(!) theorem of Borman-Eliashberg-Murphy [BEM15] without too much difficulty:

Corollary 1.3.9 (Borman-Eliashberg-Murphy). The existence h-principle holds for contact structures on closed odd-dimensional manifolds of dimension $4 k+1$, $k \in \mathbb{Z}^{+}$, i.e., there exists a genuine contact structure in the homotopy class of any almost contact structure on $M^{4 k+1}$.

Our proof of Corollary 1.3.9 follows from producing genuine contact structures representing each homotopy class of almost contact structures on $S^{2 n-1}$, i.e., proving the existence $h$-principle for $S^{2 n-1}$. This is easy for $S^{4 k+1}$ (via connected sums of Brieskorn manifolds) and is much more involved for $S^{4 k+3}$.

The proofs of Corollaries 1.3.6, 1.3.7, and 1.3 .9 will be given in Section 11. We note that the proof of [CPP] relies on [BEM15] and cannot be used to prove it. Finally, note that in contrast to the contact structures constructed in [BEM15], the contact submanifolds constructed by Corollary 1.3.6 are not a priori overtwisted.

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## 2. A CONVEXITY CRITERION

Let $\Sigma \subset\left(M^{2 n+1}, \xi\right)$ be a closed cooriented hypersurface. The goal of this section is to give a sufficient condition for the characteristic foliation $\Sigma_{\xi}$ on $\Sigma$ (see Definition 2.1.1) which guarantees the Weinstein convexity of $\Sigma$.
2.1. Characteristic foliations. Let $\alpha$ be a contact form for $\xi$. Let $(-\epsilon, \epsilon) \times \Sigma$ be a tubular neighborhood of $\Sigma=\{0\} \times \Sigma \subset M$. Fix an orientation on $\Sigma$ such that the induced orientation on $(-\epsilon, \epsilon) \times \Sigma$ agrees with the orientation determined by $\alpha \wedge(d \alpha)^{n}$. We now introduce the characteristic foliation $\Sigma_{\xi}$ on $\Sigma$.

Definition 2.1.1. The characteristic foliation $\Sigma_{\xi}$ is an oriented singular line field on $\Sigma$ defined by

$$
\Sigma_{\xi}=\left.\operatorname{ker} d \beta\right|_{\operatorname{ker} \beta}
$$

where $\beta:=\left.\alpha\right|_{\Sigma} \in \Omega^{1}(\Sigma)$. The orientation of $\Sigma_{\xi}$ is determined by requiring that the decomposition $T \Sigma=\Sigma_{\xi} \oplus \Sigma_{\xi}^{\perp}$ respect orientations, where the orthogonal complement $\Sigma_{\xi}^{\perp}$, taken with respect to an auxiliary Riemannian metric on $\Sigma$, is oriented by $\left.\beta \wedge(d \beta)^{n-1}\right|_{\Sigma_{\xi}}$.

Remark 2.1.2. The characteristic foliation depends only on the contact structure and the orientation of $\Sigma$, and not on the choice of the contact form.

Note that $x \in \Sigma$ is a singular point of $\Sigma_{\xi}$ if $T_{x} \Sigma=\xi_{x}$ as unoriented spaces. We say $x$ is positive (resp. negative) if $T_{x} \Sigma= \pm \xi_{x}$ as oriented spaces, respectively.

The significance of the characteristic foliation in 3-dimensional contact topology is that it uniquely determines the germ of contact structures on any embedded surface. The corresponding statement for hypersurfaces in contact manifolds of dimension $>3$ is unlikely to hold, i.e., the characteristic foliation by itself is not enough to determine the contact germ. Instead we have the following characterization of contact germs on hypersurfaces in any dimension. The proof is a standard application of the Moser technique and is omitted here.

Lemma 2.1.3. Suppose $\xi_{i}=\operatorname{ker} \alpha_{i}, i=0,1$, are contact structures on $M$ such that $\beta_{0}=g \beta_{1} \in \Omega^{1}(\Sigma)$ for some $g: \Sigma \rightarrow \mathbb{R}_{+}$, where $\beta_{i}=\left.\alpha_{i}\right|_{\Sigma}$. Then there exists an isotopy $\phi_{s}: M \xrightarrow{\sim} M, s \in[0,1]$, such that $\phi_{0}=\operatorname{id}_{M}, \phi_{s}(\Sigma)=\Sigma$ and $\left(\phi_{1}\right)_{*}\left(\xi_{0}\right)=\xi_{1}$ on a neighborhood of $\Sigma$.

Generally speaking, $\Sigma_{\xi}$ can be rather complicated, even when $\Sigma$ is convex with Liouville $R_{ \pm}(\Sigma)$. For our purposes of this paper, it is more convenient to regard $\Sigma_{\xi}$ as a vector field rather than an oriented line field. Of course there is no natural way to specify the magnitude of $\Sigma_{\xi}$ as a vector field, which motivates the following definition: Two vector fields $v_{1}, v_{2}$ on $\Sigma$ are conformally equivalent if there exists a positive function $h: \Sigma \rightarrow \mathbb{R}_{+}$such that $v_{1}=h v_{2}$. This is clearly an equivalence relation among all vector fields, and we will not distinguish conformally equivalent vector fields in the rest of the paper unless otherwise stated.

In order to state the convexity criterion, we need to prepare some generalities on gradient-like vector fields in the following subsection. Our treatment on this subject will be kept to a minimum. The reader is referred to the classical works of Cerf [Cer70] and Hatcher-Wagoner [HW73] for more thorough discussions. Indeed the adaptation of the techniques of Cerf and Hatcher-Wagoner to CHT is carried out in [BHH]. Note that similar techniques in symplectic topology have been developed by Cieliebak-Eliashberg in [CE12].
2.2. Morse and 1-Morse vector fields. Let $Y$ be a closed manifold of dimension $n$. A smooth function $f: Y \rightarrow \mathbb{R}$ is Morse if all the critical points of $f$ (i.e., points $p \in Y$ such that $d f(p)=0$ ) are nondegenerate, i.e., there exist local coordinates $x_{1}, \ldots, x_{n}$ about $p$ such that locally $f$ takes the form

$$
\begin{equation*}
-x_{1}^{2}-\cdots-x_{k}^{2}+x_{k+1}^{2}+\cdots+x_{n}^{2} \tag{2.2.1}
\end{equation*}
$$

Here $k$ is called the Morse index, or just the index, of the critical point $p$. We write $\operatorname{ind}(p)=k$.

A smooth function $f: Y \rightarrow \mathbb{R}$ is 1-Morse if the critical points of $f$ are either nondegenerate or of birth-death type. Here a critical point $p \in Y$ of $f$ is of birthdeath type (also called embryonic) if there exist local coordinates $x_{1}, \ldots, x_{n}$ about $p$ such that $f$ takes the form

$$
\begin{equation*}
-x_{1}^{2}-\cdots-x_{k}^{2}+x_{k+1}^{2}+\cdots+x_{n-1}^{2}+x_{n}^{3} . \tag{2.2.2}
\end{equation*}
$$

Similarly, $k$ is defined to be the (Morse) index of $p$. The birth-death type critical point fits into a 1-parameter family of 1-Morse functions

$$
-x_{1}^{2}-\cdots-x_{k}^{2}+x_{k+1}^{2}+\cdots+x_{n-1}^{2}+t x_{n}+x_{n}^{3},
$$

such that for $t<0$, there exist two nondegenerate critical points of indices $k$ and $k+1$; for $t=0$, there exists a birth-death type critical point; and for $t>0$, there are no critical points.

It is a well-known fact due to Morse that any smooth function can be $C^{\infty}$ _ approximated by a Morse function. Moreover, Cerf proved that any 1-parameter family of smooth functions can be $C^{\infty}$-approximated by a family of 1-Morse functions, where the birth-death type critical points as above occur only at isolated moments.

Given a 1-Morse function $f: Y \rightarrow \mathbb{R}$, we say a vector field $v$ on $Y$ is gradientlike for $f$ if the following two conditions are satisfied:
(GL1) Near each critical point of $f, v=\nabla f$ with respect to some Riemannian metric; and
(GL2) $f$ is strictly increasing along (non-constant) flow lines of $v$.

## Definition 2.2.1.

(1) A vector field $v$ on $Y$ is Morse (resp. 1-Morse) if there exists a Morse (resp. 1-Morse) function $f: Y \rightarrow \mathbb{R}$ such that $v$ is gradient-like for $f$.
(2) A 1-parameter family of vector fields $\left(v_{t}\right)_{t \in[0,1]}$ is a 1-Morse family if each $v_{t}$ is 1 -Morse and there exist $t_{1}<\cdots<t_{k} \in(0,1)$ such that the birthdeath type singularities occur only at $t_{i}$ and there is a single birth-death type singularity at each $t_{i}$.
We also make the slightly nonstandard definition:
Definition 2.2.2. A Liouville domain is 1-Weinstein if its Liouville vector field is gradient-like with respect to a 1-Morse function.
Remark 2.2.3. 1-Morse functions will be sufficient for the purposes of this paper since we will only encounter 1-parameter families of functions. In [BHH], we will need to deal with generic 2-parameter families of functions (called 2-Morse functions) where new singularities, i.e., the swallowtails, appear.
Convention 2.2.4. A flow line $\ell:(a, b) \rightarrow Y$ of a vector field $v$ on $Y$ is assumed to be a maximal oriented smooth trajectory $\mathbb{R} \rightarrow Y$ that has been precomposed with an orientation-preserving reparametrization $(a, b) \xrightarrow{\sim} \mathbb{R}$. A partial flow line is the restriction of a flow line to a subinterval.

Definition 2.2.5. A broken flow line (resp. possibly broken flow line) of a vector field $v$ on $Y$ is a continuous map $\ell:[a, b] \rightarrow Y$ such that there exists an increasing sequence $a=a_{0}<a_{1}<\cdots<a_{m}=b$ with $m>1(r e s p . m \geq 1)$ such that $\ell\left(a_{j}\right)$, $j=0, \ldots, m$, are zeros of $v$ and $\left.\ell\right|_{\left(a_{j}, a_{j+1}\right)}, j=0, \ldots, m-1$, are flow lines of $v$. We may also replace $[a, b]$ by half-open or open intervals.

In the rest of this subsection, we present a simple criterion for a vector field to be Morse which will be useful for our later applications. The corresponding version for 1-Morse vector fields is left to the reader as an exercise.

Proposition 2.2.6. A vector field $v$ on a closed manifold $Y$ is Morse (resp. 1Morse) if and only if the following conditions are satisfied:
(M1) For any point $x \in Y$ with $v(x)=0$, there exists a neighborhood of $x$ and a locally defined function $f$ of the form given by Eq. (2.2.1) (resp. Eq. (2.2.1) or Eq. (2.2.2)) such that $v=\nabla f$.
(M2) For any point $x \in Y$ with $v(x) \neq 0$, the flow line of $v$ passing through $x$ converges to zeros of $v$ in both forward and backward time.
(M3) There exist no possibly broken loops, i.e., a possibly broken flow line $\ell$ : $[0,1] \rightarrow Y$ such that $\ell(0)=\ell(1)$.
Proof. The "only if" direction is obvious. To prove the "if" direction, let $Z(v)=$ $\left\{x_{1}, \ldots, x_{k}\right\}$ be the finite set of zeros of $v$, where the finiteness is guaranteed by (M1) and the compactness of $Y$. Then we define a partial order on $Z(v)$ such that $x_{i} \prec x_{j}$ if there exists a flow line of $v$ from $x_{i}$ to $x_{j}$. The fact that $\prec$ is a partial order follows from (M3).

We then construct a handle decomposition of $Y$ starting from the minimal elements $Z_{0}$ of $Z(v)$ (note that a minimal element of $Z(v)$ has index 0 by (M2)) and inductively attaching handles as follows: Starting with a standard neighborhood of $Z_{0}$, suppose we have already attached the handles corresponding to $Z_{j}$. Then we attach the handles corresponding to the minimal elements of $Z(v)-Z_{j}$, and then let $Z_{j+1}$ be the union of $Z_{j}$ and the minimal elements of $Z(v)-Z_{j}$.
2.3. A convexity criterion. The goal of this subsection is to give a sufficient condition for a hypersurface to be Weinstein convex. To this end, we introduce the notions of Morse and Morse ${ }^{+}$hypersurfaces whose characteristic foliations have particularly simple dynamics.

## Definition 2.3.1.

(1) A hypersurface $\Sigma \subset(M, \xi)$ is Morse (resp. 1-Morse) if there exists a representative $v$ in the conformal equivalence class of $\Sigma_{\xi}$ which is a Morse (resp. 1-Morse) vector field on $\Sigma$. We say $\Sigma$ is Morse $^{+}$(resp. 1-Morse ${ }^{+}$) if, in addition, there exist no flow trajectories from a negative singular point of $v$ to a positive one.
(2) A 1-parameter family of hypersurfaces $\left(\Sigma_{t}\right)_{t \in[0,1]}$ is a 1-Morse family, if $\left(\left(\Sigma_{t}\right)_{\xi}\right)_{t \in[0,1]}$ is represented by a 1-Morse family.

Lemma 2.3.2. If $\Sigma$ is a Morse hypersurface, then a $C^{\infty}$-small perturbation of $\Sigma$ is Morse ${ }^{+}$.

Proof. Choose a contact form $\xi=\operatorname{ker} \alpha$. It suffices to observe that $\left.d \alpha\right|_{\Sigma}$ is nondegenerate on a neighborhood of the singular points of $\Sigma_{\xi}$. It is a standard fact (see e.g. [CE12, Proposition 11.9]) that the Morse index $\operatorname{ind}(x) \leq n$ if $x$ is a positive singular point of $\Sigma_{\xi}$, and $\operatorname{ind}(x) \geq n$ if $x$ is negative. The claim therefore follows from the usual transversality argument.

The following proposition gives a sufficient condition for convexity:

## Proposition 2.3.3.

(1) A 1-Morse ${ }^{+}$hypersurface $\Sigma$ is convex.
(2) A hypersurface $\Sigma$ is Weinstein convex if and only if it is Morse ${ }^{+}$.

Proof. (1) is a straightforward generalization of the usual proof for surfaces due to Giroux that $\Sigma$ is convex if it has a Morse ${ }^{+}$characteristic foliation.

Let $\mathbf{x}=\left\{x_{1}, \ldots, x_{m}\right\}$ (resp. $\mathbf{y}=\left\{y_{1}, \ldots, y_{\ell}\right\}$ ) be the positive (resp. negative) singular points of $\Sigma_{\xi}$. Then $d \beta$ is nondegenerate on a small open neighborhood $U(\mathbf{x})$ of $\mathbf{x}$, where $\beta:=\left.\alpha\right|_{\Sigma}$. Let $W_{x_{i}}$ be the stable manifold of $x_{i}$ with respect to the gradient of the 1 -Morse function and let the $i$ th skeleton $\mathrm{Sk}_{i}$ be the closure of $W_{x_{1}} \cup \cdots \cup W_{x_{i}}$. We order the points of $\mathbf{x}$ so that $W_{x_{i+1}}$ intersects the boundary of a small open neighborhood of $\mathrm{Sk}_{i}$ along a sphere if $x_{i+1}$ is a Morse critical point and along a disk if $x_{i+1}$ is an embryonic point. In particular we necessarily have $\operatorname{ind}\left(x_{1}\right)=0$, but we do not require $\operatorname{ind}\left(x_{i}\right) \geq \operatorname{ind}\left(x_{j}\right)$ for $i>j$. Such an arrangement is possible thanks to the assumption that there is no flow line of $\Sigma_{\xi}$ going from $\mathbf{y}$ to $\mathbf{x}$.

There exists a conformal modification $\beta \rightsquigarrow g \beta$, where $g$ is a positive function, so that it becomes Liouville on a tubular neighborhood $U\left(\mathrm{Sk}_{m}\right)$ of $\mathrm{Sk}_{m}$ and $\partial U\left(\mathrm{Sk}_{m}\right)$ is contact: Arguing by induction, suppose that $\beta$ is Liouville on $U\left(\mathrm{Sk}_{i}\right)$ such that $\partial U\left(\mathrm{Sk}_{i}\right)$ is contact. During the induction we will often reset notation, i.e., modify $\beta \rightsquigarrow g \beta$ and call the result the new $\beta$. We will explain the case where $x_{i+1}$ is Morse and $W_{x_{i+1}} \cap \partial U\left(\mathrm{Sk}_{i}\right)$ is a Legendrian sphere $\Lambda \subset \partial U\left(\mathrm{Sk}_{i}\right)$; the other cases are similar. Using the flow of $\Sigma_{\xi}$, we may identify a tubular neighborhood of $W_{x_{i+1}} \backslash\left(U\left(\mathrm{Sk}_{i}\right) \cup U\left(x_{i+1}\right)\right)$ with $[0,1]_{r} \times Y$, where $Y$ is an open neighborhood of the 0 -section in $J^{1}(\Lambda)$ such that:

- $\{0\} \times Y \subset \partial U\left(\mathrm{Sk}_{i}\right)$;
- $\{1\} \times Y \subset \partial U\left(x_{i+1}\right)$; and
- $\partial_{r}$ is identified with $\Sigma_{\xi}$.

It follows that one can write $\beta=g \lambda$ on $[0,1] \times Y$, where $\lambda$ is a contact form on $Y$ and $g$ is a positive function on $[0,1] \times Y$. Note that

$$
d \beta=\partial_{r} g d r \wedge \lambda+d_{Y} g \wedge \lambda+g d \lambda
$$

is symplectic if $\partial_{r} g>0$. By assumption we have $\partial_{r} g>0$ when $r$ is close to 0 or 1 . Rescaling $\left.\beta\right|_{U\left(x_{i+1}\right)}$ by a large constant $K \gg 0$, we can extend $\left.\beta\right|_{U\left(\mathrm{Sk}_{i}\right) \cup U\left(x_{i+1}\right)}$ to a Liouville form on $U\left(\mathrm{Sk}_{i+1}\right)$. Moreover, we can assume $\partial U\left(\mathrm{Sk}_{i+1}\right)$ is transverse to $\Sigma_{\xi}$ by slightly shrinking $U\left(\mathrm{Sk}_{i+1}\right)$. Hence by induction we can arrange so that $\beta$ is a Liouville form on $U\left(\mathrm{Sk}_{m}\right)$.

The treatment of the negative singular points of $\Sigma_{\xi}$ is similar. Let $\mathrm{Sk}_{\ell}^{\prime}$ be the closure of the union of the unstable manifolds of $\mathbf{y}$. Then by the same argument we can assume that $\beta$ is a Liouville form on $-U\left(\mathrm{Sk}_{\ell}^{\prime}\right)$, where the minus sign indicates the opposite orientation.

Using the flow of $\Sigma_{\xi}$, we can identify $\Sigma \backslash\left(U\left(\mathrm{Sk}_{m}\right) \cup U\left(\mathrm{Sk}_{\ell}^{\prime}\right)\right)$ with $\Gamma \times[-1,1]_{s}$ such that:

- $\Gamma \times\{-1\}$ is identified with $\partial U\left(\mathrm{Sk}_{m}\right)$;
- $\Gamma \times\{1\}$ is identified with $\partial U\left(\mathrm{Sk}_{\ell}^{\prime}\right)$; and
- $\mathbb{R}\left\langle\partial_{s}\right\rangle=\Sigma_{\xi}$.

We can write $\beta=h \eta$ near $\Gamma \times\{-1,1\}$, where $\eta$ is a contact form on $\Gamma$ and $h=h(s)$ is a positive function such that $h^{\prime}(s)>0$ near $\Gamma \times\{-1\}$ and $h^{\prime}(s)<0$ near $\Gamma \times\{1\}$. Extend $h$ to a positive function $\Gamma \times[-1,1] \rightarrow \mathbb{R}$ such that $h^{\prime}(s)>0$ for $s<0, h^{\prime}(0)=0$, and $h^{\prime}(s)<0$ for $s>0$. Let $f=f(s): \Gamma \times[-1,1] \rightarrow \mathbb{R}$ be a nonincreasing function of $s$ such that $f(-1)=1, f(0)=0, f(1)=-1$, $f^{\prime}(0)<0$, and $f^{(n)}(-1)=0=f^{(n)}(1)$ for all $n \geq 1$. Then define $\rho=f d t+h \eta$ on $\mathbb{R}_{t} \times \Gamma \times[-1,1], \rho=d t+\beta$ on $\mathbb{R} \times U\left(\mathrm{Sk}_{m}\right)$, and $\rho=-d t+\beta$ on $\mathbb{R} \times U\left(\mathrm{Sk}_{\mathbf{y}_{\ell}}^{\prime}\right)$. We leave it to the reader to check that $\rho$ is contact and that $\left.\rho\right|_{\{0\} \times \Sigma}$ agrees with $\left.\alpha\right|_{\Sigma}$ up to an overall positive function. (1) now follows from Lemma 2.1.3.
(2) The "if" direction follows from the proof of (1) and the "only if" direction is clear.

## 3. Construction of mushrooms in dimension 3

In order to make a hypersurface $\Sigma \subset(M, \xi)$ Weinstein convex, we would like to modify the characteristic foliation $\Sigma_{\xi}$ so it is directed by a Morse vector field and then apply Proposition 2.3.3. (Note that going from Morse to Morse ${ }^{+}$is a $C^{\infty}$-generic operation which can always be done by Lemma 2.3.2.) This will be achieved by certain $C^{0}$-small perturbations of $\Sigma$ which we call mushrooms. The mushrooms are most easily described in dimension 3 and the general case will be constructed in Section 5 using 3-dimensional mushrooms. It turns out that mushrooms alone are enough to make any $\Sigma_{\xi}$ Morse if $\operatorname{dim} \Sigma=2$. If $\operatorname{dim} \Sigma>2$, then mushrooms are not quite sufficient and we will need an additional technical construction in Section 7.

The standard model of a mushroom will be constructed in a Darboux chart

$$
\left(\mathbb{R}_{z, s, t}^{3}, \xi=\operatorname{ker} \alpha\right), \quad \alpha=d z+e^{s} d t
$$

Let $\Sigma=\{z=0\}$ be the surface under consideration with normal orientation $\partial_{z}$ and characteristic foliation $\Sigma_{\xi}$ directed by $\partial_{s}$. The goal of this section is to "fold" $\Sigma$ to obtain another surface $Z$ which coincides with $\Sigma$ outside of a compact set, and analyze the change in the dynamics of the characteristic foliations.

In $\S 3.1$ we construct a piecewise linear (PL) model $Z_{\mathrm{PL}}$ and then in $\S 3.2$ we round the corners of $Z_{\mathrm{PL}}$ to obtain a suitably generic smooth surface $Z$ such that the characteristic foliation $Z_{\xi}$ has the desired properties.

Remark 3.0.1. In an earlier version of the paper we constructed a mushroom whose base was smaller than the cap and discussed "mushroom packing ratios". In the current version they approximately have the same size and can be packed tightly.
3.1. Construction of $Z_{\mathrm{PL}}$. Choose a rectangle $\square=\left[0, s_{0}\right] \times\left[0, t_{0}\right] \subset \Sigma$, where $s_{0}, t_{0}>0$. We define $Z_{\mathrm{PL}}$ to coincide with $\Sigma$ outside of $\square$.

Remark 3.1.1. The more general case $\left[s_{-1}, s_{0}\right] \times\left[t_{-1}, t_{0}\right]$ can be computed analogously. In what follows we can replace $1-e^{-s_{0}}$ by $e^{-s_{-1}}-e^{-s_{0}}$ and the $s$-width $\mathscr{S}\left(Z_{P L}\right)$ becomes $s_{0}-s_{-1}$.

Choose $z_{0}>0$. We construct three rectangles $P_{0}, P_{2}, P_{4}$ and two parallelograms $P_{1}, P_{3}$ in $\mathbb{R}^{3}$, which, together with $\mathbb{R}_{s, t}^{2} \backslash \square$, glue to give $Z_{\mathrm{PL}}$, i.e., we define

- $P_{0}:=\left[0, s_{0}\right] \times\left[-e^{-s_{0} / 2} z_{0},-e^{-s_{0} / 2} z_{0}+t_{0}\right] \subset\left\{z=z_{0}\right\} ;$
- $P_{i}, i=1, \ldots, 4$, are the faces $\left(\neq P_{0}, \square\right)$ of the convex hull of $P_{0} \cup \square$ (which is a parallelepiped), ordered counterclockwise so that $P_{1} \subset\{s=$ $0\}$.

Definition 3.1.2. We define the PL surface

$$
Z_{\mathrm{PL}}:=(\Sigma \backslash \square) \cup\left(\cup_{0 \leq i \leq 4} P_{i}\right)
$$

The rectangle $\square \subset \Sigma$ (resp. $P_{0}$ ) is called the base (resp. cap) of the PL mushroom and $\cup_{0 \leq i \leq 4} P_{i}$ is called the PL mushroom. We refer to the modification $\Sigma \rightsquigarrow Z_{\mathrm{PL}}$ as "growing a PL mushroom".


Figure 3.1.1. The PL model $Z_{\mathrm{PL}}$. Here $P_{0}$ is the top face; $P_{1}$ and $P_{3}$ are the front and back faces, respectively; $P_{2}$ and $P_{4}$ are the right and left faces, respectively; and $\square$ is the bottom face which is not part of $Z_{\mathrm{PL}}$.

Observe that, away from the corners, the characteristic foliation $\left(Z_{\mathrm{PL}}\right)_{\xi}$ on $Z_{\mathrm{PL}}$ satisfies

- $\left(Z_{\mathrm{PL}}\right)_{\xi}=\mathbb{R}\left\langle\partial_{s}\right\rangle$ on $\Sigma \backslash \square$ and $P_{0}$;
- on $P_{2}$ (resp. $\left.P_{4}\right),\left(Z_{\mathrm{PL}}\right)_{\xi}$ is directed by $\partial_{s}$ (resp. $-\partial_{s}$ ) for $s \in\left[0, \frac{s_{0}}{2}\right.$ ), is singular along $s=\frac{s_{0}}{2}$, and is directed by $-\partial_{s}$ (resp. $\partial_{s}$ ) for $s \in\left(\frac{s_{0}}{2}, s_{0}\right]$;
- $\left(Z_{\mathrm{PL}}\right)_{\xi}$ is the linear foliation on $P_{1}$ and $P_{3}$ with "slopes" -1 and $-e^{-s_{0}}$, respectively, where "slope" refers to the value of $d t / d z=-e^{-s}$. See Figure 3.1.2.
We refer to the singular line segments on $P_{2}$ and $P_{4}$ by $\mathscr{S}_{-}$and $\mathscr{S}_{+}$indicating their signs.


Figure 3.1.2. The linear characteristic foliations on $P_{1}$ (left) and $P_{3}$ (right) where $Z_{\mathrm{PL}}$ is sufficiently thin.

We now analyze the dynamics of the PL flow on $Z_{\mathrm{PL}}$. Note that the flow lines are not necessarily uniquely determined by the initial conditions due to the presence of corners.

We begin by introducing a few quantities which characterize the various sizes of the mushrooms.

Definition 3.1.3. Given $Z_{\mathrm{PL}}$ as above, its $z$-height, $s$-width, and $t$-width are given by:

$$
\mathscr{Z}\left(Z_{\mathrm{PL}}\right):=z_{0}, \quad \mathscr{S}\left(Z_{\mathrm{PL}}\right):=s_{0}, \quad \mathscr{T}\left(Z_{\mathrm{PL}}\right):=t_{0} .
$$

The following lemma characterizes a key feature of $\left(Z_{\mathrm{PL}}\right)_{\xi}$ when the parameters of the mushrooms are appropriately adjusted.
Lemma 3.1.4. Fix $s_{0}, z_{0}>0$. If $t_{0}<\left(1-e^{-s_{0}}\right) z_{0}$, then
(1) the unique flow line of $\left(Z_{\mathrm{PL}}\right)_{\xi}$ passing through $(-1, a) \in \mathbb{R}_{s, t}^{2}$, where $a \in\left(0, t_{0}\right)$, either hits $P_{0} \cap P_{2}$ or converges to $\mathscr{S}_{-}$in forward time;
(2) the unique flow line of $\left(Z_{\mathrm{PL}}\right)_{\xi}$ passing through $\left(s_{0}+1, a\right), a \in\left(0, t_{0}\right)$, either hits $P_{0} \cap P_{4}$ or converges to $\mathscr{S}_{+}$in backward time; and
(3) all the flow lines of $\left(Z_{\mathrm{PL}}\right)_{\xi}$ passing through $(-1, a), a \notin\left[0, t_{0}\right]$, or $\left(s_{0}+\right.$ $1, a), a \notin\left[0, t_{0}\right]$, are unaffected.
Proof. (1) Since $t_{0}-z_{0}\left(1-e^{-s_{0}}\right)<0$, it follows that the unique flow line passing through the point $(-1, a) \in \Sigma$ with $a \in\left(0, t_{0}\right)$ does one of three things in forward time:

- travels over $P_{0}$, enters $P_{1}$, and ends at $\mathscr{S}_{-}$;
- travels over $P_{0}$ into $P_{0} \cap P_{2}$ (this happens with only one flow line); or
- travels over $P_{0}, P_{1}, P_{2}$ in that order, enters $P_{1}$, and ends at $\mathscr{S}_{-}$.

See Figure 3.1.3. (2) is similar and (3) is clear.


Figure 3.1.3. Flow lines of $\left(Z_{\mathrm{PL}}\right)_{\xi}$ limiting to $\mathscr{S}_{-}$in red (the generic case) and one flow line limiting to $P_{0} \cap P_{2}$.
3.2. Smoothing of $Z_{\mathrm{PL}}$. In this subsection we construct a smoothing of $Z_{\mathrm{PL}}$.

Construction of the smoothing $Z$. Choose a small smoothing parameter $\delta>0$ and a smooth "profile function" $\phi:\left[0, z_{0}\right] \rightarrow[-\delta, \delta]$ such that $\phi(0)=\delta, \phi\left(z_{0}\right)=-\delta$ and has "derivative $-\infty$ " at $z=0, z_{0}$.

For each $z^{\prime} \in\left(0, z_{0}\right)$, the slices $R_{z^{\prime}}:=Z_{\mathrm{PL}} \cap\left\{z=z^{\prime}\right\}$ are rectangles. When $z^{\prime}=0$ or $z_{0}$, we take $R_{z^{\prime}}=\partial\left(Z_{\mathrm{PL}} \cap\left\{z=z^{\prime}\right\}\right)$. For each $\delta^{\prime} \in \mathbb{R}$ with $\left|\delta^{\prime}\right|$ small, let $R_{z^{\prime}}^{\delta^{\prime}} \subset\left\{z=z^{\prime}\right\}$ be the rectangle concentric to $R_{z^{\prime}}$ and whose side lengths are $\delta^{\prime}$ larger. Then let $\widetilde{R}_{z^{\prime}}^{\delta^{\prime}} \subset\left\{z=z^{\prime}\right\}$ be a smoothing of $R_{z^{\prime}}^{\delta^{\prime}}$ contained in the closure of the region between $R_{z^{\prime}}^{\delta^{\prime}}$ and $R_{z^{\prime}}^{\delta^{\prime}-\delta}$, which:
(R1) agrees with $R_{z^{\prime}}^{\delta^{\prime}}$ on the complement of the open $\delta$-neighborhoods of the edges of $R_{z^{\prime}}^{\delta^{\prime}}$ parallel to $t=$ const;
(R2) has nonzero curvature on these $\delta$-neighborhoods and is tangent to $R_{z^{\prime}}^{\delta^{\prime}}$ precisely at the midpoints of the edges of $R_{z^{\prime}}^{\delta^{\prime}}$ (i.e., when $s=s_{0} / 2$ ); and
(R3) is smoothly varying with $z^{\prime}$ and $\delta^{\prime}$.
The $\phi$-smoothing $Z$ of $Z_{\mathrm{PL}}$ with smoothing parameter $\delta>0$ and profile function $\phi$ is obtained by modifying $Z_{\mathrm{PL}}$ as follows:
(1) replace $R_{z^{\prime}}$ by $\widetilde{R}_{z^{\prime}}^{\phi\left(z^{\prime}\right)}$ for $z \in\left(0, z_{0}\right)$;
(2) remove the bounded component of $\{z=0\} \backslash \widetilde{R}_{0}^{\phi(0)}$; and
(3) adjoin the closure of the bounded component of $\left\{z=z_{0}\right\} \backslash \widetilde{R}_{z_{0}}^{\phi\left(z_{0}\right)}$.

The base of the mushroom $\tilde{\square} \subset \Sigma$ of $Z$ is the closure of $\Sigma \backslash Z$ and the mushroom is the closure of $Z \backslash \Sigma$. By construction $\widetilde{\square}$ converges to $\square$ when all the parameters tend to zero.

The following proposition describes the key dynamical properties of $Z_{\xi}$.
Proposition 3.2.1. Given $Z_{P L}$ with parameters $s_{0}, z_{0}, t_{0}, \epsilon$ satisfying $t_{0}<(1-$ $\left.e^{-s_{0}}\right) z_{0}$, there exists a smoothing $Z$ of $Z_{P L}$ with small smoothing parameter $\delta>$ 0 and profile function $\phi:\left[0, z_{0}\right] \rightarrow[-\delta, \delta]$ whose vector field $Z_{\xi}$ satisfies the following properties:
(TZO) $Z_{\xi}$ is gradient-like with respect to a Morse function $f_{a}: \widetilde{Z} \rightarrow \mathbb{R}$ which agrees with s outside of $\widetilde{\square}$.
(TZ1) $Z_{\xi}$ has four nondegenerate singularities: a positive source $e_{+}$and a positive saddle $h_{+}$near the midpoint of $\mathbb{R}_{s, t}^{2} \cap P_{4}$, and a negative sink $e_{-}$and a negative saddle $h_{-}$near the midpoint of $\mathbb{R}_{s, t}^{2} \cap P_{2}$.
(TZ2) There is a unique flow line each from $e_{+}$to $h_{+}$, from $e_{+}$to $h_{-}$, from $h_{+}$ to $e_{-}$, and from $h_{-}$to $e_{-}$, and the four flow lines bound a quadrilateral whose interior consists of flow lines from $e_{+}$to $e_{-}$.
(TZ3) There exist $\kappa_{1}>\kappa_{2}>\kappa_{3}>0>\kappa_{4}>\kappa_{5}>\kappa_{6}$ such that all $\kappa_{i} \rightarrow 0$ as $\delta \rightarrow 0$ and the following hold:
(1) the stable manifold of $h_{+}$(resp. $h_{-}$) intersects the line $\{s=-1\}$ at $\left(-1, t_{0}+\kappa_{2}\right)\left(\right.$ resp. $\left.\left(-1, \kappa_{4}\right)\right)$,
(2) the unstable manifold of $h_{+}$(resp. $h_{-}$) intersects the line $\left\{s=s_{0}+1\right\}$ at $\left(s_{0}+1, t_{0}+\kappa_{3}\right)\left(\right.$ resp. $\left.\left(s_{0}+1, \kappa_{5}\right)\right)$,
(3) any flow line passing through $(-1, a), a \in\left(\kappa_{4}, t_{0}+\kappa_{2}\right)$, converges to $e_{-}$in forward time,
(4) any flow line passing through $\left(s_{0}+1, a\right), a \in\left(\kappa_{5}, t_{0}+\kappa_{3}\right)$, converges to $e_{+}$in backward time,
(5) a flow line passes through $(-1, a), a \notin\left[\kappa_{6}, t_{0}+\kappa_{1}\right]$, if and only if it passes through $\left(s_{0}+1, a\right), a \notin\left[\kappa_{6}, t_{0}+\kappa_{1}\right]$,
(6) a flow line passes through $(-1, a), a \in\left(t_{0}+\kappa_{2}, t_{0}+\kappa_{1}\right)$, if and only if it passes through $\left(s_{0}+1, a^{\prime}\right), a^{\prime} \in\left(t_{0}+\kappa_{3}, t_{0}+\kappa_{1}\right)$,
(7) a flow line passes through $(-1, a), a \in\left(\kappa_{6}, \kappa_{4}\right)$ if and only if it passes through $\left(s_{0}+1, a^{\prime}\right), a^{\prime} \in\left(\kappa_{6}, \kappa_{5}\right)$.
(TZ4) The flow lines described in (TZ2) and (TZ3) are all the flow lines that nontrivially intersect the mushroom.

In words, $Z_{\xi}$ blocks all flow lines that pass through an open interval that is close to $\{-1\} \times\left(0, t_{0}\right)$, slightly bends unblocked flow lines that pass through points close to $(-1,0)$ and $\left(-1, t_{0}\right)$ (with bending $\rightarrow 0$ as the smoothing parameter $\delta \rightarrow 0$ ), and leaves all other flow lines passing through $s=-1$ untouched.

See Figure 3.2.1 for an illustration of the effect of a mushroom on the characteristic foliation and also the values of $t$ where certain flow lines in (TZ3) intersect $s=-1$ or $s=s_{0}+1$.


Figure 3.2.1. The characteristic foliations before and after a mushroom. The numbers $t_{0}+\kappa_{1}$ etc. are the $t$-coordinates where the indicated flow lines intersect $s=-1$ or $s=s_{0}+1$; see (TZ3).

Proof. Let $Z$ be the smoothing of $Z_{P L}$ from above.
(TZ1) We determine the singular points of $Z_{\xi}$ as follows: The first requirement for $\xi$ to be tangent to $Z$ is for $Z$ to be tangent to $\partial_{s}$. Since $\widetilde{R}_{z^{\prime}}^{\delta^{\prime}}$ is tangent to $\partial_{s}$ exactly at two points by (R2), i.e., when $s=s_{0} / 2$, the singular points lie on the restriction of $Z$ to the slice $s=s_{0} / 2$.

We now choose $\phi$ such that $\phi^{\prime}<0$ on $\left(0, c_{0}\right)$ and $\left(c_{1}, z_{0}\right)$ and $\phi^{\prime}>0$ on $\left(c_{0}, c_{1}\right)$, where $0<c_{0}<c_{1} \ll z_{0}$. By the choice of $\phi$, there are four points where $Z \cap\left\{s=s_{0} / 2\right\}$ is tangent to $\xi \cap\left\{s=s_{0} / 2\right\}$. They occur as described in (TZ1) since $\phi^{\prime}=0$ at $z=c_{0}, c_{1}$, which are both close to $z=0$.
(TZ0), (TZ2)-(TZ5) then follow from Lemma 3.1.4. The conditions $\kappa_{2}>\kappa_{3}$ and $\kappa_{4}>\kappa_{5}$ are required since the surface $Z$ was obtained from $\Sigma$ by pushing in the positive $z$-direction.

The following remark will also be very useful later:
Remark 3.2.2. Proposition 3.2.1 also holds with (TZ1) replaced by:
(TZ1') $Z_{\xi}$ has two singularities: a positive birth-death singularity near the midpoint of $\mathbb{R}_{s, t}^{2} \cap P_{4}$ and a negative birth-death singularity near the midpoint of $\mathbb{R}_{s, t}^{2} \cap P_{2}$.
Moreover, there is a foliated 1-parameter family of surfaces from $\{z=0\}$ to a slight upward translate of $Z$ (i.e., in the $z$-direction) whose characteristic foliations have no singularities except for $Z$. For this $Z$ we take $\phi$ such that $\phi^{\prime}<0$ for $z \neq c_{0}$ and $\phi^{\prime}\left(c_{0}\right)=0$. Note that the singularities vanish if we take $\phi$ such that $\phi^{\prime}<0$ for all $z$.

Convention 3.2.3. In view of Proposition 3.2.1, from now on we assume that all mushrooms satisfy $t_{0}<\left(1-e^{-s_{0}}\right) z_{0}$.

## 4. Convex surface theory revisited

The goal of this section is to give elementary, Morse-theoretic proofs of Theorems 1.2.3 and 1.2.5 in dimension 3 using the folding techniques developed in Section 3. In dimension 3, Theorem 1.2.3 was proved by Giroux in [Gir91] in a stronger form where $C^{0}$ is replaced by $C^{\infty}$. Theorem 1.2.5 can be inferred from Giroux's work on bifurcations [Gir00] and the bypass-bifurcation correspondence. The technical heart of Giroux's work is based on the study of dynamical systems of vector fields on surfaces, a.k.a., Poincaré-Bendixson theory. In particular, one invokes a deep theorem of Peixoto [Pei62] to prove the $C^{\infty}$-version of Theorem 1.2.3 and much more work to establish Theorem 1.2.5.

Our proof strategy is the following: First apply a $C^{\infty}$-small perturbation of $\Sigma \subset\left(M^{3}, \xi\right)$ such that the singularities of $\Sigma_{\xi}$ become Morse. There exists a finite collection of pairwise disjoint transverse arcs $\gamma_{i}, i \in I$, in $\Sigma$ such that any flow line of $\Sigma_{\xi}$ passes through some $\gamma_{i}$. In $\S 4.1$ we will construct a 3 -dimensional plug supported on a small flow box $B_{i}=\left[0, s_{i}\right] \times\left[0, t_{i}\right]$, where $\gamma_{i}=\left\{\frac{s_{i}}{2}\right\} \times\left(\epsilon, t_{i}-\epsilon\right)$ and $\epsilon>0$ is small, such that no flow line that enters through $\{0\} \times\left(\epsilon, t_{i}-\epsilon\right)$ can leave through $\left\{s_{0}\right\} \times\left[0, t_{i}\right]$, i.e., they all necessarily converge to singularities in the plug. Each plug consists of a large number of mushrooms constructed in Section 3. This proves Theorem 1.2.3. To prove Theorem 1.2.5, we slice $\Sigma \times[0,1]$ into thin layers using $\Sigma_{i}:=\Sigma \times\left\{\frac{i}{N}\right\}, 0 \leq i \leq N$, for large $N$ such that the difference between $\left(\Sigma_{i}\right)_{\xi}$ and $\left(\Sigma_{i+1}\right)_{\xi}$ is small. (By "small" we mean the vector fields in question are $C^{0}$-close to each other. The global dynamics of $\left(\Sigma_{i}\right)_{\xi}$ may still drastically differ from that of $\left(\Sigma_{i+1}\right)_{\xi}$.) Within each layer we insert plugs on $\Sigma_{i}$ as in the case of a single surface so that the isotopy from $\Sigma_{i}$ to $\Sigma_{i+1}$ is through a 1-Morse family of surfaces, i.e., $\left(\Sigma_{t}\right)_{\xi}$ is a 1 -Morse family for all $\frac{i}{N} \leq t \leq \frac{i+1}{N}$. For technical reasons, it is desirable to eliminate the plugs created on $\Sigma_{i}$ when we reach $\Sigma_{i+1}$, replacing them by new plugs on $\Sigma_{i+1}$, so that one can inductively run from $i=0$ to $i=N$ and make all intermediate surfaces 1-Morse. Then the only obstructions to convexity occur at finitely many instances where the surface is 1-Morse but not 1 -Morse ${ }^{+}$, corresponding to bypass attachments.

This section is organized as follows: In $\S 4.1$ we describe 3 -dimensional plugs and in $\S 4.4$ we explain how to "install" and "uninstall" plugs. The higher-dimensional plugs will be described in Section 7. We then use this technology to prove Theorem 1.2.3 in §4.3 and Theorem 1.2.5 in §4.5.
4.1. 3-dimensional plugs. The construction of a plug is local. Consider $M=$ $\left[-z_{0}, z_{0}\right] \times\left[0, s_{0}\right] \times\left[0, t_{0}\right]$ with coordinates $(z, s, t)$, contact form $\alpha=d z+e^{s} d t$, and contact structure $\xi=\operatorname{ker} \alpha$. Here $z_{0}, s_{0}, t_{0}>0$ are arbitrary, but for most of our applications, we should think of $z_{0}, s_{0}$ as being much smaller than $t_{0}$. In other words, the condition $t_{0}<\left(1-e^{-s_{0}}\right) z_{0}$ in Proposition 3.2.1 will not be satisfied.

Consider the surface $B=\{0\} \times\left[0, s_{0}\right] \times\left[0, t_{0}\right]$ with $B_{\xi}=\mathbb{R}\left\langle\partial_{s}\right\rangle$. We will refer to $(M, \alpha)$ as a standard contact neighborhood of $B$ with parameters $s_{0}, t_{0}, z_{0}$. Let $\partial_{-} B=\{0\} \times\{0\} \times\left[0, t_{0}\right]$ and $\partial_{+} B=\{0\} \times\left\{s_{0}\right\} \times\left[0, t_{0}\right]$. Pick a large integer
$N$ such that $N \equiv 3 \bmod 4$. Let $\square_{k, l} \subset B$ be boxes defined by

$$
\square_{k, l}:=\left[\frac{2 l-1}{5} s_{0}, \frac{2 l}{5} s_{0}\right] \times\left[\frac{4 k+2 l-1}{N} t_{0}, \frac{4 k+2 l+2}{N} t_{0}\right],
$$

where $0 \leq k<\lfloor N / 4\rfloor, l=1,2$. See Figure 4.1.1. In words, since the mushrooms can be packed "tightly", it suffices to arrange two rows of mushrooms.


Figure 4.1.1. Bases of the mushrooms on $B$. Here $N=15$.
Applying Proposition 3.2.1, we install pairwise disjoint mushrooms $Z_{k, l}$ on $B$ such that the base of each $Z_{k, l}$ approximately equals $\square_{k, l}$. The key property of the dynamics of the resulting $B_{\xi}^{\vee}$ is as follows: Suppose $\epsilon>0$ is small and we choose $N \gg 1 / \epsilon$. Let $x=(z(x), s(x), t(x)) \in \partial_{-} B$.
(1) If $t(x) \in\left(\epsilon, t_{0}-\epsilon\right)$, then the maximal possibly broken flow line of $B_{\xi}^{\vee}$ passing through $x$ converges to a sink in forward time.
(2) If a flow line passes through $x$ and does not converge to sink, then it exits along $y \in \partial_{+} B$ and $|t(y)-t(x)|<\epsilon$.
4.2. Barricades. In this subsection only, let $\Sigma$ be a manifold of dimension $m$ and $v$ a vector field on $\Sigma$ with only Morse singularities. We also fix a Riemannian metric on $\Sigma$.
Definition 4.2.1. $A$ flow box is an embedded cylinder $B=\left[0, s_{0}\right] \times D^{m-1} \subset \Sigma$ with coordinates $(s, x)$ over the disk $D^{m-1}$ such that $\left.v\right|_{B}=\partial_{s}$.

Given $\epsilon>0$ small, let $B^{\epsilon}=\left[\epsilon, s_{0}-\epsilon\right] \times D_{\epsilon}^{m-1}$ be a slightly smaller flow box (called a shrinkage of $B$ ), where $D_{\epsilon}^{m-1} \subset D^{m-1}$ is a disk such that every point $x \in D^{m-1} \backslash D_{\epsilon}^{m-1}$ has metric distance $<\epsilon$ from $\partial D_{\epsilon}^{m-1}$ and $\partial D^{m-1}$.
Definition 4.2.2 (Barricade). A collection $B_{I}=\left\{B_{i}=\left[0, s_{i}\right] \times D^{m-1}\right\}_{i \in I}$ of pairwise disjoint flow boxes for $\Sigma$ is a barricade for $v$ if for each $B_{i}$ there exist a small constant $\epsilon_{i}>0$ and a shrinkage $B_{i}^{\epsilon_{i}}$ and the following hold:
${ }^{(*)}$ each flow line of $v$ intersects some $B_{i}^{\epsilon}$ and for any $x \in \Sigma$ which neither is a singularity of $v$ nor is contained in any $B_{i}^{\epsilon}$, the flow line of $v$ passing through $x$ enters some $B_{i}^{\epsilon}$ or limits to some Morse singularity in forward time (resp. in backward time).
(**) $B_{I}$ is locally finite, i.e., each compact subset of $\Sigma$ intersects only a finite number of $B_{i}$.

The following theorem of Wilson [Wil66, Theorem A], slightly adapted to our situation, guarantees the existence of barricades:

Theorem 4.2.3 (Wilson). A vector field $v$ on a manifold $\Sigma$ with only Morse singularities has a barricade $B_{I}$.

The following lemma on taking refinements is immediate:
Lemma 4.2.4 (Refinement). Given a flow box $B=\left[0, s_{0}\right] \times D^{m-1}$, a properly embedded submanifold $Z$ of $D^{m-1}$, and $\epsilon, \delta>0$ small, there exists a finite disjoint collection $\mathscr{B}$ of flow boxes $B_{1 j}, j=1, \ldots, k_{1}$, and $B_{2 j}, j=1, \ldots, k_{2}$, such that:
(1) $B_{1 j} \subset\left(0, s_{0} / 2\right) \times D^{m-1}$ and $B_{2 j} \subset\left(s_{0} / 2, s_{0}\right) \times D^{m-1}$;
(2) $\pi_{D^{m-1}}\left(\cup_{j=1}^{k_{1}} B_{1 j}\right)$ does not intersect $Z$ and $\pi_{D^{m-1}}\left(\cup_{j=1}^{k_{2}} B_{2 j}\right)$ is contained in a $\delta$-neighborhood of $Z$; and
(3) $\mathscr{B}$ is a barricade for the shrinkage $B^{\epsilon}$.

Here $\pi_{D^{m-1}}: B \rightarrow D^{m-1}$ is the projection onto the second factor.
4.3. Proof of Theorem 1.2.3 in dimension 3. Given any closed surface $\Sigma \subset$ $(M, \xi)$, it is well-known (see e.g. [Gei08, Section 4.6]) that, after a $C^{\infty}$-small perturbation, we can assume that $\Sigma_{\xi}$ has only Morse type singularities. By Theorem 4.2.3, a barricade $B_{I}=\left\{B_{i}=\left[0, s_{i}\right] \times\left[0, t_{i}\right]\right\}_{i \in I}$ exists for $\Sigma_{\xi}$; moreover $I$ can be taken to be finite since $\Sigma$ is closed. Each $B_{i}$ has a standard contact neighborhood with parameters $s_{i}, t_{i}, z_{i}$, where $z_{i}>0$ is small and we construct a $C^{0}$-small modification $\Sigma^{\vee}$ of $\Sigma$ by replacing every $B_{i}$ by the plug $B_{i}^{\vee}$. The characteristic foliation $\Sigma_{\xi}^{\vee}$ satisfies Conditions (M1)-(M3) of Proposition 2.2.6 and is Morse.

After a further $C^{\infty}$-small perturbation if necessary, $\Sigma^{\vee}$ can be made Morse ${ }^{+}$: If there exists a "retrogradient" flow line $\ell$ from a negative index 1 singularity to a positive index 1 singularity, we take a flow box $B=\{0\} \times\left[0, s_{0}\right] \times\left[0, t_{0}\right] \subset \Sigma^{\vee}$ such that $\ell$ intersects $B^{\epsilon}$ exactly once and $B$ has a standard contact neighborhood with parameters $s_{0}, t_{0}, z_{0}$. The retrogradient flow line can be eliminated by taking a small nonnegative function $h: B \rightarrow \mathbb{R}_{\geq 0}$ with support on $B^{\epsilon / 2}$ and replacing $B$ by $z=h(s, t)$; the modification has the effect of pushing the holonomy from $s=0$ to $s=s_{0}$ in the negative $t$-direction.

Since $\Sigma^{\vee}$ is now Morse ${ }^{+}$, it is Weinstein convex by Proposition 2.3.3.
4.4. Installing and uninstalling plugs. The construction of a plug $B^{\vee}$ was sufficient to prove Theorem 1.2.3 in dimension 3. In order to prove Theorem 1.2.5, we also need to interpolate between $B$ and $B^{\vee}$ with some control of the intermediate dynamics. We now explain this procedure.

Let $\left(M=\left[-z_{0}, z_{0}\right] \times\left[0, s_{0}\right] \times\left[0, t_{0}\right], \xi=\operatorname{ker}\left(d z+e^{s} d t\right)\right)$ be as before and let $B_{z}:=\{z\} \times\left[0, s_{0}\right] \times\left[0, t_{0}\right]$. Replace $B_{z_{0} / 2}$ by a plug $B_{z_{0} / 2}^{\vee}$ such that its analogously defined $z$-height satisfies $\mathscr{Z}\left(B_{z_{0} / 2}^{\vee}\right) \ll \frac{z_{0}}{2}$ and in particular $B_{z_{0} / 2}^{\vee}$ is still contained in $M$.

For the moment consider the PL model of the plug $B_{z_{0} / 2}^{\vee}$, i.e., each mushroom $Z$ involved in the construction is replaced by the corresponding $Z_{\mathrm{PL}}$. It is fairly straightforward to foliate the regions bounded between $B_{0}$ and $B_{z_{0} / 2}^{\vee}$ and between $B_{z_{0} / 2}^{\vee}$ and $B_{z_{0}}$ by a family of PL surfaces; see Figure 4.4.1 for a schematic picture. Then one can apply the smoothing scheme from $\S 3.2$ to smooth the corners of the


Figure 4.4.1. The interpolation between $B_{0}, B_{z_{0} / 2}^{\vee}$ and $B_{z_{0}}$; the $z$-height of $B_{z_{0} / 2}^{\vee}$ not drawn to scale.
leaves simultaneously and obtain the desired foliation $M=\cup_{0 \leq a \leq z_{0}} \widetilde{B}_{a}$, where $\widetilde{B}_{0}=B_{0}, \widetilde{B}_{z_{0}}=B_{z_{0}}$, and $\widetilde{B}_{z_{0} / 2}$ is the smoothed version of $B_{z_{0} / 2}^{\vee}$.

To analyze the dynamics of $\left(\widetilde{B}_{a}\right)_{\xi}$ for each $a \in\left[0, z_{0}\right]$, we introduce the partially defined and possibly multiple-valued holonomy map $\rho_{a}: \partial_{-} \widetilde{B}_{a} \rightarrow \partial_{+} \widetilde{B}_{a}$, where $\partial_{+} \widetilde{B}_{a}$ (resp. $\partial_{-} \widetilde{B}_{a}$ ) is the side $s=s_{0}$ (resp. $s=0$ ) of $\widetilde{B}_{a}$ as before: Given $x \in \partial_{-} \widetilde{B}_{a}$, if there exists a possibly broken partial flow line of $\left(\widetilde{B}_{a}\right)_{\xi}$ starting from $x$ and ending at $y \in \partial_{+} \widetilde{B}_{a}$, then $y \in \rho_{a}(x)$. Note that such $y$ may not be unique. If there is no such flow line, then $\rho_{a}(x)$ is not defined.

Define the internal discrepancy

$$
\left\|\rho_{a}\right\|:=\sup _{x \in \partial_{-} \widetilde{B}_{a}}\left|t(x)-t\left(\rho_{a}(x)\right)\right|
$$

where $t(x)$ refers to the $t$-coordinate of the point $x ;\left|t(x)-t\left(\rho_{a}(x)\right)\right|=0$ if $\rho_{a}(x)$ is not defined; and the supremum is taken over all possible $\rho_{a}(x)$ if $\rho_{a}$ is not singlevalued at $x$.

The following lemma will be important for our applications.
Lemma 4.4.1. The internal discrepancies $\sup _{0 \leq a \leq z_{0}}\left\|\rho_{a}\right\| \rightarrow 0$ as $N \rightarrow \infty$.
Proof. The lemma is not a statement about blocking and is rather a statement about the $t$-widths of the mushrooms $Z_{k, l}$ : We will treat the case where $a \in\left[0, \frac{z_{0}}{2}\right]$. If a flow line enters a box $\square_{k, l}$ along the bottom and exits from the top, the maximum it is moved in the $t$-direction is the width $\frac{3 t_{0}}{N}$ of the box. Since a flow line or broken flow line passes through at most 2 boxes, $\left\|\rho_{a}\right\| \leq \frac{6 t_{0}}{N}$.

We call the foliation from $\widetilde{B}_{0}=B_{0}$ to $\widetilde{B}_{z_{0} / 2}$ installing a plug and the foliation from $\widetilde{B}_{z_{0} / 2}$ to $\widetilde{B}_{z_{0}}=B_{z_{0}}$ uninstalling a plug. Then Lemma 4.4.1 basically says that neither installing nor uninstalling a plug affects the local holonomy by much. For the rest of Section 4, we assume that $N \gg 0$ without further mention.

The following is based on the construction of mushrooms in $\S 3.2$ and its slight generalization to 1-parameter families:

Lemma 4.4.2. $\left(\widetilde{B}_{a}\right)_{\xi}$, $a \in\left[0, z_{0}\right]$, is gradient-like with respect to a 1 -Morse function $f_{a}: \widetilde{B}_{a} \rightarrow \mathbb{R}$ which agrees with s on $\partial \widetilde{B}_{a}$.
4.5. Proof of Theorem 1.2.5 in dimension 3. Consider a contact structure $\xi$ on $\Sigma \times[0,1]$ such that $\Sigma \times\{0,1\}$ is Morse ${ }^{+}$in the sense of Definition 2.3.1. The goal is to show that up to an isotopy relative to the boundary, $\left(\left(\Sigma_{t}\right)_{\xi}\right)_{t \in[0,1]}$ is a 1-Morse family, where $\Sigma_{t}:=\Sigma \times\{t\}$.

Define $L:=\cup_{t \in[0,1]}\left\{x \in \Sigma_{t} \mid \xi_{x}=T_{x} \Sigma_{t}\right\}$. Up to a $C^{\infty}$-small perturbation of $\xi$, we can assume that $L$ satisfies the following:
(S1) $L$ is a properly embedded 1 -submanifold;
(S2) the singularities of $\left(\Sigma_{t}\right)_{\xi}$ are 1-Morse for all $t$; and
(S3) the restricted coordinate function $\left.t\right|_{L}: L \rightarrow[0,1]$ is Morse and all its critical points have distinct critical values.
Suppose $0<a_{1}<\cdots<a_{m}<1$ are the critical values of $\left.t\right|_{L}$, which we assume to be irrational. For each $t \in[0,1]$ there exists a barricade $B_{I^{t}}$ for $\Sigma_{t}$; moreover $B_{I^{t}}$ is a barricade for any vector field that is sufficiently close to $\left(\Sigma_{t}\right)_{\xi}$. By the compactness of $[0,1]$, there exists an integer $K \gg 0$ such that, for $i=0,1, \ldots, K$, $B_{I_{i}}, I_{i}=I^{i / K}$, is a barricade for all $\Sigma_{t}, t \in\left[\frac{i-1}{K}, \frac{i+1}{K}\right] \cap[0,1]$. Note that, for each $a_{j}$, there exist unique $j_{+}, j_{-}$such that $j_{+}=j_{-}+1$ and $\frac{j_{-}}{K}<a_{j}<\frac{j_{+}}{K}$.

Let $\pi: \Sigma \times[0,1] \rightarrow \Sigma$ be the projection onto the first factor.
We claim that for each $i=0, \ldots, K-1$ there are refinements of $B_{I_{i}}$ and $B_{I_{i+1}}$ (by abuse of notation we keep the same notation for the refinements) such that

$$
\begin{equation*}
\pi\left(B_{I_{i+1}}\right) \cap \pi\left(B_{I_{i}}\right)=\varnothing \tag{4.5.1}
\end{equation*}
$$

and moreover we may choose the refinement so that the new $B_{I_{i+1}}$ remains a barricade for all $\Sigma_{t}, t \in\left[\frac{i}{K}, \frac{i+2}{K}\right] \cap[0,1]$. The claim follows from viewing $B_{I_{i}}$ and $B_{I_{i+1}}$ as thin neighborhoods of collections $\gamma_{i}$ and $\gamma_{i+1}$ of arcs, taking their intersection $Z=\gamma_{i} \cap \gamma_{i+1}$ which we may take to be transverse, and applying Lemma 4.2.4.

We divide the proof into several steps.
STEP 1. From $\Sigma_{0}$ to $\Sigma_{1 /\left(N^{\prime} K\right)}^{\vee}$ where $N^{\prime}>0$ is a large integer.
Consider a flow box $B_{i}=\left[0, s_{0}\right] \times\left[0, t_{0}\right]$ of $B_{I_{0}}$. Let $\partial_{+} B_{i}=\left\{s_{0}\right\} \times\left[0, t_{0}\right]$ and $\partial_{-} B_{i}=\{0\} \times\left[0, t_{0}\right]$.

For each positive integer $r$, define the external holonomy $\widehat{\rho}_{i, r}: \partial_{+} B_{i} \rightarrow \partial_{-} B_{i}$ — a multiple-valued, partially defined, $r$ th return map from $\partial_{+} B_{i}$ to $\partial_{-} B_{i}$ of $\left(\Sigma_{0}\right)_{\xi}$ - as follows: For any $x \in \partial_{+} B_{i}$, a point $y \in \partial_{-} B_{i}$ is in the image $\widehat{\rho}_{i, r}(x)$ if there exists a possibly broken flow line $c:[0,1] \rightarrow \Sigma_{0}$ such that $c(0)=x, c(1)=y$, and $c$ passes through $\operatorname{int}\left(B_{i}\right)(r-1)$ times. Of course $\widehat{\rho}_{i, r}$ is not necessarily defined on all of $\partial_{+} B_{i}$ and when it is defined, it is not necessarily single-valued.

Since $\left(\Sigma_{0}\right)_{\xi}$ is Morse by assumption, (A) there exists $\delta>0$ such that

$$
\left\|\widehat{\rho}_{i, r}\right\|:=\inf _{x \in \partial_{+} B_{i}}\left|t(x)-t\left(\widehat{\rho}_{i, r}(x)\right)\right|>\delta,
$$

where we are taking $t\left(\widehat{\rho}_{i, r}(x)\right)=\infty$ if $\widehat{\rho}_{i, r}(x)$ does not exist. Otherwise, there is a sequence of points $x_{j} \in \partial_{+} B_{i}$ such that $\left|t\left(x_{j}\right)-t\left(\widehat{\rho}_{i, r}\left(x_{j}\right)\right)\right| \rightarrow 0$ and the compactness of the sequence of broken flow lines gives us $x_{\infty} \in \partial_{+} B_{i}$ such that
$\left|t\left(x_{\infty}\right)-t\left(\widehat{\rho}_{i, r}\left(x_{\infty}\right)\right)\right|=0$, which contradicts (M3) from Proposition 2.2.6. Moreover, the Morse condition implies that (B) there exists $r_{0}>0$ finite such that $\left\|\widehat{\rho}_{i, r_{0}}\right\|=\infty$ for all $i \in I$.

We then install plugs $B_{I_{0}}^{\vee}$ on $B_{I_{0}}$ as described in $\S 4.1$ and $\S 4.4$ and obtain a foliation between $\Sigma_{0}$ and $\Sigma_{0}^{\vee}$. As long as we take $N \gg 0$, i.e., the individual mushrooms are very small, the internal discrepancies are $\ll \delta$ by Lemma 4.4.1. Together with (A) and (B), it follows that all the leaves of the foliation are Morse.

For convenience we assume that $\Sigma_{0}^{\vee}$ agrees with $\Sigma_{0}$ on the complement of $B_{I_{0}}$, and that the difference is contained in a small invariant neighborhood of $B_{I_{0}}$. Also, below we construct 1-parameter families of embedded surfaces that are disjoint away from a subset on which they all agree; the perturbation into a family of disjoint embedded surfaces is done by flowing in the transverse direction for a short time and will not be done explicitly.

In order to interpolate between $\Sigma_{0}^{\vee}$ and $\Sigma_{1 /\left(N^{\prime} K\right)}^{\vee}$ for a large integer $N^{\prime}>0$ which we take to be odd, we use $B_{I_{1}}$ satisfying Eq. (4.5.1). If $N^{\prime} \gg 0$, then there is a 1-parameter family of embedded surfaces $F_{s} \subset \Sigma \times\left[0, \frac{1}{N^{\prime} K}\right], s \in[0,1]$, such that $F_{0}=\Sigma_{0} \backslash N\left(B_{I_{0}}\right), F_{1} \cap \Sigma_{1 /\left(N^{\prime} K\right)} \supset N\left(B_{I_{1}}\right) \times\left\{\frac{1}{N^{\prime} K}\right\}, \partial F_{s}=\partial N\left(B_{I_{0}}\right) \times\{0\}$ for all $s \in[0,1]$, the interiors of $F_{s}$ are disjoint, and the $\left(F_{s}\right)_{\xi}, s \in[0,1]$, are $\epsilon$-close to one other so that $B_{I_{0}}$ is a barricade for all $\left(N\left(B_{I_{0}}\right) \times\{0\}\right) \cup F_{s}$; in particular, no new singularities are introduced in this process. The barricading condition can be guaranteed by having chosen $N^{\prime} \gg 0$. See the upper-left corner of Figure 4.5.1 for an illustration of this procedure. By the barricading condition the surfaces $\left(N\left(B_{I_{0}}\right) \times\{0\}\right)^{\vee} \cup F_{s}$ are Morse for all $s \in[0,1]$, where $\left(N\left(B_{I_{0}}\right) \times\{0\}\right)^{\vee}$ is $N\left(B_{I_{0}}\right) \times\{0\}$ with $B_{I_{0}}^{\vee}$ installed.


Figure 4.5.1. Interpolation between $\Sigma_{0}$ and $\Sigma_{1 /\left(N^{\prime} K\right)}^{\vee}$ by Morse surfaces. The blue parts represent $B_{I_{i}}, i=0,1$.

Next we install a plug on $B_{I_{1}} \times\left\{\frac{1}{N^{\prime} K}\right\} \subset\left(N\left(B_{I_{0}}\right) \times\{0\}\right)^{\vee} \cup F_{1}$, uninstall the plug on $B_{I_{0}}^{\vee}$, and lift the resulting surface up to $\Sigma_{1 /\left(N^{\prime} K\right)}$, as shown in the upperright, lower-right, and lower-left corners of Figure 4.5.1, respectively. Moreover all the intermediate surfaces are Morse by analogous reasons. This finishes our construction of the foliation from $\Sigma_{0}$ to $\Sigma_{1 /\left(N^{\prime} K\right)}^{\vee}$.

STEP 2. From $\Sigma_{1 /\left(N^{\prime} K\right)}^{\vee}$ to $\Sigma_{1_{-} / K}^{\vee}$, where $1_{-} / K<a_{1}<1_{+} / K$.
Switching back and forth between $B_{I_{0}}$ and $B_{I_{1}}$, we similarly construct the Morse foliation from $\Sigma_{1 /\left(N^{\prime} K\right)}^{\vee}$ to $\Sigma_{2 /\left(N^{\prime} K\right)}^{\vee}$, from $\Sigma_{2 /\left(N^{\prime} K\right)}^{\vee}$ to $\Sigma_{3 /\left(N^{\prime} K\right)}^{\vee}$, and so on as in Step 1, until we get to $\Sigma_{1 / K}^{\vee}$. Between $\Sigma_{1 / K}^{\vee}$ and $\Sigma_{2 / K}^{\vee}$ we use $B_{I_{1}}$ and $B_{I_{2}}$. Continuing in this manner we get to $\Sigma_{1_{-/ K}}^{\vee}$.
STEP 3. From $\Sigma_{1_{-} / K}^{\vee}$ to $\Sigma_{1_{+} / K}^{\vee}$.
The only modification needed in this step is due to the fact that the vector fields $\left(\Sigma_{1_{-} / K}\right)_{\xi}$ and $\left(\Sigma_{1_{+} / K}\right)_{\xi}$ are not $C^{\infty}$-close to each other in the usual sense. Rather, one observes either the birth or the death of a pair of nearby Morse singularities as we go from $\left(\Sigma_{1_{-} / K}\right)_{\xi}$ to $\left(\Sigma_{1_{+} / K}\right)_{\xi}$. In either case, we slightly modify the notion of barricades $B_{I_{1} \pm}$ so that the unique (short) flow line connecting the pair of Morse singularities is the only flow line that does not pass through $B_{I_{1_{ \pm}}}$. Similar remarks apply to all $a_{i}, 1 \leq i \leq m$.

STEP 4. From $\Sigma_{(K-1) / K}^{\vee}$ to $\Sigma_{1}$.
In this final step, the only new ingredient is to uninstall the plugs as we go from $\Sigma_{1}^{\vee}$ to $\Sigma_{1}$. By assumption $\Sigma_{1}$ is Morse and in fact convex. Hence by the same holonomy bound as in Step 1, all the intermediate surfaces are Morse.

Finally we have foliated $\Sigma \times[0,1]$ by surfaces of the form $\Sigma_{t}$ which are all Morse. The only obstruction to convexity occurs when $\left(\Sigma_{t}\right)_{\xi}$ is Morse but not Morse $^{+}$and this corresponds to a bypass attachment (see Proposition 8.3.2). This concludes the proof of Theorem 1.2.5 in dimension 3.
4.6. Further remarks. Compared to earlier groundbreaking works of Bennequin [Ben83] and Eliashberg [Eli92], convex surface theory is a more systematic framework for studying embedded surfaces in contact 3 -manifolds. It is sufficiently powerful that basically all known classification results of contact structures or Legendrian knots in this dimension follow from this theory.

The only "drawback" of convex surface theory, at least in its original form [Gir91, Gir00], is that the monster of dynamical systems on surfaces is always lurking behind the story. More precisely, if one just wants to classify contact structures or Legendrian knots up to isotopy, then the problem often reduces to a combinatorial one by combining Giroux's theory with, say, the bypass approach of [Hon00]. However, if one wants to obtain higher homotopical information of the space of contact structures (say $\pi_{n}$ for $n \geq 1$ ), then some serious work on higher codimensional degenerations of Morse-Smale flows seems inevitable.

As an example, in [Eli92] Eliashberg outlined the proof that the space of tight contact structures on $S^{3}$ is homotopy equivalent to $S^{2}$. This particular result is based on the study of characteristic foliations on $S^{2} \subset S^{3}$, which is particularly simple since we never have periodic orbits. In more general contact manifolds such as $T^{3}$, one cannot necessarily rule out periodic orbits from characteristic foliations, and hence the bifurcation theory quickly becomes unwieldy (the work
[Ngo] probably comes close to the limit of what one can do). However, in light of our reinterpretation/simplification of Giroux's theory, it suffices to understand the space of Morse gradient vector fields, instead of general Morse-Smale vector fields.

We hope our techniques can be applied to future studies of homotopy types of the space of contact structures. This topic however will not be pursued any further in this paper.

## 5. CONSTRUCTION OF MUSHROOMS IN DIMENSION $>3$

The goal of this section is to generalize the construction of mushrooms in dimension 3 in Section 3 to higher dimensions.

Notation. Throughout this section, we will write $Z^{3} \subset \mathbb{R}^{3}$ for the mushroom constructed in Section 3 and write $Z$ for the higher-dimensional mushroom to be constructed.
5.1. Introduction. We first introduce some notation which will be used throughout this paper.

Definition 5.1.1 (Contact handlebodies and generalized contact handlebodies).
(1) A contact handlebody over a Weinstein domain $(X, \mu)$ is a contact manifold contactomorphic to

$$
\left([0, C]_{t} \times X, \operatorname{ker}(d t+\mu)\right)
$$

where $C>0$ is the thickness of the handlebody.
(2) A generalized contact handlebody over a Weinstein domain ( $X, \mu$ ) is a contact manifold contactomorphic to

$$
\left\{(t, x) \mid f_{0}(x)<t<f_{1}(x)\right\} \subset\left(\mathbb{R}_{t} \times X, \operatorname{ker}(d t+\mu)\right)
$$

where there exists a 1-parameter family $f_{t}: X \rightarrow \mathbb{R}, t \in[0,1]$, of smooth functions such that $f_{t}(x)<f_{t^{\prime}}(x)$ for all $t<t^{\prime}, x \in X$ and the graphs $\left\{t=f_{t_{0}}(x)\right\}$ are Weinstein for all $t_{0} \in[0,1]$.

A contact handlebody is a compact contact manifold with a contact form such that all its Reeb orbits are chords of the same length and a generalized contact handlebody is one such that that all the Reeb orbits are chords but they need not have the same length.

Let $(W, \lambda)$ be a complete Weinstein manifold of dimension $2 n-2>0$ and $\mathbb{R}_{t} \times W$ be the contactization of $W$ with contact form $\beta=d t+\lambda$. Let $W^{c} \subset W$ be a compact subdomain such that $W=W^{c} \cup\left([0, \infty)_{\tau} \times \Gamma\right), \Gamma:=\partial W^{c}$ is the contact boundary, and $[0, \infty)_{\tau} \times \Gamma$ is the positive half-symplectization of $\Gamma$. Let $\eta:=\left.\lambda\right|_{\Gamma}$ be the contact form on $\Gamma$; then $\left.\lambda\right|_{[0, \infty) \times \Gamma}=e^{\tau} \eta$. For $\tau^{\prime}>0$ we also define

$$
\begin{equation*}
W_{\tau^{\prime}}^{c}:=W^{c} \cup\left(\left[0, \tau^{\prime}\right] \times \Gamma\right) \tag{5.1.1}
\end{equation*}
$$

The ambient contact manifold of a mushroom is

$$
\left(M=\mathbb{R}_{z, s, t}^{3} \times W, \xi=\operatorname{ker} \alpha\right), \quad \alpha=d z+e^{s} \beta
$$

The hypersurface on which we construct the mushroom is $\Sigma=\{z=0\} \subset M$ with characteristic foliation $\Sigma_{\xi}=\partial_{s}$.

Remark 5.1.2. For ease of notation, we will not distinguish between the characteristic foliation, which is an oriented singular line field, and a trivializing vector field.

A mushroom $Z$ with contact handlebody profile $H=\left(\left[0, t_{0}\right] \times W_{\tau_{0}}^{c}, d t+\lambda\right)$, $\tau_{0}>0$, is constructed by first taking the product hypersurface $Z^{3} \times W^{c}$, where $Z_{P L}^{3}$ has base $\left[0, s_{0}\right] \times\left[0, t_{0}\right]$, and then damping out the $Z^{3}$-factor on $W_{\tau_{0}}^{c}-W^{c}$. Roughly speaking, the goal is to fold $\Sigma$ using $Z$, so that the resulting characteristic foliation cannot pass through a region which approximates $H$.

Remark 5.1.3. One can think of the constructions in Section 3 as a special case where $W$ is a point and $H=\left[0, t_{0}\right]$ is equipped with the contact form $d t$.
5.2. Product hypersurface. Recall that in Section 3 we constructed the mushroom

$$
Z^{3} \subset\left(\mathbb{R}_{z, s, t}^{3}, \operatorname{ker}\left(d z+e^{s} d t\right)\right)
$$

which agrees with $\mathbb{R}_{s, t}^{2}$ outside of a rectangle $\square=\left[0, s_{0}\right] \times\left[0, t_{0}\right]$. Let $Z_{\xi}^{3}$ be the characteristic foliation on $Z^{3}$.

We will compute the characteristic foliation $Z_{\xi}^{\prime}$ on the product hypersurface $Z^{\prime}:=Z^{3} \times W^{c} \subset M$.

Choose vector fields $v$ on $Z^{3}$, defined away from the singularities of $Z_{\xi}^{3}$, such that $\left.\alpha\right|_{Z^{3}}(v)=1$ and $w$ on $W^{c}$, defined away from the zero set of $\lambda$, such that $\lambda(w)=1$.

Lemma 5.2.1. Away from the zeros of $\left.\alpha\right|_{Z^{3}}$ and $\lambda$, the characteristic foliation $Z_{\xi}^{\prime}$ is given by

$$
\begin{equation*}
Z_{\xi}^{\prime}=\mathbb{R}\left\langle Z_{\xi}^{3}+d z \wedge d s\left(Z_{\xi}^{3}, v\right) X_{\lambda}\right\rangle \tag{5.2.1}
\end{equation*}
$$

where $X_{\lambda}$ is the Liouville vector field of $\lambda$.
Proof. One can easily check that

$$
T\left(Z^{3} \times W^{c}\right) \cap \xi=\mathbb{R}\left\langle Z_{\xi}^{3}, w-e^{s} v, \text { ker } \lambda\right\rangle
$$

Basically the calculation of $Z_{\xi}^{\prime}$ is reduced to computing the kernel $K=a X+$ $b Y+c Z$ of the 3 -dimensional vector space $\mathbb{R}\langle X, Y, Z\rangle$ with a maximally nondegenerate alternating 2 -form $\langle\cdot, \cdot\rangle$. One can easily verify that

$$
K=\langle Y, Z\rangle X+\langle Z, X\rangle Y+\langle X, Y\rangle Z
$$

works. We have $e^{-s} d \alpha=d s \wedge d t+d s \wedge \lambda+d \lambda$, and if we write $\langle\cdot, \cdot\rangle:=e^{-s} d \alpha(\cdot, \cdot)$ and $Z_{\xi}^{\prime}=a X+b Y+c Z+d A$, where $X=Z_{\xi}^{3}, Y=X_{\lambda}, Z=w-e^{s} v$, and $A \in \operatorname{ker} \lambda$ and is not parallel to $X_{\lambda}$, then $d=0$ since otherwise there exists
$B \in \operatorname{ker} \lambda$ such that $\langle A, B\rangle \neq 0$ and $\langle w, B\rangle=0$. The remaining/relevant part of the pairing is given as follows:

$$
\begin{aligned}
\left\langle Z_{\xi}^{3}, X_{\lambda}\right\rangle & =0 \\
\left\langle Z_{\xi}^{3}, w-e^{s} v\right\rangle & =d s\left(Z_{\xi}^{3}\right)-e^{s} d s \wedge d t\left(Z_{\xi}^{3}, v\right) \\
\left\langle X_{\lambda}, w-e^{s} v\right\rangle & =1
\end{aligned}
$$

Hence $Z_{\xi}^{\prime}=K=Z_{\xi}^{3}-\left(d s\left(Z_{\xi}^{3}\right)-e^{s} d s \wedge d t\left(Z_{\xi}^{3}, v\right)\right) X_{\lambda}$.
Finally, since $\alpha-\left(d z+e^{s} d t\right)=0$ when evaluated on vectors on $Z^{3}$ and hence

$$
\left(d s \wedge \alpha+d z \wedge d s-e^{s} d s \wedge d t\right)\left(Z_{\xi}^{3}, v\right)=0
$$

it follows that $Z_{\xi}^{\prime}=Z_{\xi}^{3}+d z \wedge d s\left(Z_{\xi}^{3}, v\right) X_{\lambda}$.
At the zeros of $\left.\alpha\right|_{Z^{3}}$ and $\lambda$, Eq. (5.2.1) can be interpreted as saying that $Z_{\xi}^{\prime}$ contains the limit of the right-hand side as the points on $Z^{3} \times W^{c}$ approach the zero.

Remark 5.2.2. Lemma 5.2 .1 is rather general and holds for $Z^{3}$ replaced by any surface in $\mathbb{R}_{z, s, t}^{3}$.
5.3. Dynamics of $Z_{\xi}^{\prime}$. We now investigate the dynamics of $Z_{\xi}^{\prime}$.

Let us first consider the PL case $Z_{P L}^{\prime}=Z_{P L}^{3} \times W^{c}$.
Lemma 5.3.1. The flow lines of $\left(Z_{P L}^{\prime}\right)_{\xi}$ passing through $\{-1\}_{s} \times\left(0, t_{0}\right)_{t} \times W^{c}$ eventually limit to a negative singularity of $\left(Z_{P L}^{\prime}\right)_{\xi}$ and in particular do not leave $Z_{P L}^{\prime}$.
Proof. The lemma follows from two observations: (i) Since $d z \wedge d s\left(\left(Z_{P L}^{3}\right)_{\xi}, v\right)$ is positive on $P_{4}$, negative on $P_{2}$, and vanishes on $P_{0} \cup P_{1} \cup P_{3}$, the term $d z \wedge$ $d s\left(\left(Z_{P L}^{3}\right)_{\xi}, v\right) X_{\lambda}$ in Eq. (5.2.1) is a positive multiple of $X_{\lambda}$ on $P_{4}$, a negative multiple of $X_{\lambda}$ on $P_{2}$, and zero on $P_{0} \cup P_{1} \cup P_{3}$. [Sample sign calculation on $P_{4}$ (it is useful to refer to Figure 3.1.2): $\left(Z_{P L}^{3}\right)_{\xi}=-\partial_{s}$ and $v$, which we take to be parallel to the $z t$-plane, has positive $\partial_{z}$-component. Hence $d z \wedge d s\left(\left(Z_{P L}^{3}\right)_{\xi}, v\right)>0$ on $P_{4}$.] (ii) By Lemma 3.1.4, if a flow line of $\left(Z_{P L}^{\prime}\right)_{\xi}$ passes through $\{-1\}_{s} \times\left(0, t_{0}\right)_{t} \times W^{c}$, then its projection to $Z_{P L}^{3}$ only passes through $P_{0}, P_{1}, P_{2}$, and $P_{3}$.

Next we describe the smoothed version $Z_{\xi}^{\prime}$. We identify the singular points of $Z_{\xi}^{\prime}$ : Recall from Lemma 3.1.4 that $Z_{\xi}^{3}$ has four singular points $e_{ \pm}, h_{ \pm}$. By the sign calculations of $d z \wedge d s\left(\left(Z_{P L}^{3}\right)_{\xi}, v\right)$ from the proof of Lemma 5.3.1 and continuity, $d z \wedge d s\left(Z_{\xi}^{3}, v\right)>0$ on neighborhoods of $e_{+}, h_{+}$and $<0$ on neighborhoods of $e_{-}, h_{-}$. Hence for each singular point $x \in W^{c}$ of the Liouville vector field $X_{\lambda}$, there exist four singular points $e_{ \pm}^{x}, h_{ \pm}^{x}$ of $Z_{\xi}^{\prime}$ whose Morse indices are given by:

$$
\begin{array}{ll}
\operatorname{ind}\left(e_{+}^{x}\right)=\operatorname{ind}_{W}(x), & \operatorname{ind}\left(h_{+}^{x}\right)=\operatorname{ind}_{W}(x)+1, \\
\operatorname{ind}\left(e_{-}^{x}\right)=2 n-\operatorname{ind}_{W}(x), & \operatorname{ind}\left(h_{-}^{x}\right)=2 n-1-\operatorname{ind}_{W}(x),
\end{array}
$$

where $\operatorname{ind}_{W}(x)$ is the Morse index of $x \in W^{c} \subset W$ and we recall that $\operatorname{dim} W=$ $2 n-2$.

See the top left figure in Figure 5.3.1 for $Z_{\xi}^{3}$ and the regions indicating the signs of $d z \wedge d s\left(Z_{\xi}^{3}, v\right)$. The red (resp. blue, white) region indicates where $d z \wedge d s\left(Z_{\xi}^{3}, v\right)$ or $d z \wedge d s\left(S_{\tau, \xi}, v\right)$ is positive (resp. negative, zero).

Remark 5.3.2. In view of Remark 3.2.2, we may replace the nondegenerate singular points by birth-death singularities as in the top right of Figure 5.3.1. The advantage of the birth-death singularities is that Lemma 5.3.3 still holds but the singular points can be immediately eliminated; this will be useful for example when damping out in Section 5.4.


Figure 5.3.1. The top left is $Z_{\xi}^{3}$ and the top right is an alternate perturbation of $Z_{P L, \xi}^{3}$ corresponding to $\tau=0$. The top right, bottom right, and bottom left are $S_{\tau, \xi}$ for some as $\tau$ goes from 0 to $\tau_{0}$.

Let $\operatorname{Sk}(W)$ be the isotropic skeleton of $W^{c}$ with respect to $X_{\lambda}$. Let $\kappa_{1}>\kappa_{2}>$ $\kappa_{3}>0>\kappa_{4}>\kappa_{5}>\kappa_{6}$ with all $\kappa_{i}$ small as in Proposition 3.2.1, and let $a \geq 0$ be small. We define

$$
\begin{gathered}
I_{a}^{-}:=\{-1\} \times\left[\kappa_{4}, t_{0}+\kappa_{2}-a\right] \subset \mathbb{R}_{s, t}^{2}, \\
I_{a}^{+}:=\left\{s_{0}+1\right\} \times\left[\kappa_{5}-a, t_{0}+\kappa_{3}\right] \subset \mathbb{R}_{s, t}^{2},
\end{gathered}
$$

so that $I_{0}^{-}$(resp. $I_{0}^{+}$) is the maximal interval with the property that any flow line of $Z_{\xi}^{3}$ passing through the interval converges to a singularity of $Z_{\xi}^{3}$ in forward (resp. backward) time.

We now give a description of all the flow lines passing through $Z^{3} \times W^{c}$ :
Lemma 5.3.3 ( Description of all flow lines passing through $Z^{3} \times W^{c}$ ). There exist functions $\sigma_{0}^{-}, \sigma_{0}^{+}: W^{c} \rightarrow \mathbb{R}_{\geq 0}$, which vanish exactly on $\operatorname{Sk}(W)$ such that:
(1) each flow line of $Z_{\xi}^{\prime}$ passing through $I_{0}^{-} \times \operatorname{Sk}(W)$ (resp. $I_{0}^{+} \times \operatorname{Sk}(W)$ ) converges to a singularity of $Z_{\xi}^{\prime}$ in forward (resp. backward) time;
(2) for $x \in W^{c} \backslash \operatorname{Sk}(W)$, each flow line passing through $I_{\sigma_{0}^{-}(x)}^{-} \times\{x\}$ (resp. $\left.I_{\sigma_{0}^{+}(x)}^{+} \times\{x\}\right)$ converges to a singularity of $Z_{\xi}^{\prime}$ in forward (resp. backward) time;
(3) for $x \in W^{c} \backslash \operatorname{Sk}(W)$, each flow line passing through $\left(I_{0}^{-} \backslash I_{\sigma_{0}^{-}(x)}^{-}\right) \times\{x\}$ (resp. $\left.\left(I_{0}^{+} \backslash I_{\sigma_{0}^{+}(x)}^{+}\right) \times\{x\}\right)$ exits $Z^{3} \times W^{c}$ along $Z^{3} \times \partial W^{c}$ in finite forward (resp. backward) time;
(4) for $x \in W^{c}$, each flow line passing through $\{-1\} \times\left(t_{0}+\kappa_{2}, t_{0}+\kappa_{1}\right] \times\{x\}$ (resp. $\left.\left\{s_{0}+1\right\} \times\left[\kappa_{6}, \kappa_{5}\right) \times\{x\}\right)$ exits from either $\left\{s_{0}+1\right\} \times\left(t_{0}+\kappa_{3}, t_{0}+\right.$ $\left.\kappa_{1}\right] \times W^{c}\left(\right.$ resp. $\{-1\} \times\left[\kappa_{6}, \kappa_{4}\right) \times W^{c}$ ) or $Z^{3} \times \partial W^{c}$ in finite forward (resp. backward) time; in the former case, the $W^{c}$-coordinate $x^{\prime}$ of the exit point of the flow line is on the time $\geq 0$ flow line of $X_{\lambda}$ starting at $x$;
(5) for $x \in W^{c}$, each flow line passing through $\{-1\} \times\left[\kappa_{6}, \kappa_{4}\right) \times\{x\}$ (resp. $\left.\left\{s_{0}+1\right\} \times\left(t_{0}+\kappa_{3}, t_{0}+\kappa_{1}\right] \times\{x\}\right)$ exits from $\left\{s_{0}+1\right\} \times\left[\kappa_{6}, \kappa_{5}\right) \times W^{c}$ (resp. $\left.\{-1\} \times\left(t_{0}+\kappa_{2}, t_{0}+\kappa_{1}\right] \times W^{c}\right)$ in finite forward (resp. backward) time; the $W^{c}$-coordinate $x^{\prime}$ of the exit point of the flow line is on the time $\leq 0$ flow line of $X_{\lambda}$ starting at $x$;
(6) each flow line outside of $\mathbb{R}_{s} \times\left[\kappa_{6}, t_{0}+\kappa_{1}\right] \times W^{c}$ has trivial holonomy;
(7) all other flow lines are (i) flow lines between singularities, (ii) flow lines from a singularity to $Z^{3} \times \partial W^{c}$, or (iii) flow lines from $Z^{3} \times \partial W^{c}$ to a singularity.
Moreover, as $Z^{3} \rightarrow Z_{P L}^{3}$, all $\kappa_{i} \rightarrow 0$ and $\left|\sigma_{0}^{ \pm}\right|_{C^{0}} \rightarrow 0$.
Proof. This is an immediate consequence of Lemmas 5.2.1 and 5.3.1, taking the limit $Z^{3} \rightarrow Z_{P L}^{3}$, and a case-by-case analysis of the various regions of the top left figure of Figure 5.3.1.

Suppose the flow line passes through the red region times $W^{c}$. Then either the flow line exits from $Z^{3} \times \partial W^{c}$ or escapes to the white region times $W^{c}$. Once in the white region, the flow line either reaches $s=s_{0}+1$ or enters the blue region times $W^{c}$ and reaches a negative singularity.

Suppose the flow line passes through the white region (e.g., passes through $s=$ -1 ). Then the flow line reaches $s=s_{0}+1$, enters the blue region times $W^{c}$ (and hence reaches a negative singularity), or enters the red region times $W^{c}$ (already considered).

All $\kappa_{i} \rightarrow 0$ and $\left|\sigma_{0}^{ \pm}\right|_{C^{0}} \rightarrow 0$ as $Z^{3} \rightarrow Z_{P L}^{3}$ by construction.
Technically, the functions $\sigma_{0}^{ \pm}$account for the speed of convergence of flow lines of $\left(Z^{3}\right)_{\xi}$ towards its singularities and those of $X_{\lambda}$ in $W^{c}$ towards $\operatorname{Sk}(W)$.
5.4. Damping. In order for the mushroom to be the image of a continuous map $\Sigma \rightarrow M$, one must damp out the $Z^{3}$-fiber over $W_{\tau_{0}}^{c} \backslash \operatorname{int}\left(W^{c}\right)=\left[0, \tau_{0}\right] \times \Gamma$ as $\tau$ grows. Recall the notation from Section 5.1.

The damping procedure amounts to choosing an isotopy of surfaces $S_{\tau}, \tau \in$ $\left[0, \tau_{0}\right]$, in $\mathbb{R}_{z, s, t}^{3}$ from $S_{0}=Z^{3}$ to the flat $S_{\tau_{0}}=\mathbb{R}_{s, t}^{2}$. We take $S_{0}=Z^{3}$ to have
birth-death type singular points as in Figure 5.3.1; see Remark 5.3.2. In practice we also take $\tau_{0}>0$ to be arbitrarily small. We then set

$$
\mathscr{I}_{0}:=\cup_{0 \leq \tau \leq \tau_{0}}\left(S_{\tau} \times\{\tau\}\right) \subset \mathbb{R}_{z, s, t, \tau}^{4}
$$

and the actual hypersurface in $M$ will be $\mathscr{I}_{0} \times \Gamma$.
The PL model of $S_{\tau}$ is obtained by replacing $P_{0}$ by the rectangle $\left[0, s_{0}\right] \times$ $\left[-e^{-s_{0} / 2} z,-e^{-s_{0} / 2} z+t_{0}\right]$ as in Section 3.1, where the parameter $z$ ranges from $z_{0}$ to 0 as $\tau$ goes from 0 to $\tau_{0}$ (in other words, we are pushing the top face of the parallelepiped into the parallelepiped); its smoothing for $\tau>0$ will use a profile function $\phi$ such that $\phi^{\prime}<0$ everywhere so that there are no singularities of the characteristic foliation. We have

$$
T \mathscr{I}_{0}=\mathbb{R}\left\langle T S_{\tau}, \partial_{\tau}+f w_{0}\right\rangle,
$$

where $w_{0}=\partial_{z}-K_{0} \partial_{t}$ is parallel to $P_{2}$ and $P_{4}$ and $f \leq 0$ is a $\tau$-dependent smooth function on $\mathbb{R}_{z, s, t}^{3}$ which vanishes when $\tau$ is close to $\left\{0, \tau_{0}\right\}$ or $z=0$.

We are now ready to compute the characteristic foliation $\left(\mathscr{I}_{0} \times \Gamma\right)_{\xi}$. Let $S_{\tau, \xi}$ be the characteristic foliation on $S_{\tau}$, i.e., $\left.\alpha\right|_{\mathbb{R}^{3}}\left(S_{\tau, \xi}\right)=0$, and let $v$ be a vector field on $S_{\tau}$, defined away from the singularities of $S_{\tau, \xi}$, such that $\left.\alpha\right|_{\mathbb{R}^{3}}(v)=1$.
Lemma 5.4.1. The characteristic foliation $\left(\mathscr{I}_{0} \times \Gamma\right)_{\xi}$ is given by

$$
\begin{align*}
\left(\mathscr{I}_{0} \times \Gamma\right)_{\xi}= & S_{\tau, \xi}+d z \wedge d s\left(S_{\tau, \xi}, v\right)\left(\partial_{\tau}+f w_{0}\right)+f\left(d s\left(S_{\tau, \xi}\right) v-d s(v) S_{\tau, \xi}\right)  \tag{5.4.1}\\
& +e^{-\tau} f\left(-d s \wedge d t\left(S_{\tau, \xi}, v\right)+K_{0} d z \wedge d s\left(S_{\tau, \xi}, v\right)\right) R_{\eta}
\end{align*}
$$

on the subset of $\mathscr{I}_{0} \times \Gamma$ where $v$ is defined. Here $R_{\eta}$ is the Reeb vector field of $\eta$.
Proof. This is similar to the calculation of Lemma 5.2.1. We compute
$T\left(\mathscr{I}_{0} \times \Gamma\right) \cap \xi=\mathbb{R}\left\langle S_{\tau, \xi}, e^{\tau+s} v-R_{\eta}, \partial_{\tau}+f w_{0}+\left(-e^{-\tau-s}+K_{0} e^{-\tau}\right) f R_{\eta}\right.$, ker $\left.\eta\right\rangle$.
Next we have

$$
\begin{aligned}
\alpha & =d z+e^{s}\left(d t+e^{\tau} \eta\right) \\
e^{-s} d \alpha & =d s \wedge d t+e^{\tau} d s \wedge \eta+e^{\tau}(d \tau \wedge \eta+d \eta)
\end{aligned}
$$

Setting $X=S_{\tau, \xi}, Y=e^{\tau+s} v-R_{\eta}, Z=\partial_{\tau}+f w_{0}+\left(-e^{-\tau-s}+K_{0} e^{-\tau}\right) f R_{\eta}$,

$$
\begin{aligned}
\langle X, Y\rangle= & e^{\tau+s} d s \wedge d t\left(S_{\tau, \xi}, v\right)-e^{\tau} d s\left(S_{\tau, \xi}\right)=e^{\tau}\left(d z \wedge d s\left(S_{\tau, \xi}, v\right)\right), \\
\langle X, Z\rangle= & e^{\tau} d s\left(S_{\tau, \xi}\right)\left(-e^{-\tau-s} f\right)=-e^{-s} f d s\left(S_{\tau, \xi}\right), \\
\langle Y, Z\rangle= & e^{\tau+s}\left(e^{\tau} d s(v)\left(-e^{-\tau-s} f\right)\right)+e^{\tau}=e^{\tau}(1-f d s(v)) . \\
\left(\mathscr{I}_{1} \times \Gamma\right)_{\xi}= & (1-f d s(v)) S_{\tau, \xi}+e^{-\tau-s} f d s\left(S_{\tau, \xi}\right)\left(e^{\tau+s} v-R_{\eta}\right) \\
& +d z \wedge d s\left(S_{\tau, \xi}, v\right)\left(\partial_{\tau}+f w_{0}+\left(-e^{-\tau-s}+K_{0} e^{-\tau}\right) f R_{\eta}\right) .
\end{aligned}
$$

A rearrangement of the terms gives the lemma.
Note that Eq. (5.4.1) agrees with Eq. (5.2.1) at $\tau=0$. The first two terms of $\left(\mathscr{I}_{0} \times \Gamma\right)_{\xi}$ are analogous to those of $Z_{\xi}^{\prime}$; see Lemma 5.2.1. The third term $f\left(d s\left(S_{\tau, \xi}\right) v-d s(v) S_{\tau, \xi}\right)$ lies in ker $d s$ and, away from the corners,

- vanishes on $P_{1} \cup P_{3}$,
- has negative $\partial_{t}$-component on $P_{0} \cup P_{4}$, and
- has positive $\partial_{t}$-component on $P_{2}$.

See Figure 5.4.1. In other words, the third term, when we project out the $s$ - and $W$-directions, is a flow in the clockwise direction around $\partial P_{1}$ as seen in the picture. The last term of Eq. (5.4.1) has a substantial contribution in the $R_{\eta}$-direction when the damping happens quickly, i.e., when $\tau_{0}$ is small and $f$ is large. This is something we need to be careful about, but ultimately can be finessed away by stacking the mushrooms in a particular way in Section 7.


Figure 5.4.1. The vector field $f\left(d s\left(S_{\tau, \xi}\right) v-d s(v) S_{\tau, \xi}\right)$ is depicted in blue.
5.5. Description of the characteristic foliation of the mushroom. In this subsection we summarize the dynamics of the characteristic foliation of the mushroom $Z_{H}$ of $\Sigma=\{z=0\} \subset M=\mathbb{R}_{z, s, t}^{3} \times W$ with profile $H=\left[0, t_{0}\right] \times W_{\tau_{0}}^{c}$.
Definition 5.5.1. The mushroom of $\Sigma$ with profile $H$ is the hypersurface

$$
\begin{equation*}
Z_{H}:=\left(\Sigma \backslash\left(\square \times W_{\tau_{0}}^{c}\right)\right) \cup Z_{P L}^{\prime} \cup\left(\mathscr{I}_{0} \times \Gamma\right)_{P L}, \tag{5.5.1}
\end{equation*}
$$

modulo smoothing. (The smoothed versions do not have the subscripts PL.) The region $\square \times W_{\tau_{0}}^{c} \subset \Sigma$, where $\square=\left[0, s_{0}\right] \times\left[0, t_{0}\right]$, is the base of $Z_{H}$, and the region $\mathscr{I}_{0} \times \Gamma$ is the damping region.

Let $\tau_{0}^{\prime} \in\left(0, \tau_{0}\right)$, let $\kappa_{1}>\kappa_{2}>\kappa_{3}>0>\kappa_{4}>\kappa_{5}>\kappa_{6}$ with all $\kappa_{i}$ small as in Proposition 3.2.1, and let $\sigma_{1}^{ \pm}, \sigma_{2}^{ \pm}: W_{\tau_{0}^{\prime}}^{c} \rightarrow \mathbb{R}_{\geq 0}$ be functions such that:

- $\sigma_{1}^{ \pm}$vanishes exactly on $\operatorname{Sk}(W)$ and $\sigma_{2}^{ \pm}$vanishes on $W^{c}$;
- on $\left\{0 \leq \tau \leq \tau_{0}^{\prime}\right\}$, both $\sigma_{1}^{ \pm}=\sigma_{1}^{ \pm}(\tau)$ and $\sigma_{2}^{ \pm}=\sigma_{2}^{ \pm}(\tau)$ are strictly increasing and reach their maximum at $\tau=\tau_{0}^{\prime}$;
- $\sigma_{1}^{-}\left(\tau_{0}^{\prime}\right)+\sigma_{2}^{-}\left(\tau_{0}^{\prime}\right)=t_{0}+\kappa_{2}-\kappa_{4}$ and $\sigma_{1}^{+}\left(\tau_{0}^{\prime}\right)+\sigma_{2}^{+}\left(\tau_{0}^{\prime}\right)=t_{0}+\kappa_{3}-\kappa_{5}$.

As the smooth version of $Z_{H}$ limits to the PL version, all $\kappa_{i} \rightarrow 0$ and $\left|\sigma_{i}^{ \pm}\right|_{C^{0}} \rightarrow$ 0 on $W_{\tau_{0}}^{c}$.

We then define the compact submanifolds

$$
\begin{align*}
H_{\mathrm{in}} & :=\left\{(t, x) \mid x \in W_{\tau_{0}^{\prime}}^{c}, \kappa_{4}+\sigma_{2}^{-}(x) \leq t \leq t_{0}+\kappa_{2}-\sigma_{1}^{-}(x)\right\},  \tag{5.5.2}\\
H_{\mathrm{out}} & :=\left\{(t, x) \mid x \in W_{\tau_{0}^{\prime}}^{c}, \kappa_{5}+\sigma_{1}^{+}(x) \leq t \leq t_{0}+\kappa_{3}-\sigma_{2}^{+}(x)\right\}, \tag{5.5.3}
\end{align*}
$$

which approximate $H$ when all the smoothing parameters involved in the construction tend to 0 . See Figure 5.5.1. We use the notation $X^{\circ}$ (and also $\operatorname{int}(X)$ ) as in $H_{\text {in }}^{\circ}$ to denote the interior of a space $X$.

Note that $\partial H$ (which we assume has rounded corners) is convex; this follows from observing that $\left([-1,1]_{t} \times W, d t+\lambda\right)$ has contact vector field $t \partial_{t}+X_{\lambda}$.

Proposition 5.5.2. Assuming all the corner rounding parameters are sufficiently small, there exists a tubular neighborhood $[-\epsilon, \epsilon]_{\ell} \times \partial H$ of $\partial H=\{0\} \times \partial H$ and $H_{\text {in }}$ and $H_{\text {out }}$ that approximate $H$ such that:
(Z1) $\partial H_{\text {in }}, \partial H_{\text {out }} \subset[-\epsilon, \epsilon] \times \partial H$ are graphical over $\partial H$.
(Z2) $Z_{H, \xi}$ is " 1 -Morse" in the following sense: it satisfies (M1) and (M3) of Proposition 2.2.6 and
(M2') every flow line passing through $x \in Z_{H}$ with $Z_{H, \xi}(x) \neq 0$ converges to a singularity or goes to $\{s=+\infty\}$ in forward time and converges to a singularity or goes to $\{s=-\infty\}$ in backward time.
(Z3) Any flow line of $Z_{H, \xi}$ that passes through $H_{\mathrm{in}}^{\circ} \subset\{s=-1\}$ converges to a negative singularity of $Z_{H, \xi}$ in forward time. Similarly, any flow line of $Z_{H, \xi}$ that passes through $H_{\text {out }}^{\circ} \subset\left\{s=s_{0}+1\right\}$ converges to a positive singularity of $Z_{H, \xi}$ in backward time.
(Z4) Any flow line of $Z_{H, \xi}$ that does not pass through $H \cup([-\epsilon, \epsilon] \times \partial H) \subset$ $\{s=-1\}$ has trivial holonomy.
(Z5) There exists a Morse function $F$ on $\partial H$ such that $\partial H_{\xi}$ is gradient-like for $F$ (and hence flows "from $R_{+}(\partial H)$ to $R_{-}(\partial H)$ ") and such that any flow line of $Z_{H, \xi}$ that passes through $(\ell, x) \in[-\epsilon, \epsilon] \times \partial H \subset\{s=-1\}$ and does not converge to a singularity of $Z_{H, \xi}$ :
(1) passes through $\left(\ell^{\prime}, y\right) \in[-\epsilon, \epsilon] \times \partial H \subset\left\{s=s_{0}+1\right\}$ with $F(y) \geq$ $F(x)$; and
(2) is parallel to $X_{\lambda}\left(\right.$ resp. $\left.-X_{\lambda}\right)$ on $[-\epsilon, \epsilon] \times W_{+}^{c}\left(\right.$ resp. $\left.[-\epsilon, \epsilon] \times W_{-}^{c}\right)$ when projected to $[-\epsilon, \epsilon] \times \partial H$.
Here $W_{+}^{c}$ is the portion of $\partial H$ corresponding to $\left\{t_{0}\right\} \times W^{c}$ and $W_{-}^{c}$ is the portion corresponding to $\{0\} \times W^{c}$.
Proof. (Z1) is by construction. (Z4) is clear. (Z2)-(Z5) follow from Lemmas 5.3.3 and 5.4.1. In (Z5) we take $\{\tau\} \times \Gamma$ to be level sets of $F$ so that the component of $Z_{H, \xi}$ in the direction of the Reeb vector field $R_{\eta}$ vanishes on $d F$.

## 6. QUANTITATIVE STABILIZATION OF AN OPEN BOOK DECOMPOSITION FOR $S^{2 n-1}$

6.1. Some definitions. Let $M$ be a closed manifold. An open book decomposition (abbreviated OBD) of $M$ is a pair $(B, \pi)$, where $B \subset M$ is a closed codimension 2 submanifold and

$$
\pi: M \backslash B \rightarrow S^{1} \subset \mathbb{C}
$$

is a fibration which agrees with the angular coordinate $\theta$ on a neighborhood $B \times D^{2}$ of $B=B \times\{0\}$. We call $S_{\theta}:=\pi^{-1}\left(e^{i \theta}\right), e^{i \theta} \in S^{1}$, the pages of the OBD, and call $B$ the binding.


Figure 5.5.1. The shaded regions are $H_{\text {in }}$ and $H_{\text {out }}$, respectively. The area of the complements of $H_{\text {in }}$ and $H_{\text {out }}$ in the rectangles tend to 0 as all the parameters involved in the construction tend to 0 .

Let $\xi$ be a contact structure on $M$.

## Definition 6.1.1.

(1) An $O B D(B, \pi)$ is $\xi$-compatible if there exists a contact form $\alpha$ for $\xi$, called an adapted contact form, such that the Reeb vector field $R_{\alpha}$ of $\alpha$ is transverse to all the pages and is tangent to $B$, and $\left.\alpha\right|_{B}$ is a contact form on $B$. We also say " $\alpha$-compatible" or simply "compatible" if the contact structure is understood.
(2) An $\alpha$-compatible $(B, \pi)$ is strongly Weinstein if all its pages $\left(S_{\theta},\left.\alpha\right|_{S_{\theta}}\right)$ are Weinstein.

Let $(B, \pi)$ be an $\alpha$-compatible OBD. Let $\arg : S^{1} \rightarrow \mathbb{R} / 2 \pi \mathbb{Z}$ be the map $e^{i \theta} \mapsto \theta$. Then define $\rho: M \backslash B \rightarrow \mathbb{R}_{>0}$ by

$$
\begin{equation*}
\rho(x)=d(\arg \circ \pi)\left(R_{\alpha}(x)\right) . \tag{6.1.1}
\end{equation*}
$$

Roughly speaking, $\rho(x)$ measures infinitesimally how fast the orbit of $R_{\alpha}$ through $x$ traverses the pages.

Definition 6.1.2. The infinitesimal variation on the page $S_{\theta}$ is

$$
\begin{equation*}
V_{\theta}:=\sup _{x \in S_{\theta}} \rho(x) / \inf _{x \in S_{\theta}} \rho(x) \in[1, \infty), \tag{6.1.2}
\end{equation*}
$$

and the total infinitesimal variation is $V:=\sup _{\theta \in[0,2 \pi]} V_{\theta}$.
The following is standard:

## Lemma 6.1.3. If $V \equiv 1$, then $M \backslash \bar{S}_{0}$ is the interior of a contact handlebody.

Proof. Let $t$ be the coordinate obtained by flowing in the direction of $R_{\alpha}$ starting from $S_{0}$. Then $M \backslash \bar{S}_{0} \simeq(0, C)_{t} \times S_{0}$ and $\alpha=f_{t} d t+\beta_{t}$, where $f_{t}\left(\right.$ resp. $\left.\beta_{t}\right)$ is a function (resp. 1-form) on $S_{0}$ that depends on $t$.

We claim that $R_{\alpha}=\partial_{t}$ implies that $f_{t}=1$ and $\dot{\beta}_{t}=0$, where the dot denotes the derivative in the $t$-direction: Since $\alpha\left(R_{\alpha}\right)=1$, we have $f_{t}=1$. Then $d \alpha$ becomes $d t \wedge \dot{\beta}_{t}+d_{S_{0}} \beta_{t}$, where $d_{S_{0}}$ is the exterior derivative in the $S_{0}$-direction. Finally, $i_{R_{\alpha}} d \alpha=0$ forces $\dot{\beta}_{t}=0$.
6.2. Quantitative stabilization for $S^{2 n-1}$. Before starting, we warn the reader that the type of stabilization in this section is different from the notion of stabilization of an OBD in the Giroux correspondence which involves changing the topology of the page by a handle attachment and composing the monodromy with a suitable Dehn twist.

Let $\xi_{s t d}$ be the standard contact structure on $S^{2 n-1}=\left\{\left|z_{1}\right|^{2}+\cdots+\left|z_{n}\right|^{2}=\right.$ $1\} \subset \mathbb{C}^{n}$ given by the restriction of $\alpha=\frac{1}{2} \sum_{i=1}^{n}\left(x_{i} d y_{i}-y_{i} d x_{i}\right)$ which we denote by $\alpha_{s t d}$. The standard OBD for $\xi_{s t d}$ can be constructed as follows: Starting with $z_{1}: S^{2 n-1} \rightarrow \mathbb{C}$ which is a submersion away from $\left|z_{1}\right|=1$, we set $B=z_{1}^{-1}(0)=$ $S^{2 n-3}=\left\{\left|z_{2}\right|^{2}+\cdots+\left|z_{n}\right|^{2}=1\right\}$ and

$$
\pi=\frac{z_{1}}{\left|z_{1}\right|}: S^{2 n-1} \backslash B \rightarrow S^{1} \subset \mathbb{C}
$$

The pages $S_{\theta}$ are Weinstein $(2 n-2)$-disks and the Reeb vector field is

$$
R_{\alpha_{s t d}}=\sum_{i=1}^{n}\left(x_{i} \partial_{y_{i}}-y_{i} \partial_{x_{i}}\right)=\sum_{i=1}^{n} \partial_{\theta_{i}}
$$

where $\theta_{i}$ is the $i$ th angular coordinate. Then $\rho\left(z_{1}, \ldots, z_{n}\right)=d \theta_{1}\left(\partial_{\theta_{1}}\right)=1, V \equiv 1$, and $S^{2 n-1} \backslash \bar{S}_{0}$ is the interior of a genuine contact handlebody of thickness $2 \pi$ by Lemma 6.1.3. Note that $B$ has an analogous OBD derived from $z_{2}: S^{2 n-3}=$ $\left\{\left|z_{2}\right|^{2}+\cdots+\left|z_{n}\right|^{2}=1\right\} \rightarrow \mathbb{C}$.

Lemma 6.2.1. For any positive integer $k>0$ and $\epsilon^{\prime}>0$ small, there exists an $\alpha_{\text {std-compatible, strongly Weinstein } O B D}\left(B_{k}, \pi_{k}\right)$ of $S^{2 n-1}$ such that each page is $C^{\infty}$-close to

$$
S_{\theta_{0}} \cup S_{\theta_{0}+2 \pi / k} \cup \cdots \cup S_{\theta_{0}+2 \pi(k-1) / k}
$$

(i.e., the union of $k$ evenly spaced pages for some $\theta_{0}$ ) outside an $\epsilon^{\prime}$-small neighborhood of $B$ (with respect to the standard Euclidean metric on $\mathbb{C}^{n}$ ) and such that $S^{2 n-1}$, cut open along a new page, is the interior of a contact handlebody of thickness $2 \pi / k$.

Proof. We would like to "stabilize" $(B, \pi)$ by replacing $z_{1}$ by $z_{1}^{k}$. However, since 0 is not a regular value of $z_{1}^{k}$, we use

$$
f_{k}:=z_{1}^{k}+\epsilon z_{2}^{k}+\epsilon^{2} z_{3}^{k}+\cdots+\epsilon^{n-1} z_{n}^{k}
$$

where $\epsilon>0$ is small. We are thinking of $f_{k}$ as inductively defined as $z_{1}^{k}$ plus $\epsilon$ times $f_{k}$ corresponding to the binding $\left|z_{2}\right|^{2}+\cdots+\left|z_{n}\right|^{2}=1$. We write $R=R_{\alpha_{s t d}}$.

Step 1. Verification that 0 is a regular value of $f_{k}$. Let $F_{k}$ be $f_{k}$ viewed as a $\operatorname{map} \mathbb{C}^{n} \rightarrow \mathbb{C}$. Then $d F_{k}\left(z_{1}, \ldots, z_{k}\right)=k\left(z_{1}^{k-1}, \epsilon z_{2}^{k-1}, \ldots, \epsilon^{n-1} z_{n}^{k-1}\right)$. Next we precompose with the derivative of the inclusion map $i: S^{2 n-1} \hookrightarrow \mathbb{C}^{n}$. Let $z=\left(z_{1}, \ldots, z_{n}\right) \in S^{2 n-1}$. Suppose there exists $z_{j} \neq 0,1$. Then there exists $v \in T_{z} S^{2 n-1}$ with a nontrivial component in the $z_{j}$-direction and $d f_{k}$ is surjective since $\epsilon^{j-1} z_{j}^{k-1} \neq 0$. Otherwise, some $z_{j}=1$ and $z_{i}=0$ for all $i \neq j$. Then $f_{k}(z)=\epsilon^{j-1} z_{j}^{k} \neq 0$. Hence 0 is a regular value of $f_{k}$.

We set $B_{k}=f_{k}^{-1}(0) \pi_{k}=\frac{f_{k}}{\left|f_{k}\right|}: S^{2 n-1} \backslash B_{k} \rightarrow S^{1}$, and $S_{k, \theta}=\pi_{k}^{-1}(\theta)$.

Step 2. Computation of $d f_{k}(R)$. For $\left(z_{1}, \ldots, z_{k}\right) \in S^{2 n-1} \backslash B_{k}$, we use polar coordinates $\left(r_{i}, \theta_{i}\right)$ for $z_{i}$ and compute:

$$
\begin{align*}
d f_{k}(R) & =d\left(r_{1}^{k} e^{i k \theta_{1}}+\epsilon r_{2}^{k} e^{i k \theta_{2}}+\cdots+\epsilon^{n-1} r_{n}^{k} e^{i k \theta_{n}}\right)\left(\sum_{j=1}^{n} \partial_{\theta_{j}}\right) \\
& =i k\left(r_{1}^{k} e^{i k \theta_{1}}+\epsilon r_{2}^{k} e^{i k \theta_{2}}+\cdots+\epsilon^{n-1} r_{n}^{k} e^{i k \theta_{n}}\right)=i k f_{k} . \tag{6.2.1}
\end{align*}
$$

Observe that this equation is the version of the equation in [Gir00, p.411, last line] when $F_{k}$ is holomorphic.

Step 3. Verification of the properties. Eq. (6.2.1) implies that $R$ is tangent to $B_{k}$ and transverse to $S_{k, \theta}$ and that $\pi_{k}$ is a fibration. Moreover, for $\left(B_{k}, \pi_{k}\right), \rho(z)=1$ and $V \equiv 1$, and $S^{2 n-1} \backslash \bar{S}_{k, 0}$ is the interior of a contact handlebody of thickness $2 \pi / k$. Since $R$ is tangent to $B_{k}$, for each $z \in B_{k}, d f_{k}\left(\operatorname{ker} \alpha_{s t d}(z)\right)=\mathbb{C}$. This, together with the invariance of ker $\alpha_{\text {std }}$ under the standard almost complex structure on $\mathbb{C}^{n}$, implies that $B_{k}$ is a codimension two contact submanifold of $\operatorname{ker} \alpha_{s t d}$.

Next we apply Lemma 6.3.1, proved in $\S 6.3$, to show that $S_{k, \theta}$ is Weinstein after perturbing $f_{k}$ by adding $\sum_{i=1}^{n} c_{i} z_{i}^{k}$ for $c_{i}$ small. Using the notation from $\S 6.3$, we have $d \theta(R)=1$ and $d \phi(R)=0$; hence $d \phi\left(X_{\beta}\right)=0$ if and only if $d \phi=0$, where $\phi$ is viewed as a function on $S_{k, \theta=0}$. One can compute that if $z \in S_{k, 0}$ is a critical point of $\phi$, then all $\left|z_{i}\right| \neq 0$; at such a point $d\left(z_{1}^{k}\right), \ldots, d\left(z_{n}^{k}\right)$ are linearly independent. This provides enough perturbations to make $\phi$ Morse.

Finally, the $C^{\infty}$-closeness property is immediate from taking $\epsilon>0$ small.
6.3. Verification of the Weinstein property. Let $f: \mathbb{C}^{n} \rightarrow \mathbb{C}$ be a holomorphic function such that $f(0)=0$ with an isolated critical point at the origin. For $\kappa>0$ sufficiently small, $f$ defines an $\operatorname{OBD}(B, \theta)$ of the sphere $S_{\kappa}$ of radius $\kappa$, where $B=S_{\kappa} \cap f^{-1}(0)$ and $\theta=\arg f: S_{\kappa} \backslash B \rightarrow \mathbb{R} / 2 \pi \mathbb{Z}$. Also let $\phi:=-\log |f|$ on $S_{\kappa} \backslash B$.

Let

- $\alpha=\frac{1}{2} \sum_{i=1}^{n}\left(x_{i} d y_{i}-y_{i} d x_{i}\right)=\frac{1}{2} \sum_{i=1}^{n} r_{i}^{2} d \theta_{i}$ be the standard Liouville form on $\mathbb{C}^{n}$ with Liouville vector field $X_{\alpha}=\frac{1}{2} \sum_{i=1}^{n} r_{i} \frac{\partial}{\partial r_{i}}$, where $\left(r_{i}, \theta_{i}\right)$ are polar coordinates corresponding to $\left(x_{i}, y_{i}\right)$;
- $\alpha_{\text {std }}$ be the induced contact form on $S_{\kappa}$ for $\kappa>0$ small and $R=\frac{2}{r^{2}} \sum_{i} \frac{\partial}{\partial \theta_{i}}$ be the Reeb vector field on $S_{\kappa}$, where $r^{2}=\sum_{i} r_{i}^{2}$ (note that $R$ is defined on all of $\mathbb{C}^{n}$, not just on $S_{\kappa}$ );
- $\beta$ be the 1 -form induced by $\alpha_{s t d}$ on any page of the OBD with Liouville vector field $X_{\beta}$; we view $X_{\beta}$ as a vector field on $S_{\kappa} \backslash B$ that is tangent to the pages;
- $X_{\theta}$ be the Hamiltonian vector field of $\theta$ (satisfying $i_{X_{\theta}} d \alpha=d \theta$ ), viewed as a function on $\mathbb{C}^{n} \backslash f^{-1}(0)$; and
- $J$ be the standard complex structure on $\mathbb{C}^{n}$.

The following is due to Emmanuel Giroux (presented here with his permission):
Lemma 6.3.1 (Giroux). On $S_{\kappa} \backslash B$, if $d \phi(R) \equiv 0$, then $d \phi\left(X_{\beta}\right)=\frac{1}{d \theta(R)}|d \phi|^{2}$.

Proof. We first claim that, at every point of $S_{\kappa} \backslash B$, the following identity holds:

$$
\begin{equation*}
X_{\alpha}=X_{\beta}+a X_{\theta}+b R, \tag{6.3.1}
\end{equation*}
$$

where $a=\frac{1}{d \theta(R)}$ and $b=\frac{d \theta\left(X_{\alpha}\right)}{d \theta(R)}$. First note that the $d \alpha$-symplectic orthogonal complement of the tangent space $T S$ of a page $S$ is spanned by $R$ and $X_{\theta}$. [Verification: $d \alpha(R, T S)=0$ since $R, T S$ are tangent to $S_{\kappa} ; d \alpha\left(X_{\theta}, T S\right)=d \theta(T S)=$ 0 ; and $d \alpha\left(X_{\theta}, R\right)=d \theta(R)>0$.] Hence we can write $X_{\alpha}=Y+a X_{\theta}+b R$ for some $Y \in T S$. Evaluating on $T S$ gives $Y=X_{\beta}$. We can then determine $a$ and $b$ by applying $d \theta$ and $d\left(r^{2}\right)$ to Eq. (6.3.1): $d\left(r^{2}\right)\left(X_{\alpha}\right)=2 r d r\left(\frac{1}{2} \sum_{i} r_{i} \frac{\partial}{\partial r_{i}}\right)=r^{2}$ and $d\left(r^{2}\right)\left(X_{\alpha}\right)=a d\left(r^{2}\right)\left(X_{\theta}\right)=-a d \theta\left(X_{r^{2}}\right)=-a d \theta\left(-r^{2} R\right)$, so $a=\frac{1}{d \theta(R)}$. Similarly, $d \theta\left(X_{\alpha}\right)=b d \theta(R)$, so $b=\frac{d \theta\left(X_{\alpha}\right)}{d \theta(R)}$.

It follows from Eq. (6.3.1) that

$$
d \phi\left(X_{\beta}\right)=-\frac{d \phi\left(X_{\theta}\right)}{d \theta(R)}+d \phi\left(X_{\alpha}\right)-d \phi(R) \frac{d \theta\left(X_{\alpha}\right)}{d \theta(R)} .
$$

Now, since the function $-\phi+i \theta$ is holomorphic when viewed as a function on $\mathbb{C}^{n} \backslash\{f=0\}$, we have $d \phi \circ J=d \theta$. Hence

$$
\begin{aligned}
d \phi\left(X_{\alpha}\right) & =-d \phi \circ J\left(J X_{\alpha}\right)=-d \theta\left(\frac{1}{2} \sum_{i} \frac{\partial}{\partial \theta_{i}}\right)=-\frac{r^{2}}{4} d \theta(R), \\
d \phi\left(X_{\theta}\right) & =-d \phi \circ J\left(J X_{\theta}\right)=-d \theta\left(J X_{\theta}\right)=-d \alpha\left(X_{\theta}, J X_{\theta}\right) \\
& =-\left|X_{\theta}\right|^{2}=-|d \theta|^{2}=-|d \phi|^{2}, \\
d \phi(R) & =-d \theta(J R)=d \theta\left(\frac{4}{r^{2}} X_{\alpha}\right),
\end{aligned}
$$

and

$$
d \phi\left(X_{\beta}\right)=\frac{1}{d \theta(R)}\left(|d \phi|^{2}-\left(d \phi\left(\frac{r}{2} R\right)\right)^{2}-\left(d \phi\left(\frac{2}{r} X_{\alpha}\right)\right)^{2}\right) .
$$

Observe that $\frac{2}{r} X_{\alpha}$ and $\frac{r}{2} R$ are orthonormal unit vectors. Finally, since $d \phi(R)=0$, $d \phi\left(X_{\beta}\right)=\frac{1}{d \theta(R)}|d \phi|^{2}$ on each page.

## 7. Construction of the plug

The goal of this section is to generalize the 3-dimensional plug constructed in §4.1 to higher dimensions. This is the key construction that will allow us to prove Theorems 1.2.3 and 1.2 .5 in Section 9 in essentially the same way as in the 3dimensional case.
7.1. Definition of a plug. Let us rephrase the 3-dimensional case considered in $\S 4.1$ in a way that is amenable to higher-dimensional generalization. Consider the standard contact space $\left(\mathbb{R}^{3}, \operatorname{ker}\left(d z+e^{s} d t\right)\right)$ and the surface $\Sigma=\{z=0\} \subset \mathbb{R}^{3}$. The plug is obtained by growing a mushroom along a box $U=\left[0, s_{0}\right] \times\left[0, t_{0}\right] \subset \Sigma$, where we are viewing $U$ as the truncated symplectization of the 1 -dimensional compact contact manifold $\partial_{-} U=\{0\} \times\left[0, t_{0}\right]$ with contact form $d t$.

In higher dimensions, let $(Y, \operatorname{ker} \eta)$ be a compact contact manifold of dimension $2 n-1$ with convex boundary. Let

$$
\left(N_{\epsilon_{0}}(Y):=Y \cup\left(\left[0, \epsilon_{0}\right] \times \partial Y\right), \text { ker } \eta\right)
$$

be a small extension of $(Y, \operatorname{ker} \eta)$. Now we consider

$$
\left(M^{2 n+1}:=\mathbb{R}_{z, s}^{2} \times N_{\epsilon_{0}}(Y), \xi=\operatorname{ker}\left(d z+e^{s} \eta\right)\right)
$$

and the hypersurface $\Sigma:=\{z=0\}$. Let $U:=\left[0, s_{0}\right] \times N_{\epsilon_{0}}(Y)$ and let

$$
\partial_{-} U:=\{-1\} \times N_{\epsilon_{0}}(Y) \quad \text { and } \quad \partial_{+} U:=\left\{s_{0}+1\right\} \times N_{\epsilon_{0}}(Y) .
$$

From now on, we fix a Riemannian metric on $M$, which induces a metric on any submanifold and such that $\left[0, \epsilon_{0}\right] \times \partial Y$ has thickness $\epsilon_{0}$ with respect to this metric.
Definition 7.1.1. $A Y$-shaped plug with parameter $\epsilon>0$ is a $C^{0}$-small perturbation $\widetilde{U}$ of $U$ supported in the interior $U^{\circ}$ of $U$ such that:
(1) all the flow lines of $\widetilde{U}_{\xi}$ that pass through $\{-1\} \times Y^{\circ}$ flow to a negative singularity;
(2) all the flow lines of $\widetilde{U}_{\xi}$ that pass through $\left\{s_{0}+1\right\} \times Y^{\circ}$ flow from a positive singularity;
(3) for all possibly broken flow lines of $\widetilde{U}_{\xi}$ that go from $\partial_{-} U$ to $\partial_{+} U$, the holonomy map is $\epsilon$-close to the identity when defined;
(4) $\widetilde{U}_{\xi}$ is gradient-like with respect to a Morse function $f: \widetilde{U} \rightarrow \mathbb{R}$ which agrees with s on $\partial \widetilde{U}$; in particular there are no possibly broken loops of $\widetilde{U}_{\xi}$.
Definition 7.1.2. $A Y$-shaped pre-plug $\widetilde{U}$ satisfies Definition 7.1.1 with (3) replaced by:
(3') for each possibly broken flow line of $\widetilde{U}_{\xi}$ that goes from $\partial_{-} U$ to $\partial_{+} U$, the holonomy map is obtained by following a small perturbation of $(\partial Y)_{\operatorname{ker}} \eta$.
7.2. A Peter-Paul contactomorphism. Let $(Y, \eta)$ be a contact manifold with a fixed choice of contact form $\eta$. Let $S$ be a hypersurface of $Y$ transverse to the Reeb vector field $R_{\eta}$. Then $S$ has a neighborhood $S \times[-\epsilon, \epsilon]_{\tau} \subset Y$ on which $R_{\eta}=\partial_{\tau}$.

The following is well-known:
Lemma 7.2.1. If $R_{\eta}=\partial_{\tau}$ on $S \times[a, b]_{\tau} \subset Y, a<b$, then $\eta=d \tau+\beta$, where $\beta$ is the pullback of a 1 -form on $S$. Moreover, $d \beta$ is symplectic on $S$.

In other words, $\eta$ is the contactization of $(S, \beta)$. In particular, if $(S, \beta)$ is Weinstein then $S \times[a, b]$ is a contact handlebody.
Proof. We first write $\eta=f d \tau+\beta$, where $f(\tau) \in \Omega^{0}(S)$ and $\beta(\tau) \in \Omega^{1}(S)$. Since $\eta\left(R_{\eta}\right)=1$, we have $f=1$. Also, since $\mathscr{L}_{R_{\eta}} \eta=0, \beta(\tau)$ must be $\tau$-independent. Finally, $d \beta$ is symplectic on $S$ due to the contact condition on $Y$.

Given $(Y, \eta)$, let $(M, \alpha)=\left(\mathbb{R}_{z, s}^{2} \times Y, d z+e^{s} \eta\right)$ and let $\phi_{t}: Y \xrightarrow{\sim} Y$ be the time- $t$ flow of $R_{\eta}$.
Lemma 7.2.2. The diffeomorphism

$$
\begin{equation*}
\Psi: M \xrightarrow{\sim} M, \quad(z, s, y) \mapsto\left(e^{(-1+1 / C) s} \cdot C z, s / C, \phi_{(1-C) e^{-s} z}(y)\right), \tag{7.2.1}
\end{equation*}
$$

where $C>0$, is a contactomorphism.

Proof. We compute

$$
\begin{aligned}
\Psi^{*}(\alpha)= & d\left(e^{(-1+1 / C) s} \cdot C z\right)+e^{s / C}\left(\eta+d\left((1-C) e^{-s} z\right)\right) \\
= & e^{(-1+1 / C) s}(1-C) z d s+e^{(-1+1 / C) s} C d z+e^{s / C} \eta \\
& -e^{s / C}(1-C) e^{-s} z d s+e^{s / C}(1-C) e^{-s} d z \\
= & e^{(-1+1 / C) s}\left(d z+e^{s} \eta\right)=e^{(-1+1 / C) s} \alpha .
\end{aligned}
$$

We explain the first line: By Lemma 7.2.1, $\eta$ can locally be written as $d \tau+\beta$, where $\beta$ is a 1-form on a hypersurface $S \subset Y$ transverse to $R_{\eta}=\partial_{\tau}$. (Note that we can use the immersion $i: S \times \mathbb{R}_{\tau} \rightarrow Y$ with $i^{*} R_{\eta}=\partial_{\tau}$ instead in Lemma 7.2.1.) Then $\phi_{(1-C) e^{-s} z}(\tau, x)=\left(\tau+(1-C) e^{-s} z, x\right)$, where $x$ is the coordinate on $S$, and $\phi_{(1-C) e^{-s} z}^{*} \eta=\eta+d\left((1-C) e^{-s} z\right)$.

Let $M_{\left(z_{0}, s_{0}\right)}=\left[-z_{0}, z_{0}\right] \times\left[0, s_{0}\right] \times Y \subset(M, \alpha)$. As an immediate corollary, by taking $C \gg 0$, we have:

Lemma 7.2.3. For any $0<s_{0}^{\prime} \leq s_{0}$ and $z_{0}^{\prime}>0$, there exists $0<z_{0} \leq z_{0}^{\prime}$, such that $M_{\left(z_{0}, s_{0}\right)}$ contactly embeds into $M_{\left(z_{0}^{\prime}, s_{0}^{\prime}\right)}$ and takes $\{z=0\} \cap M_{\left(z_{0}, s_{0}\right)}$ to $\{z=0\} \cap M_{\left(z_{0}^{\prime}, s_{0}^{\prime}\right)}$.

We call the contactomorphism $\Psi$ given by (7.2.1) a Peter-Paul contactomorphism for the following reason: In Lemma 7.2.3, $\Sigma=\{0\} \times[0, s] \times Y \subset M_{\left(z_{0}, s_{0}\right)}$ is the hypersurface on which we want to create mushrooms. The length of the interval $\left[-z_{0}, z_{0}\right]$ can be regarded as the given size of a neighborhood of $\Sigma$. The map $\Psi$ then allows us to rob the (already small) size of the neighborhood of $\Sigma$ to pay for a large size in the $s$-direction.

Observe that the Peter-Paul contactomorphism was not needed in Section 4 to make any 2 -dimensional surface convex.
7.3. A pre-plug. Given a standard Darboux ball $\left(Y^{2 n-1}, \eta^{\prime}\right)$ with convex boundary, we explain how to construct a $Y$-shaped pre-plug $\widetilde{U}$ on

$$
\Sigma:=\{z=0\} \subset M=\mathbb{R}_{z, s}^{2} \times N_{\epsilon_{0}}(Y) .
$$

The $Y$-shaped pre-plug will be upgraded to a $Y$-shaped plug with parameter $\epsilon>0$ in the next two subsections.

Modulo corner rounding, we may assume that ( $Y, \operatorname{ker} \eta^{\prime}$ ) is contactomorphic to ( $Y_{0} \cup Y_{1}$, ker $\eta$ ) such that:
(1) $Y_{0}=\left(z_{1}^{-1}\left(\left\{|\zeta| \leq \epsilon^{\prime}\right\}\right), \eta=\alpha_{s t d}\right)$, where $z_{1}: S^{2 n-1} \rightarrow \mathbb{C}_{\zeta}$ and $\alpha_{s t d}=$ $\frac{1}{2} \sum_{i}\left(x_{i} d y_{i}-y_{i} d x_{i}\right)$ are as in $\S 6.2$ and $\epsilon^{\prime}>0$ is small;
(2) $Y_{1}=\left([0,2 \pi]_{t} \times D^{2 n-2}, \eta=d t+\beta_{D^{2 n-2}}\right)$, where $\beta_{D^{2 n-2}}$ is a standard Liouville form on $D^{2 n-2}$ with one elliptic singular point; and
(3) for each $t \in[0,2 \pi],\{t\} \times \partial D^{2 n-2}$ is glued to $z_{1}^{-1}\left(\epsilon^{\prime} e^{i t}\right)$ so that the contact forms (and the Reeb vector fields) match.
(In particular, $Y_{0} \cup Y_{1}$ has a partial open book structure, where $Y_{0}$ is a neighborhood of the binding and $\{t\} \times D^{2 n-2}$ is a retraction of a page; see Definition 8.4.2.)

We now apply the Peter-Paul contactomorphism to realize $\left(Y_{0} \cup Y_{1}, \eta\right)$ as a transverse slice (i.e., transverse to the characteristic foliation) on $\Sigma$; we then blur the distinction between $Y$ and $Y_{0} \cup Y_{1}$ and write $Y=Y_{0} \cup Y_{1}$. The price that we pay is that we lose control of the $z$-height. (Recall that the $z$-height restricts the thickness of the contact handlebody of the mushroom that we want to grow.)

To remedy this we apply quantitative stabilization (Lemma 6.2.1) with $k \gg 0$ and $0<\epsilon \ll \epsilon^{\prime}$ to $Y_{0}$ to obtain $\pi_{k}: Y_{0} \rightarrow S^{1}$. After a $C^{\infty}$-small perturbation of $Y_{1}$ whose size depends on $k$ and which we still denote by $Y_{1}$, there is a smooth extension $\pi_{k}: Y \rightarrow S^{1}$ which agrees with $(t, x) \mapsto k t$ on $Y_{1}$.

Choose a finite number, say $N=5$, and a small constant $0<\epsilon^{\prime \prime}<\epsilon^{\prime}$. Then, for $j=0, \ldots, 4$, let $H_{j}^{\prime}$ be the sector $\pi_{k}^{-1}\left(\left[\frac{2 \pi j}{5}, \frac{2 \pi(j+1)}{5}\right]\right)$. Assuming we chose a suitable extension of $\pi_{k}, H_{j}^{\prime}$ is a (complete) contact handlebody of thickness $\frac{2 \pi}{k}$. Next let $H_{j}$ be a slight modification of $H_{j}^{\prime}, j=0, \ldots, 4$, obtained by removing an $\frac{\epsilon^{\prime \prime}}{2}$-neighborhood $\pi_{k}^{-1}\left(\left\{z \leq \frac{\epsilon^{\prime \prime}}{2}\right\}\right) \cap Y_{0}$ of the binding $B_{k}$ and thickening the contact handlebody by flowing forward and backward by $\frac{\epsilon^{\prime \prime}}{k}$ in the Reeb direction.

Finally, we slightly (i.e., in a $C^{\infty}$-small manner) modify $\eta$ by "shifting the binding" $B_{k}$ away from $H_{0}$ using a contactomorphism $\phi$ and construct the analogous contact handlebody $\phi\left(H_{0}\right)$ of thickness $\frac{2 \pi+2 \epsilon^{\prime \prime}}{k}$ such that $\phi\left(H_{0}\right)$ contains the $\epsilon^{\prime \prime}$ neighborhood of $B_{k}$, as follows:

Brief explanation of "shifting the binding". By flowing along the characteristic foliation of each page, one can normalize $\alpha_{s t d}$ on a small neighborhood $B_{k} \times D^{2}$ of the binding $B_{k}$ as

$$
\alpha_{s t d}=f \beta+g(x d y-y d x)
$$

where $D^{2}$ is a disk with Euclidean/polar coordinates $(x, y) /(r, \theta)$ and small radius; $\beta=\left.\alpha_{s t d}\right|_{B_{k}} ; f=f(r)$ and $f(0)=1 ; g=g_{0} r^{2}+O\left(r^{3}\right)$ and $g_{0}$ is a nonvanishing function on $B_{k}$; and the pages are still $\theta=$ const. Letting

$$
\alpha:=f^{-1} \alpha_{s t d}=\beta+h(x d y-y d x), \quad \text { where } \quad h=f^{-1} g
$$

and $X=\frac{\partial}{\partial x}-h y R_{\beta}$, where $R_{\beta}$ is the Reeb vector field of $\beta$, we compute:

$$
\begin{gathered}
d \alpha=d \beta+d h \wedge(x d y-y d x)+2 h d x d y \\
i_{X} d \alpha=2 h d y+O(r), \quad d(\alpha(X))=-2 h d y+O(r)
\end{gathered}
$$

and hence $\mathscr{L}_{X} \alpha=O(r)$. Therefore, the translation by $a X, a<0$ small, on $B_{k} \times D^{2}$ (where defined), is close to a contactomorphism and the modification needed to make it into a contactomorphism $\phi$ is on the order of $a \cdot O(r)$, which is an order of magnitude smaller. Hence $\phi\left(H_{0}\right)$ has the property of being close to the $x \mapsto x+a$-translate of $H_{0}$ with error much smaller than $a$; in particular $\phi\left(H_{0}\right)$ contains the $\epsilon^{\prime \prime}$-neighborhood of $B_{k}$. Finally, one may adjust the contact form on $\phi\left(H_{0}\right)$ by multiplying by a function that is close to 1 so that $\phi\left(H_{0}\right)$ becomes a contact handlebody.

In what follows we abuse notation and refer to $\phi\left(H_{0}\right)$ by $H_{0}$.

The pre-plug $\widetilde{U}$ consists of mushrooms $Z_{j}, 0 \leq j \leq 4$, with bases

$$
B_{j}:=\left[\frac{8-2 j}{10}, \frac{9-2 j}{10}\right]_{s} \times H_{j}
$$

see Figure 7.3.1. (Strictly speaking, the base $B_{0}$ is obtained by shifting $B_{0}^{\prime}=$ $\left[\frac{8}{10}, \frac{9}{10}\right] \times H_{0}$ by the map $(s, x) \mapsto\left(s+f_{0}(x), x\right)$, where $f_{0}$ is a $C^{0}$-small smooth function on $H_{0}$. We will assume that we started with a slightly smaller $B_{0}^{\prime}=$ $\left[\frac{8}{10}+\delta, \frac{9}{10}-\delta\right] \times H_{0}$ and that the base for the contact handlebody $H_{0}$ contains $\left.B_{0}.\right)$ We also require the damping of $Z_{1}, \ldots, Z_{4}$ to occur inside the $\frac{3 \epsilon^{\prime \prime}}{4}$-neighborhood of $B_{k}$ so that the damping regions of $H_{1}, \ldots, H_{4}$ are contained in $H_{0}$ and the damping region of $H_{0}$ to occur inside $\mathrm{H}_{2}$.


Figure 7.3.1. $H_{0}, \ldots, H_{4}$, from top to bottom.

Verification of (3') in Definition 7.1.2. This is a direct consequence of Proposition 5.5.2 and our choice of the ordering of the mushrooms in the $s$-direction.

We first give names to regions of $\partial_{-} U$ : Viewing $H_{j}$ as a subset of $\partial_{-} U$, let $\widetilde{H}_{j}$ be the closure of the union of $H_{j}$ and the set of points such that the holonomy from $s=\frac{8-2 j}{10}-\delta$ to $\frac{9-2 j}{10}+\delta($ for $\delta>0$ small) is not trivial or does not exist; note that $\widetilde{H}_{j}$ is contained in a small neighborhood of $H_{j}$. We denote the portion of $\widetilde{H}_{j}$ that closely approximates it and acts as a sink by $H_{j, \text { in }}$ and $\widetilde{H}_{j} \backslash H_{j, \text { in }}$ by $H_{j, \partial}$, and the corresponding products with $\left[\frac{8-2 j}{10}, \frac{9-2 j}{10}\right]$ by $B_{j, \text { in }}$ and $B_{j, \partial}$.

The dynamics of $\widetilde{U}$ in forward time is described as follows: Let $x \in \partial_{-} U$ and let $\ell_{x}$ be the flow line of $\widetilde{U}$ passing through $x$.
(A) If $\ell_{x}$ enters $B_{j, \text { in }}, j=0, \ldots, 4$, then $\ell_{x}$ converges to a singularity in $Z_{j}$.
(B) If $\ell_{x}$ enters $B_{j, \partial}$, then $\ell_{x}$ exits $Z_{j}$ at a point near $\partial H_{j}$ - recall that as $\ell_{x}$ passes through $Z_{j}$ it flows "from $R_{+}\left(\partial H_{j}\right)$ to $R_{-}\left(\partial H_{j}\right)$ " in the sense of $(Z 5)$ of Proposition 5.5.2 with possibly large nontrivial components in the Reeb direction of $\Gamma_{\partial H_{j}}$ on the damping region - and one of the following will happen:
(1) $\ell_{x}$ follows $\partial_{s}$ until $\{s=1\}$ and exits the plug; the holonomy map is obtained by following a small perturbation of $(\partial Y)_{\operatorname{ker} \eta}$;
(2) $\ell_{x}$ follows $\partial_{s}$ until $B_{i, \text { in }}$ for $i<j$; we then apply (A) after letting the new $j$ equal $i$;
(3) $\ell_{x}$ follows $\partial_{s}$ until $B_{i, \partial}$ for $i<j$; we then apply (B) after letting the new $j$ equal $i$.
Note that $H_{0}$ was chosen so that $Z_{0}$ captures all the trajectories that enter $\partial_{-} U$, "survives to" $s=\frac{8}{10}$, and is not close to $\partial Y$. A similar analysis can be applied to the dynamics of $\widetilde{U}$ in backward time and Definition 7.1.2(3') then holds.

In light of Proposition 2.2.6 and Lemma 2.3.2, we conclude that $\widetilde{U}$ is Morse.
7.4. $\epsilon$-short hypersurfaces. In this subsection we strengthen Theorem 1.2.3 in a quantitative way. Namely, in addition to requiring that $\Sigma_{\xi}$ be Morse $^{+}$or 1-Morse ${ }^{+}$, we also require all the smooth flow lines of $\Sigma_{\xi}$ to be short.

Definition 7.4.1. A closed hypersurface $\Sigma \subset(M, \xi)$ is $\epsilon$-short if the length of any smooth flow line of $\Sigma_{\xi}$ is shorter than $\epsilon$ with respect to the induced metric on $\Sigma$.

Observe that any closed Morse ${ }^{+}$hypersurface $\Sigma$ is $\epsilon$-short for a sufficiently large $\epsilon$ which depends on $\Sigma$. For our purposes we take $\epsilon>0$ to be a small number which is independent of the choice of convex hypersurface. Theorem 1.2.3 can be strengthened as follows:

Theorem 7.4.2. Given $\epsilon>0$, any closed hypersurface $\Sigma$ in a contact manifold can be $C^{0}$-approximated by an $\epsilon$-short and Morse ${ }^{+}$one. Moreover, if $\Sigma$ is Weinstein convex, then there exists a t-invariant neighborhood $[-\delta, \delta]_{t} \times \Sigma$ of $\{0\} \times \Sigma$ and a 1-parameter family of pairwise disjoint embeddings $\phi_{t}: \Sigma \rightarrow[-\delta, \delta] \times \Sigma$, $t \in[-\delta, 0]$, such that:
(1) $\phi_{-\delta}(\Sigma)=\{-\delta\} \times \Sigma$,
(2) $\phi_{t}(\Sigma)$ is $C^{0}$-close to $\{t\} \times \Sigma$ for all $t \in[-\delta, 0]$,
(3) $\phi_{0}(\Sigma)$ is 1-Morse ${ }^{+}$and $\epsilon$-short, and
(4) $\phi_{t}(\Sigma)$ have the same number and type of singular points for all $t \in[-\delta, 0)$ and the whole family $\phi_{t}(\Sigma), t \in[-\delta, 0]$ is Weinstein convex.

In words, what (3) and (4) are saying is that the singular points of the characteristic foliation of $\phi_{t}(\Sigma)$ remain the same for $t<0$ and all the singular points necessarily of birth-death type - are created at the same time when $t$ reaches 0 .

Theorem 7.4.2 holds in dimension 3 by $\S 4.3$, together with a slightly more careful analysis of the characteristic foliation when installing a mushroom. This will be the base case of our inductive argument.
7.5. Construction of the plug. The goal of this subsection is to prove the following:

Theorem 7.5.1. Let $Y$ be the standard Darboux ball of dimension $2 n-1$ with convex boundary. Then for any $\epsilon>0$ small there exists a $Y$-shaped plug with parameter $\epsilon$, provided Theorem 7.4.2 holds for contact manifolds of any dimension $\leq 2 n-1$.

We adopt a simplification due to Eliashberg, Fauteux-Chapleau, and Pancholi, used with their permission, while retaining certain elements of the original version of the paper.

Proof. Given $\epsilon>0$ small, choose $0<\epsilon_{0} \ll \epsilon$. Let $\Sigma$ be the pushoff of $\partial Y$ towards $\partial N_{\epsilon_{0}}(Y)$ such that $\Sigma$ is the standard Weinstein convex sphere, which we take to satisfy:
(Sph1) $\Sigma \simeq S^{2 n-2}$ and $\Gamma(\Sigma) \simeq S^{2 n-3}$, where $\simeq$ means "diffeomorphic";
(Sph2) $R_{ \pm}(\Sigma) \simeq D^{2 n-2}$ is a Weinstein domain with precisely one singular point, and the singular point on $R_{+}(\Sigma)$ (resp. $R_{-}(\Sigma)$ ) has index 0 (resp. index $2 n-2$ ).

Applying Theorem 7.4.2, there exists a foliation by $\Sigma_{t}^{\prime}:=\phi_{t}(\Sigma)$ for $t \in[-\delta, 0]$ satisfying (1)-(4). We are starting with $\Sigma_{-\delta}^{\prime}$ which is the standard Weinstein convex sphere by (1); for $t \in[-\delta, 0), \Sigma_{t}^{\prime}$ remains the standard Weinstein convex sphere by (4); and for $t=0, \Sigma_{0}^{\prime}$ is $\epsilon$-short and Weinstein convex with birth-death singular points by (3).

Setup for $\widetilde{U}$. The $Y$-shaped plug $\widetilde{U}$ consists of two mushrooms $Z_{0}, Z_{2}$ with contact handlebody profiles $H_{0}, H_{2}$; a pre-plug $Z_{1}$ with profile $H_{1}$; and bases $\left[\frac{2 j}{6} s_{0}, \frac{2 j+1}{6} s_{0}\right] \times H_{j}$ (i.e., $Z_{0}$ has the smallest $s$-value, followed by $Z_{1}$, and then by $Z_{2}$ ), such that the following hold:
(P1) $N_{\epsilon_{0}}(Y) \supset H_{0} \cup H_{1} \cup H_{2} \supset Y$ and $H_{0} \cap H_{2}=\varnothing$;
(P2) the damping region of $H_{0}$ is a subset of $H_{1, \text { in }}$ and the damping region of $H_{2}$ is a subset of $H_{1, \text { out }}$.

Note that $H_{1}$ does not need a damping region. See Figure 7.5.1. Recall that on the damping regions the characteristic foliation may have large components in the Reeb direction of the dividing set. Since these can potentially create trouble with the $\epsilon$-shortness, we require ( P 2 ).


Figure 7.5.1. A schematic diagram of the plug, with the $s$ direction projected out (the boundaries on the right-hand side do not actually exist). The arrows indicate the direction of the flow along the characteristic foliations of the handlebodies (i.e., from the positive side to the negative side). The circular region represents $N(\Gamma)$ and the shaded regions are the damping regions.

Analysis of the dynamics of $\widetilde{U}_{\xi}$. We will specify $H_{0}, H_{1}, H_{2}$ later, but for the moment we analyze the effect of the plug $\widetilde{U}$ when we stack the mushrooms as above.
(1) Given a flow line of $\tilde{U}_{\xi}$ that enters $Z_{0}$, either it flows to a negative singularity of $Z_{0}$ or passes through $Z_{0}$ after flowing parallel to the characteristic foliation of $\partial H_{0}$ as given by (Z5) of Proposition 5.5.2.
(2) Given a flow line of $\widetilde{U}_{\xi}$ that enters $Z_{1}$ (including one that exits $Z_{0}$ in the previous step), either it flows to a negative singularity of $Z_{1}$ or passes through $Z_{1}$ after flowing parallel to the characteristic foliation of $\partial H_{1}$. Notice that all the potentially problematic orbits that had large components in the Reeb direction in $Z_{0}$ (i.e., those that flowed out of the damping region of $H_{0}$ or near $R_{-}\left(\partial H_{0}\right)$ ) flow to a negative singularity of $Z_{1}$ by ( P 2 ).
(3) A similar consideration holds for a flow line of $\widetilde{U}_{\xi}$ that enters $Z_{2}$ (including one that exits $Z_{1}$ in the previous step).
(4) Backwards flow lines can be analyzed similarly.

The orbits that pass through $\tilde{U}$ flow along the characteristic foliations of $\partial H_{0}, \partial H_{1}$, and $\partial \mathrm{H}_{2}$ without ever entering damping regions; this means they flow closely along $\partial\left(H_{0} \cup H_{1} \cup H_{2}\right)$.

Description of $H_{0}, H_{1}, H_{2}$. Let us write $\Gamma=\Gamma\left(\Sigma_{0}^{\prime}\right)$ and let $N(\Gamma)=\Gamma \times D_{\rho, \theta}^{2}$ be a tubular neighborhood of $\Gamma$ with contact form $\beta_{\Gamma}+C\left(\frac{\rho^{2}}{2} d \theta\right)$ for $C>0$ small and such that the thickness in the $D^{2}$-direction is $\leq \epsilon^{\prime \prime}$ for $\epsilon^{\prime \prime}>0$ small. We take $H_{0}$ (resp. $H_{2}$ ) to be a thin contact handlebody over an $\epsilon^{\prime \prime \prime}$-retraction of $R_{+}\left(\Sigma_{0}^{\prime}\right)$ (resp. $R_{-}\left(\Sigma_{0}^{\prime}\right)$ ), where $0<\epsilon^{\prime \prime \prime} \ll \epsilon^{\prime \prime}$, and take the damping region to be small. There exists $\delta_{0}^{\prime}>0$ small such that

$$
R_{+}\left(\Sigma_{-\delta_{0}}^{\prime}\right) \backslash N(\Gamma) \subset H_{0, \text { out }} \quad \text { and } \quad R_{-}\left(\Sigma_{-\delta_{0}}^{\prime}\right) \backslash N(\Gamma) \subset H_{2, \text { in }}
$$

We then take $H_{1}$ to be union of $N(\Gamma)$ and the region of $N_{\epsilon_{0}}(Y)$ bounded by $\Sigma_{-\delta_{0}}^{\prime}$. There is some flexibility in choosing the dividing set and we take $\Gamma_{\partial H_{1}}$ and also the damping region to be a subset of $H_{2, \text { in }}$. (P1) and (P2) hold by construction and also $\partial\left(H_{0} \cup H_{1} \cup H_{2}\right)$ is $\epsilon$-short.

Lemma 7.5.2. $H_{1}$, after a $C^{\infty}$-small perturbation, is contactomorphic to a standard Darboux ball with convex boundary.

Proof of Lemma 7.5.2. Let $H_{1}^{\prime}:=[0,1]_{t} \times D$, where $D=D^{2 n-1}$, and let $D_{t}:=$ $\{t\} \times D$. We sketch the proof that if $\alpha$ is a contact form on $H_{1}^{\prime}$ such that:
(1) on a neighborhood of $[0,1] \times \partial D_{t}, \alpha=d t+\beta$, where $\beta$ is independent of $t$ and is the symplectization of a standard contact form on $\partial D$.
(2) the characteristic foliation on each $D_{t}$ consists of a positive elliptic singularity $e_{t}$ and all its trajectories go from $e_{t}$ to $\partial D_{t}$;
then $\left(H_{1}^{\prime}, \operatorname{ker} \alpha\right)$ is contactomorphic to $\left(H_{1}^{\prime}, \operatorname{ker}(d t+\beta)\right)$ for some extension of $\beta$ to $D$. The assumptions imply that $\alpha=f_{t} d t+\beta_{t}$, where $f_{t}$ is a function on $D_{t}$ and $\beta_{t}$ is a 1-form on $D_{t}$. We slightly perturb $\alpha$ so that the $\left(N\left(e_{t}\right),\left.\beta_{t}\right|_{N\left(e_{t}\right)}\right)$
are all diffeomorphic, where $N\left(e_{t}\right) \subset D_{t}$ is a neighborhood of $e_{t}$, and apply a 1parameter family of diffeomorphisms to straighten the characteristic foliation and make $\beta_{t}$ independent of $t$, which we now write as $\beta^{\prime}$. The contact condition implies that $f_{t}>0$; applying the Reeb flow gives the normalization $f d t+\beta^{\prime}$; and we divide by $f$.

The lemma then follows modulo adjusting/rounding corners.
In view of Lemma 7.5.2 and the construction of a pre-plug from $\S 7.3, \widetilde{U}$ is a $Y$-shaped plug with parameter $\epsilon$.

## 8. BIFURCATIONS OF CHARACTERISTIC FOLIATIONS AND BYPASS ATTACHMENTS

The goal of this section is to relate certain codimension 1 degenerations of Morse ${ }^{+}$hypersurfaces to bypass attachments introduced in [HH]. Such a correspondence is fundamental in bridging the more dynamical approach [Gir91, Gir00] and the more combinatorial approach [Hon00] of convex surface theory in dimension 3 . Unfortunately the details have never existed in the literature.

To this end, we slightly repackage the Morse theory on Morse hypersurfaces from Section 2 in terms of folded Weinstein hypersurfaces.
8.1. Definitions and examples. In this subsection we define folded Weinstein hypersurfaces and examine a few examples.

Definition 8.1.1. An oriented hypersurface $\Sigma \subset(M, \xi)$ is a folded Weinstein hypersurface if the characteristic foliation $\Sigma_{\xi}$ satisfies the following properties:
(FW1) There exist pairwise disjoint closed codimension 1 submanifolds $K_{i} \subset \Sigma$, $i=1, \ldots, 2 m-1$, which cut $\Sigma$ into $2 m$ pieces, i.e.,

$$
\Sigma=W_{1} \cup_{K_{1}} \cdots \cup_{K_{2 m-1}} W_{2 m}
$$

where $W_{i}$ are compact with boundary $\partial W_{i}=K_{i} \cup K_{i-1}$. Here we are setting $K_{0}=K_{2 m}=\varnothing$. We call $K_{i}$ the folding loci of $\Sigma$.
(FW2) The singular points of $\Sigma_{\xi}$ in each $W_{i}$ have the same sign, and the sign changes when crossing $K_{i}$. We assume the singular points in $W_{1}$ are positive and that each $W_{i}$ has at least one singular point.
(FW3) There exists a Morse function $f_{i}$ on each $W_{i}$ such that $K_{i-1}$ and $K_{i}$ are regular level sets and $\left(W_{i}\right)_{\xi}$ is gradient-like with respect to $f_{i}$. In particular, $\Sigma_{\xi}$ is transverse to all the $K_{i}$.

Observe that if $\Sigma=W_{1} \cup \cdots \cup W_{2 m} \subset(M, \xi)$ is a folded Weinstein hypersurface, then there exists a contact form $\alpha$ for $\xi$ whose restriction to the interior of each $W_{i}$ defines a Weinstein cobordism (the argument is similar to that of Proposition 2.3.3). Moreover, the orientation on $W_{i}$ given by the Weinstein structure agrees with (resp. is opposite to) the orientation inherited from $\Sigma$ if the singular points of $\left(W_{i}\right)_{\xi}$ are positive (resp. negative). We say a folding locus $K_{i}$ is maximal (resp. minimal) if the Liouville vector fields on $W_{i}$ and $W_{i+1}$ are pointing towards (resp. away from) $K_{i}$.


Figure 8.1.1. A schematic picture of a folded Weinstein hypersurface $\Sigma$. The top arrows indicate the direction of the Liouville vector fields on the $W_{i}$ and the bottom arrow indicates the direction of the characteristic foliation.

Note that by definition any folded Weinstein hypersurface is Morse, and any Morse hypersurface can be equipped with the structure of a folded Weinstein hypersurface.

We now give examples of folded Weinstein hypersurfaces and explain why they are called "folded".

Example 8.1.2. If $\Sigma$ is a convex hypersurface such that $R_{ \pm}(\Sigma)$ are Weinstein manifolds, then $\Sigma$ can be given the structure of a folded Weinstein hypersurface where the folding locus coincides with the dividing set $\Gamma_{\Sigma}$ and is maximal. On the other hand, a folded Weinstein hypersurface is not always convex because there may exist trajectories of $\Sigma_{\xi}$ from a negative singularity to a positive one. Nevertheless, since a $C^{\infty}$-small perturbation of a Morse hypersurface is Morse ${ }^{+}$by Lemma 2.3.2, any folded Weinstein hypersurface is $C^{\infty}$-generically convex by Proposition 2.3.3.

Example 8.1.3. Consider $\left(\mathbb{R}^{2 n+1}, \xi_{\text {std }}\right)$ with contact form $\alpha=d z+\sum_{i=1}^{n} r_{i}^{2} d \theta_{i}$. The unit sphere $S^{2 n}$ is convex with respect to the contact vector field $2 z \partial_{z}+$ $\sum_{i=1}^{n} r_{i} \partial_{r_{i}}$.

We slightly generalize this example as follows, which motivates our definition of a "folded" Weinstein hypersurface: We refer the reader to Definition A.3.1 for the definition of a $v$-folded hypersurface and a seam, where $v$ is a vector field; for example, the graph of $y=x^{2}$ is a $\partial_{x}$-folded hypersurface in $\mathbb{R}^{2}$. Taking $v$ to be $R_{\alpha}=\partial_{z}$, we consider closed $R_{\alpha}$-folded hypersurfaces $\Sigma \subset \mathbb{R}^{2 n+1}$ with seam $C$ and decomposition $\Sigma \backslash C=\Sigma_{+} \cup \Sigma_{-}$such that $R_{\alpha}$ is positively (resp. negatively) transverse to $\Sigma_{ \pm}$. It follows that $\Sigma_{ \pm}$are naturally exact symplectic manifolds with symplectic forms $\left.d \alpha\right|_{\Sigma_{ \pm}}$.

Now consider the following nontrivial condition:
(WC) Each component of $\Sigma_{ \pm}$is a (completed) Weinstein cobordism.
For example (WC) holds if each component of $C$ is contained in $\{z=$ const $\} \cong$ $\mathbb{R}^{2 n}$ and is transverse to the radial vector field. Moreover, $\Sigma_{ \pm}$are graphical over $\mathbb{R}^{2 n}$ and are Weinstein homotopic to subdomains of $\left(\mathbb{R}^{2 n}, \omega_{\text {std }}\right)$.

Any $R_{\alpha}$-folded hypersurface satisfying (WC) is clearly folded Weinstein with the folding locus equal to $C$, and hence can be made convex by a $C^{\infty}$-small perturbation. Note, however, that if $\Sigma$ happens to be convex, $R_{ \pm}(\Sigma) \neq \Sigma_{ \pm}$in general.

This explains our terminology but at the same time raises a hard problem:
Question 8.1.4. Characterize or classify convex hypersurfaces (e.g., spheres) in a Darboux chart.

Any answer to this question will be of fundamental importance in understanding contact manifolds. See [Eli92] for a complete answer to this question in the case $S^{2} \subset\left(\mathbb{R}^{3}, \xi_{\text {std }}\right)$.
8.2. Normalization of contact structure near a folded Weinstein hypersurface.

Recall that if $\Sigma \subset(M, \xi)$ is a convex hypersurface, then there exists a "standard" tubular neighborhood $U(\Sigma) \simeq \mathbb{R}_{t} \times \Sigma$ of $\Sigma$ such that $\left.\xi\right|_{U(\Sigma)}=\operatorname{ker}(f d t+\beta)$, where $f \in C^{\infty}(\Sigma)$ and $\beta \in \Omega^{1}(\Sigma)$.

The goal of this subsection is to generalize this to folded Weinstein hypersurfaces, namely associate to a folded Weinstein surface $\Sigma \subset(M, \xi=\operatorname{ker} \alpha)$ a standard tubular neighborhood, i.e., a tubular neighborhood $U(\Sigma)$ of $\Sigma$ and a contactomorphism $\phi: U(\Sigma) \xrightarrow{\sim}\left((-\epsilon, \epsilon)_{t} \times \Sigma, \operatorname{ker} \alpha_{\Sigma}\right)$, where $\alpha_{\Sigma}$ is a normalized contact form (specified below), $\epsilon>0$ is sufficiently small, and $\phi(\Sigma)=\Sigma_{0}$. Here $\Sigma_{t}:=\{t\} \times \Sigma$.

In the following three steps we construct the normalized contact form $\alpha_{\Sigma}$ on $\mathbb{R} \times \Sigma$ such that $\left.\alpha_{\Sigma}\right|_{\Sigma}=\left.\alpha\right|_{\Sigma_{0}}$, up to rescaling by a positive function. Lemma 2.1.3 then gives the desired contactomorphism.

Following Definition 8.1.1, we write $\Sigma=W_{1} \cup_{K_{1}} \cup \cdots \cup_{K_{2 m-1}} W_{2 m}$. Choose a tubular neighborhood $U\left(K_{i}\right)$ for each $K_{i}$ and identify it with $[-1,1]_{\tau} \times K_{i}$ such that $\Sigma_{\xi}$ is directed by $\partial_{\tau}$ on $U\left(K_{i}\right)$. (In particular, this means that $\{-1\} \times K_{i} \subset W_{i}$ and $\{1\} \times K_{i} \subset W_{i+1}$.)
STEP 1. Construct the contact form on $\mathbb{R} \times\left(\Sigma \backslash \cup_{i=1}^{2 m-1} U\left(K_{i}\right)\right)$.
Let $W_{i}^{\circ}:=W_{i} \backslash\left(U\left(K_{i-1}\right) \cup U\left(K_{i}\right)\right)$. After possibly rescaling $\alpha$ by a positive function as in Proposition 2.3.3, we may assume that $\beta_{i}:=\left.\alpha\right|_{W_{i}^{\circ}}$ is Liouville for all $i$. (Here $d \beta_{i}>0$ for $i$ odd and $d \beta_{i}<0$ for $i$ even.) Moreover, we can arrange so that the Liouville vector field $X_{\beta_{i}}$ equals $\partial_{\tau} /(2 \tau)$ near $\partial U\left(K_{i}\right)$ if $i$ is even, and equals $-\partial_{\tau} /(2 \tau)$ if $i$ is odd. (This is a purely technical arrangement which makes the gluing of contact forms below easier.) We define

$$
\begin{equation*}
\alpha_{\Sigma}:=(-1)^{i+1} d t+\beta_{i} \tag{8.2.1}
\end{equation*}
$$

on $\mathbb{R} \times\left(\Sigma \backslash \cup_{i=1}^{2 m-1} U\left(K_{i}\right)\right)$.
Step 2. Construct the contact form on $\mathbb{R} \times U\left(K_{i}\right)$ for $i$ even.
In this case $K_{i}$ is minimal. Assume without loss of generality that $\left.\alpha\right|_{U\left(K_{i}\right)}=$ $e^{\tau^{2}} \lambda$, where $\lambda$ is a contact form on $K_{i}$. We will choose $\alpha_{\Sigma}$ of the form

$$
\begin{equation*}
\alpha_{\Sigma}=-f(\tau) d t-\operatorname{tg}(\tau) d \tau+e^{\tau^{2}} \lambda \tag{8.2.2}
\end{equation*}
$$

on $\mathbb{R} \times U\left(K_{i}\right)$. Clearly $\left.\alpha_{\Sigma}\right|_{U\left(K_{i}\right)}=\left.\alpha\right|_{U\left(K_{i}\right)}$. A straightforward computation shows that $\alpha_{\Sigma}$ is positively contact if and only if

$$
\begin{equation*}
f^{\prime}-2 \tau f-g>0 \tag{8.2.3}
\end{equation*}
$$

We choose $f$ to be a decreasing odd function which equals $\pm 1$ when $\tau$ is close to $\mp 1$, and then choose $g$ to be a nonpositive even function which equals 0 when $\tau$ is close to $\pm 1$, subject to (8.2.3); see Figure 8.2.1.



Figure 8.2.1. The graph of functions used in the contact form given by Eq. (8.2.2).

For later use, note that $\alpha_{\Sigma}$ restricts to the Liouville form $\beta_{i, t}=-\operatorname{tg}(\tau) d \tau+e^{\tau^{2}} \lambda$ on $\{t\} \times\left(U\left(K_{i}\right) \backslash K_{i}\right)$ for any $t \in \mathbb{R}$. We compute the Liouville vector fields

$$
X_{\beta_{i, t}}=1 /(2 \tau)\left(\partial_{\tau}+t e^{-\tau^{2}} g(\tau) R_{\lambda}\right)
$$

where $R_{\lambda}$ denotes the Reeb vector field on $\left(K_{i}, \lambda\right)$.
It follows that

$$
\begin{equation*}
U\left(K_{i}\right)_{t, \xi}:=\left(\{t\} \times U\left(K_{i}\right)\right)_{\xi}=\partial_{\tau}+t e^{-\tau^{2}} g(\tau) R_{\lambda} \tag{8.2.4}
\end{equation*}
$$

STEP 3. Construct the contact form on $\mathbb{R} \times U\left(K_{i}\right)$ for $i$ odd.
In this case $K_{i}$ is maximal. This step is analogous to the construction of the contact form on $\Gamma \times[-1,1]$ in the proof of Proposition 2.3.3. Assume without loss of generality that $\left.\alpha\right|_{U\left(K_{i}\right)}=e^{-\tau^{2}} \lambda$. We define the 1-form

$$
\begin{equation*}
\alpha_{\Sigma}=f(\tau) d t+e^{-\tau^{2}} \lambda \tag{8.2.5}
\end{equation*}
$$

on $\mathbb{R} \times U\left(K_{i}\right)$. Since $\alpha_{\Sigma}$ is positively contact if and only if $-f^{\prime}-2 \tau f>0$, taking $f(\tau)$ as in Figure 8.2.1 suffices.

We compute the Liouville vector fields $X_{\beta_{i, t}}=-1 /(2 \tau) \partial_{\tau}$ on $\{t\} \times\left(U\left(K_{i}\right) \backslash\right.$ $K_{i}$ ), and note that it is independent of $t$. In fact $U\left(K_{i}\right)$ is convex with respect to the contact vector field $\partial_{t}$.

Combining (8.2.1), (8.2.2), and (8.2.5), we obtain a contact form $\alpha_{\Sigma}$ on $\mathbb{R} \times \Sigma$ such that $\left.\alpha_{\Sigma}\right|_{\Sigma}=\left.\alpha\right|_{\Sigma_{0}}$, up to rescaling by a positive function.

Remark 8.2.1. A crucial difference between the normal forms of contact structures near a convex hypersurface and a folded Weinstein hypersurface is the following: For convex hypersurfaces, since $\partial_{t}$ is a transverse contact vector field, any small
neighborhood of $\Sigma$ is in fact contactomorphic to the entire $\mathbb{R}_{t} \times \Sigma$. On the other hand, the above-constructed $\alpha_{\Sigma}$ is not $t$-invariant, and hence only specified by the data on $\Sigma$ for $|t|$ sufficiently small.
8.3. Bypass attachment as a bifurcation. Let $\xi$ be a contact structure on $\Sigma \times$ $[0,1]$ such that $\Sigma \times\{0,1\}$ are Weinstein convex, or equivalently Morse ${ }^{+}$. Then by Proposition 2.3.3, $\Sigma_{t}$ is Morse ${ }^{+}$for all $t \in[0, \epsilon) \cup(1-\epsilon, 1]$ with $\epsilon>0$ small. It will be shown in Section 9 that after a boundary-relative isotopy, $\Sigma_{t}$ is 1-Morse for all $t \in[0,1]$. However, the 1 -Morse ${ }^{+}$condition may fail at isolated instances. In particular, there may exist a first instance $t_{0}>0$ such that $\Sigma_{t}$ is convex for any $t \neq t_{0}$ sufficiently close to $t_{0}$ but there exists a "retrogradient" flow line of $\left(\Sigma_{t_{0}}\right)_{\xi}$ from a negative index $n$ singularity to a positive one. Such a phenomenon is called a bifurcation of the characteristic foliation in [Gir00]. ${ }^{2}$ As we will see, crossing such $t_{0}$ corresponds precisely to a bypass attachment as introduced in [ HH ].

To set up the "bypass-bifurcation correspondence", it is convenient to reformulate the bypass attachment in the language of folded Weinstein hypersurfaces.
8.3.1. Bypass attachments. We briefly review bypass attachments from [HH], leaving the details of contact handle attachments and Legendrian (boundary) sums to [HH].

Let $\Sigma$ be a Weinstein convex hypersurface with the usual decomposition $\Sigma \backslash$ $\Gamma=R_{+} \cup R_{-}$. A bypass attachment data $\left(\Lambda_{ \pm} ; D_{ \pm}\right)$is given as follows: Let $D_{ \pm} \subset R_{ \pm}$be Lagrangian disks with cylindrical ends which are regular in the sense of [EGL18], i.e., the complement in $R_{ \pm}$of a standard neighborhood of $D_{ \pm}$is still Weinstein. Let $\Lambda_{ \pm}=\partial \bar{D}_{ \pm}$be Legendrian spheres in $\Gamma$ equipped with the contact form $\left.\alpha\right|_{\Gamma}$, which we assume have a unique $\left.\xi\right|_{\Gamma}$-transversal intersection point (i.e., they intersect transversely when projected to $\left.\xi\right|_{\Gamma}$ ) .

Next we discuss Reeb pushoffs. If $\Lambda$ is a Legendrian submanifold of $\Gamma$, then let $\Lambda^{\epsilon}$ be the Reeb pushoff of $\Lambda$ in the Reeb direction by $\epsilon$. Clearly $\Lambda^{\epsilon}$ is embedded for $|\epsilon|$ sufficiently small. Moreover, if $\Lambda$ bounds a Lagrangian disk $D$ in some Weinstein filling, then there exists a corresponding Lagrangian $D^{\epsilon}$ in the same filling with $\partial D^{\epsilon}=\Lambda^{\epsilon}$.

We now explain how to attach a bypass to $\Sigma$ using the bypass attachment data $\left(\Lambda_{ \pm} ; D_{ \pm}\right)$to obtain a contact structure on $\Sigma \times[0,1]$. The bypass attachment is a smoothly canceling pair of contact handle attachments in the middle dimensions. The first is a contact $n$-handle attachment to $\Sigma_{0}$ along the Legendrian sphere $\Lambda_{-} \uplus \Lambda_{+} \subset \Gamma$ obtained by Legendrian sum. This step produces a new convex hypersurface $S$. It turns out the pushoffs $\Lambda_{ \pm}^{\mp \epsilon}$ of $\Lambda_{ \pm}$become Legendrian isotopic when viewed on $\Gamma_{S}$. Hence we can attach a contact ( $n+1$ )-handle to $S$ along the Legendrian sphere that we denote by $D_{+}^{-\epsilon} \cup D_{-}^{+\epsilon}$ and is obtained by gluing $D_{+}^{-\epsilon}$ and $D_{-}^{+\epsilon}$ via the Legendrian isotopy.

[^1]Remark 8.3.1. It is not necessary to assume that $D_{ \pm}$are regular in the definition of a bypass attachment. It is an outstanding, and of course hard, problem to even find an irregular Lagrangian disk in any Weinstein domain. One consequence of our work in this paper is that, as far as convex hypersurface theory and open book decompositions are concerned, one can completely stay in the world of Morse theory, e.g., avoid using any irregular Lagrangian disks, regardless of their very existence, without losing any generality. Compare this with the work of Lazarev [Laz20] which proves the existence $h$-principle for regular Lagrangians.

Let $(\Sigma \times[0,1], \xi)$ be the contact manifold resulting from the bypass attachment. Write $\Sigma_{t}:=\Sigma \times\{t\}$, where $\Sigma=\Sigma_{0}$. We have the usual decomposition $\Sigma_{i} \backslash \Gamma_{i}=$ $R_{+}^{i} \cup R_{-}^{i}, i=0,1$. Then by [HH, Theorem 5.1.3]:

- $R_{+}^{1}$ is obtained from $R_{+}^{0}$ by removing a standard neighborhood of $D_{+}^{-\epsilon}$ and attaching a Weinstein handle along $\Lambda_{-} \uplus \Lambda_{+}$.
- $R_{-}^{1}$ is obtained from $R_{-}^{0}$ by removing a standard neighborhood of $D_{-}^{+\epsilon}$ and attaching a Weinstein handle along $\Lambda_{-} \uplus \Lambda_{+}$.
- $\Gamma_{1}$, viewed as the boundary of $R_{+}^{1}$, is obtained from $\Gamma_{0}$ by a contact $(+1)-$ surgery along $\Lambda_{+}^{-\epsilon}$ and a contact ( -1 )-surgery along $\Lambda_{-} \uplus \Lambda_{+} . \Gamma_{1}$, viewed as the boundary of $R_{-}^{1}$, is obtained from $\Gamma_{0}$ by a contact $(+1)$-surgery along $\Lambda_{-}^{+\epsilon}$ and a contact ( -1 )-surgery along $\Lambda_{-} \uplus \Lambda_{+}$. These two presentations of $\Gamma_{1}$ are canonically identified by a handleslide.
8.3.2. Folded Weinstein description. We will now describe the Morse ${ }^{+}$hypersurface $\Sigma$ as a folded Weinstein hypersurface.

Let $W_{1} \subset R_{+}$be the Weinstein subdomain obtained by digging out a standard neighborhood of $D_{+}^{-\epsilon}$. Then $D_{+}^{-\epsilon}$ is the unstable manifold of an index $n$ critical point $q_{+}$with respect to the Liouville flow on $R_{+}$and $R_{+}$can be viewed as the concatenation $W_{1} \cup W_{3}^{\prime}$, where $W_{3}^{\prime}$ is a Weinstein cobordism with a unique critical point $q_{+}$. Similarly, let $W_{4} \subset R_{-}$be the Weinstein subdomain such that $R_{-}$is the concatenation of $W_{4}$ and a Weinstein cobordism $W_{2}^{\prime}$ with a unique index $n$ critical point $q_{-}$, whose unstable manifold is $D_{-}^{\epsilon}$ with respect to the Liouville flow on $R_{-}$. Since $\Lambda_{ \pm}^{\mp \epsilon}=\partial D_{ \pm}^{\mp \epsilon}$ are disjoint, we can shuffle the critical values of $q_{ \pm}$to obtain the following decomposition

$$
\Sigma=W_{1} \cup_{K_{1}} W_{2} \cup_{K_{2}} W_{3} \cup_{K_{3}} W_{4},
$$

where $W_{2}, W_{3}$ are Weinstein cobordisms (slight variants of $W_{2}^{\prime}, W_{3}^{\prime}$ ) associated with the critical points $q_{-}, q_{+}$, respectively. See Figure 8.3.1.

In particular we have:
(FBP1) As contact manifolds, $K_{1}$, oriented as $\partial W_{1}$ (and also as $\partial W_{2}$ ), is obtained from $\Gamma$ by a contact $(+1)$-surgery along $\Lambda_{+}^{-\epsilon} ; K_{3}$, oriented as $\partial W_{3}$ (and also as $\partial W_{4}$ ), is obtained from $\Gamma$ by a contact ( +1 )-surgery along $\Lambda_{-}^{\epsilon}$; and $K_{2}$, oriented as $-\partial W_{2}$ (and also as $-\partial W_{3}$ ), is obtained from $\Gamma$ by contact $(+1)$-surgeries along $\Lambda_{+}^{-\epsilon}$ and $\Lambda_{-}^{\epsilon}$.
(FBP2) Let $\left(D_{+}^{-\epsilon}\right)^{\dagger}$ and $\left(D_{-}^{\epsilon}\right)^{\dagger}$ be the stable manifolds of $q_{+}$and $q_{-}$in $W_{3}$ and $W_{2}$ with respect to the Liouville flows. Then $\left(D_{+}^{-\epsilon}\right)^{\dagger} \cap K_{2}=\Lambda_{+}^{-\epsilon}\left(K_{2}\right)$ and
$\left(D_{-}^{\epsilon}\right)^{\dagger} \cap K_{2}=\Lambda_{-}^{\epsilon}\left(K_{2}\right)$, where we write $\Lambda_{+}^{-\epsilon}\left(K_{2}\right)$ and $\Lambda_{-}^{\epsilon}\left(K_{2}\right)$ for the core Legendrians of the contact $(+1)$-surgeries.
(FBP3) The $\left.\xi\right|_{\Gamma}$-transverse intersection point between $\Lambda_{+}$and $\Lambda_{-}$turns into a short Reeb chord $\gamma \subset K_{2}$ from $\Lambda_{+}^{-\epsilon}\left(K_{2}\right)$ to $\Lambda_{-}^{\epsilon}\left(K_{2}\right)$.


Figure 8.3.1. A Morse ${ }^{+}$hypersurface viewed as a folded Weinstein hypersurface.

### 8.3.3. Bypass-bifurcation correspondence.

Proposition 8.3.2 (Bypass-bifurcation correspondence). Let $\left(\Sigma \times[-\delta, \delta], \alpha_{\Sigma}\right)$ with $\delta>0$ small be a standard neighborhood of a $2 n$-dimensional folded Weinstein hypersurface

$$
\begin{equation*}
\Sigma_{0}=W_{1} \cup_{K_{1}} W_{2} \cup_{K_{2}} W_{3} \cup_{K_{3}} \cup W_{4}, \tag{8.3.1}
\end{equation*}
$$

where we are writing $\Sigma_{t}:=\Sigma \times\{t\}$ as before, such that:
(F1) $W_{2}$ and $W_{3}$ are Weinstein cobordisms associated with ind $=n$ negative and positive critical points $q_{-}$and $q_{+}$. Let $D_{-, 0}^{\dagger}$ and $D_{+, 0}^{\dagger}$ (resp. $D_{-, 0}$ and $D_{+, 0}$ ) be the stable (resp. unstable) manifolds of $q_{-}$and $q_{+}$in $W_{2}$ and $W_{3}$ with respect to the Liouville flows.
(F2) $\Lambda_{-}^{\prime}:=\partial D_{-, 0}^{\dagger}$ and $\Lambda_{+}^{\prime}:=\partial D_{+, 0}^{\dagger}$ are Legendrians which intersect at a unique $\left.\xi\right|_{K_{2}}$-transversal point.
Then the contact manifold $\left(\Sigma \times[-\delta, \delta], \operatorname{ker} \alpha_{\Sigma}\right)$ is contactomorphic, relative boundary, to the bypass attachment to $\Sigma_{-\delta}$ along a quadruple $\left(\Lambda_{ \pm} ; D_{ \pm}\right)$.

Under the conditions of Proposition 8.3.2,
(Fa) $\Sigma_{t}$ is convex for all $t \neq 0$ by (8.2.1), (8.2.4), and (8.2.5);
(Fb) $\Sigma_{t}$ for $t<0($ resp. $t>0)$ has a folded Weinstein structure which satisfies (8.3.1) and (F1) with stable and unstable manifolds $D_{-, t}^{\dagger}, D_{+, t}^{\dagger}, D_{-, t}$, $D_{+, t}$, and $\partial D_{-, t}^{\dagger}$ is a positive (resp. negative) Reeb pushoff of $\partial D_{-, 0}^{\dagger}$ and $\partial D_{+, t}^{\dagger}$ is a negative (resp. positive) Reeb pushoff of $\partial D_{+, 0}^{\dagger}$.
Moreover, comparing with the notation from §8.3.2, $\Sigma=\Sigma_{-\delta}, D_{+}^{-\epsilon}=D_{+,-\delta}$, $D_{-}^{\epsilon}=D_{-,-\delta},\left(D_{+}^{-\epsilon}\right)^{\dagger}=D_{+,-\delta}^{\dagger},\left(D_{-}^{\epsilon}\right)^{\dagger}=D_{-,-\delta}^{\dagger}$, and $\Lambda_{ \pm}^{\mp \epsilon}=\partial D_{ \pm}^{\mp \epsilon}$. We then take $\Lambda_{ \pm}$and $D_{ \pm}$so that $\Lambda_{ \pm}^{\mp \epsilon}$ and $D_{ \pm}^{\mp \epsilon}$ are pushoffs of $\Lambda_{ \pm}$and $D_{ \pm}$. We will use both sets of notations interchangeably.

Since the proof of Proposition 8.3.2 is somewhat complicated, we start by explaining the key ideas involved and also highlight the difference between the usual 3-dimensional strategy and the higher-dimensional approach.

First note that the bypass attachment is a local operation, i.e., the hypersurface is only affected in a neighborhood of $D_{+} \cup D_{-}$. Let $B \subset \Sigma$ be a small neighborhood of $D_{+} \cup D_{-}$which is diffeomorphic to a ball. The question is then reduced to understanding the contact structure on $B \times I$ given by the bypass attachment, where $I=[-\delta, \delta]$.

At this point, two "miracles" happen in dimension 3 (i.e., $\operatorname{dim} \Sigma=2$ ) which greatly simplify the 3 -dimensional proof. The first is that one can take $\partial B$ to be Legendrian using the Legendrian realization principle (see [Hon00, Theorem 3.7]). This gives us good control over the contact structure near $\partial B \times I$. The second, and more significant, miracle is Eliashberg's theorem (see [Eli92, Theorem 2.1.3]) on the uniqueness of tight contact structures on the 3 -ball. Using these two facts, one can prove Proposition 8.3.2 in dimension 3 by arguing that both the bifurcation and the bypass attachment produce tight contact structures on the 3 -ball $B \times I$ with the same boundary conditions and hence must coincide.

Unfortunately, both of the above-mentioned miracles fail in dimension $>3$ : the first one fails for dimensional reasons and the second one fails by results of [Eli91, Ust99a]. Nevertheless, the proof of Proposition 8.3.2 follows the same general outline as in dimension 3 by replacing the Legendrian boundary condition on $\partial B$ by a transverse boundary condition and Eliashberg's theorem by a direct proof that both the bifurcation and the (trivial) bypass attachment produce the standard ball in a Darboux chart.

Proof of Proposition 8.3.2. The proof follows the above outline and consists of several steps.

Step 1. Localizing the problem to $B$.
By (FBP1), the contact manifold $K_{2}$ is obtained from $\Gamma$ by a contact ( +1 )surgeries along $\Lambda_{+}^{-\epsilon}$ and $\Lambda_{-}^{\epsilon}$; we are also viewing $K_{2}$ as a submanifold of $\Sigma_{t}$ for all $t \in I$. Let $\Lambda_{ \pm}^{\prime}:=\partial D_{ \pm, 0}^{\dagger} \subset K_{2} \subset \Sigma_{0}$ be the Legendrian spheres which $\left.\xi\right|_{K_{2}-}$ transversely intersect at one point and let $D_{ \pm, t}^{\dagger} \subset \Sigma_{t}$ be the Lagrangian disks in $W_{3}$ and $W_{2}$ from ( Fb ) with boundary $\Lambda_{ \pm, t}^{\prime}$.

We now describe a small closed neighborhood $B$ of $D_{+, 0}^{\dagger} \cup D_{-, 0}^{\dagger}$ in $\Sigma_{0}$. We take $B \cap K_{2}$ to be a small contact handlebody neighborhood $C_{2}=[-\kappa, \kappa]_{z} \times A_{2}, \kappa>0$ small, of $\Lambda_{+}^{\prime} \cup \Lambda_{-}^{\prime}$, where $A_{2}$ is the plumbing of two copies of disk bundles $\mathbb{D}^{*} S^{n-1}$ with the canonical Liouville form and $\Lambda_{ \pm}^{\prime}$ are the 0 -sections of the corresponding $\mathbb{D}^{*} S^{n-1}$ in $\{0\} \times A_{2}$. The restriction of $B$ to a tubular neighborhood $[-1,1]_{\tau} \times K_{2}$ with the 1 -form $e^{\tau^{2}} \lambda$ is $[-1,1] \times C_{2}$. Then $B$ is obtained from $[-1,1] \times C_{2}$ by attaching Weinstein handles along $\{-1\} \times \Lambda_{-}^{\prime}$ and $\{1\} \times \Lambda_{+}^{\prime}$. The boundary decomposes as $\partial B=C_{1} \cup C_{h} \cup C_{3}$, where $C_{1}$ (resp. $C_{3}$ ) is the compact contact manifold obtained from $C_{2}$ by a contact ( -1 )-surgery along $\{-1\} \times \Lambda_{-}^{\prime}$ (resp.
$\left.\{1\} \times \Lambda_{+}^{\prime}\right)$ and $C_{h}=[-1,1] \times \partial C_{2}$. We are viewing $B \subset W_{2} \cup W_{3}, C_{1} \subset K_{1}$, and $C_{3} \subset K_{3}$.

Note that the $C_{i}, i=1,2,3$, are all contactomorphic since applying a contact $(-1)$-surgery along $\{0\} \times \Lambda_{+}^{\prime}$ to $[-\kappa, \kappa] \times T^{*} \Lambda_{+}^{\prime}$ still yields $[-\kappa, \kappa] \times T^{*} \Lambda_{+}^{\prime}$. The case for $\{0\} \times \Lambda_{-}^{\prime}$ is identical.

The characteristic foliation $B_{\xi}$ is inward-pointing along $C_{1}$, outward-pointing along $C_{3}$, and tangent to $C_{h}$. By slightly tilting $C_{h}$, we may assume that $B_{\xi}$ is outward-pointing along $C_{3} \cup C_{h}$ and inward-pointing along $C_{1}$. This results in a fold-type tangency roughly along $\partial C_{1}$.

Moreover, for $\delta>0$ sufficiently small, we can construct parallel copies $B_{t} \subset$ $\Sigma_{t}, t \in I$, of $B=B_{0}$ such that the characteristic foliation is $t$-invariant near $\partial B_{t}$. The copies $B_{t}$ are obtained from $[-1,1] \times A_{2}$ by attaching cores of the handles along $\{-1\} \times \Lambda_{-, t}^{\prime}$ and $\{1\} \times \Lambda_{+, t}^{\prime}$, and $\delta>0$ small ensures that we can attach handles (i.e., the thickened cores) to $\{ \pm 1\} \times C_{2}$ in a manner that varies smoothly with $t$.

By using certain folding techniques similar to (and in fact simpler than) those in Section 5, one can reverse the direction of the characteristic foliation on $C_{1}$ through an isotopy of $B$ in a suitably wiggled $\Sigma$ such that $\Sigma_{\xi}$ is everywhere outwardpointing along $\partial B$. This will be achieved in Step 3 . The folding technique is called the Creation Lemma which in dimension 3 is the converse of the usual Elimination Lemma (see [Gei08, §4.6.3]). This is described in Step 2.

## STEP 2. The Creation Lemma.

In this step, we describe the effect of applying a $C^{0}$-small perturbation called a box-fold. This is the content of the Creation Lemma, which we do not state formally.

We closely follow the discussion of Section 5 , except that we replace $t$ by $\tilde{t}$ here, since we are already using $t$ to parametrize the hypersurfaces $\Sigma_{t}$. Consider $\mathbb{R}_{z, s, \tilde{t}}^{3} \times V$ equipped with the contact form $\alpha=d z+e^{s}(d \tilde{t}+\lambda)$, where $(V, \lambda)$ is a complete Weinstein manifold. Let $F:=\{z=0\}$ be the hypersurface on which we will create singularities. Clearly $F_{\xi}=\partial_{s}$ where $\xi=\operatorname{ker} \alpha$.

Fix $z_{0}, s_{0}, \tilde{t}_{0}>0$. Let $\Pi^{3} \subset \mathbb{R}_{z, s, \tilde{t}}^{3}$ be a surface obtained from the flat $\mathbb{R}_{s, \tilde{t}}^{2}$ by growing a box with base $\square:=\left[0, s_{0}\right] \times\left[0, \tilde{t}_{0}\right]$ and height $z_{0}$; see Figure 8.3.2. Of course, as in the construction of $Z$ in $\S 3.2$, one needs to smooth the corners of $\Pi^{3}$ and "Morsify" the resulting characteristic foliation $\Pi_{\xi}^{3}$. These operations are suppressed from the notation. We say $\Pi^{3}$ is obtained from $\mathbb{R}_{s, \tilde{t}}^{2}$ by a 3-dimensional box-fold.

Comparing Figure 8.3.2 with Figure 3.1.1, we note that the key difference is that $\Pi_{\xi}^{3}$ admits only positive singularities: one source and one saddle. See Figure 8.3.3. This is the content of the Creation Lemma in dimension 3.

In higher dimensions, we consider the hypersurface $\Pi^{3} \times V^{c}$, where $V^{c} \subset V$ is the compact domain, i.e., $V \backslash V^{c} \cong[0, \infty)_{\tau} \times \partial V^{c}$ is symplectomorphic to a half-symplectization of $\partial V^{c}$. Following the strategy from Section 5, define the general box-fold $\Pi$ to be the hypersurface which extends $\Pi^{3} \times V^{c}$ by damping out


Figure 8.3.2. The PL model of $\Pi^{3}$.


Figure 8.3.3. The characteristic foliation before and after the box folding.
the $\Pi^{3}$-factor on $\mathbb{R}^{3} \times\left[0, \tau_{1}\right] \times \partial V^{c}$ as $\tau$ increases. Then Eq. (5.2.1) implies that $\Pi_{\xi}$ is Morse with a pair of canceling critical points for each one in $V^{c}$. In particular, let $D \subset \Pi^{3}$ be a disk containing the source $e_{+}$such that $\Pi_{\xi}^{3}$ is transverse to $\partial D$. Then $\Pi_{\xi}$ is everywhere outward-pointing along $\partial\left(D \times V^{c}\right)$. Note that instead of creating a pair of canceling critical points as in dimension 3 (see [CE12, Proposition 12.21] for the higher-dimensional version), our Creation Lemma produces many pairs of canceling critical points at once, in fact as many as the number of critical points of $V$.
STEP 3. Modification from $B$ to $\widehat{B}$.
In Step 1 we constructed the family $B_{t} \subset \Sigma_{t}, t \in I$, such that $\Sigma_{t, \xi}$ is inwardpointing along $C_{1}$. The goal of this step is to modify $B_{t}$ (and $\Sigma_{t, \xi}$ ) to $\widehat{B}_{t}$ so that $\Sigma_{t, \xi}$ is outward-pointing along $\partial \widehat{B}_{t}$. Write $C_{1}=[-\kappa, \kappa]_{\tilde{t}} \times A_{2}$, where $\kappa>0$ is sufficiently small. Since $\Sigma_{\xi}$ points into $B$ along $C_{1}$, we can choose $t_{0}, s_{0}>0$ so that there exists an embedding

$$
U\left(C_{1}\right):=I_{t_{0}} \times\left[0, s_{0}\right] \times C_{1} \subset M, \quad I_{t_{0}}=\left[-t_{0}, t_{0}\right],
$$

such that $t_{0} \ll \delta, U\left(C_{1}\right) \cap B_{t}=\left\{\left(t, s_{0}\right)\right\} \times C_{1}$, and $\{t\} \times\left[0, s_{0}\right] \times C_{1} \subset W_{1}$. Write the contact form as $\left.\alpha\right|_{U\left(C_{1}\right)}=d t+e^{s}(d \tilde{t}+\lambda)$, where $\lambda$ is the standard Liouville form on $A_{2}$. Here we are identifying the $z$-coordinate from the previous step with $t$.

We then apply the Creation Lemma with $V=A_{2}$ to install/uninstall a smoothed box-fold along $\left[0, s_{0}\right] \times C_{1}$, which we are assuming is contained in $W_{1}$. This is the higher-dimensional analog of the procedure in Section 4.4 (see Figure 4.4.1). More precisely, we modify the foliation of $\Sigma \times I$ by leaves $\Sigma_{t}$ to obtain a $C^{0}$-close foliation by leaves still denoted by $\Sigma_{t}$ such that the following hold:
(i) $\Sigma_{t, \xi}$ is unchanged for $t \in \partial I$ and on $W_{2} \cup W_{3} \cup W_{4}$ for $t \in I$.
(ii) The box-fold is installed along $\left[0, s_{0}\right] \times C_{1}$ for $t \in\left[-\delta,-t_{0}\right]$ and is uninstalled along $\left[0, s_{0}\right] \times C_{1}$ for $t \in\left[t_{0}, \delta\right]$. The characteristic foliation on $W_{1}$ is $t$-invariant for $t \in I_{t_{0}}$. (This is possible since the smoothed box-folds can be taken to be graphical.)
(iii) For $t \in I_{t_{0}}, W_{1}$ contains a subdomain symplectomorphic to $D \times A_{2}$, where $D \subset \mathbb{R}^{2}$ is a disk containing $e_{+}$as in Step 2 and there exists an arc $\mu \subset \partial D$ such that $\mu \times A_{2}$ is identified with $\left\{s_{0}\right\} \times C_{1}$; see Figure 8.3.4.


Figure 8.3.4. A schematic picture for isotoping $B$ to $\widehat{B}$ so it encompasses $e_{+}$. The blue arc represents $\mu$ and the red arc represents $\partial D \backslash \mu$.

In order to achieve the transversal boundary condition on $B_{t}, t \in I_{t_{0}}$, it remains to isotop $\mu$ through $D$ to $\partial D \backslash \mu$ and use the fact from Step 2 that $\Sigma_{t, \xi}$ is everywhere transverse to $\partial\left(D \times A_{2}\right)$, to obtain the new $\widehat{B}_{t} \subset \Sigma_{t}$ such that $\Sigma_{t, \xi}$ is everywhere outward-pointing along $\partial \widehat{B}_{t}$. In particular $\partial \widehat{B}_{t}, t \in I_{t_{0}}$, are contact submanifolds of $M$.

Remark 8.3.3. Similar ideas will be exploited in greater generality in Section 11.
Claim 8.3.4. $\Sigma_{t, \xi}$ is 1 -Morse for all $t \in I$ and 1 -Morse ${ }^{+}$(hence $\Sigma_{t}$ is convex) for $t \neq 0$.

Proof of Claim 8.3.4. By the folded Weinstein structure for $\Sigma_{t}$ and the fact that installing/uninstalling the box-fold induces a Weinstein homotopy on $W_{1}$, it follows that $\Sigma_{t, \xi}$ is 1 -Morse for all $t \in I$. As for the 1 -Morse ${ }^{+}$property, it suffices to consider the stable manifold of the unique singular point in $W_{3}$. There is a stable flow line that comes from a negative singularity precisely when $t=0$ (this is the
same as the situation before $\Sigma_{t}$ was perturbed). The convexity of $\Sigma_{t}, t \neq 0$, then follows from Proposition 2.3.3.

Hence we may restrict attention to the new $\Sigma_{t}, t \in I_{t_{0}}$.
STEP 4. Triviality of the contact structure on $I_{t_{0}} \times \widehat{B}$ for $t_{0}>0$ small.
Let $S:=\partial\left(I_{t_{0}} \times \widehat{B}\right)=\widehat{B}_{-t_{0}} \cup \widehat{B}_{t_{0}} \cup\left(I_{t_{0}} \times \partial \widehat{B}\right)$. Suppose $t_{0}>0$ is small.
Claim 8.3.5. After corner rounding, $S$ is convex and $R_{ \pm}(S)$ are Weinstein homotopic to the standard ball with a unique critical point. ${ }^{3}$

Proof of Claim 8.3.5. For $R_{+}(S)$ we describe the positive critical points and the stable manifolds of these critical points; the situation for $R_{-}(S)$ is similar.

The critical points of $R_{+}(S)$ are as follows:
(1) sitting over $e_{+}$in $\widehat{B}_{t_{0}}$ are one index 0 critical point $q_{0}$ and two index $(n-1)$ critical points $q_{ \pm}$corresponding to $\Lambda_{ \pm}^{\prime} \subset A_{2}$; and
(2) the critical points $p_{+}$on $\widehat{B}_{t_{0}} \cap W_{3}$ and $p_{-}$on $\widehat{B}_{-t_{0}} \cap W_{2}$ have index $n$, where + indicates being on the "top sheet" $\widehat{B}_{t_{0}}$.
We denote the analogous critical points of $R_{-}(S)$ by $q_{0}^{\prime}, q_{ \pm}^{\prime}, p_{ \pm}^{\prime}$.
We denote the stable manifold of a critical point $p$ by $\mathscr{W}_{p}$. For definiteness, we assume that there exists $\epsilon^{\prime}>0$ small such that
(i) $\mathscr{W}_{p_{+}}, \mathscr{W}_{p_{+}^{\prime}}$ intersect $C_{2} \subset \widehat{B}_{a_{\gamma}+t_{0}}$ along the pushoffs $\left(\Lambda_{+}^{\prime}\right)^{2 \epsilon^{\prime}},\left(\Lambda_{-}^{\prime}\right)^{-2 \epsilon^{\prime}}$;
(ii) $\mathscr{W}_{p_{-}}, \mathscr{W}_{p_{-}^{\prime}}$ intersect $C_{2} \subset \widehat{B}_{a_{\gamma}-t_{0}}$ along the pushoffs $\left(\Lambda_{-}^{\prime}\right)^{\epsilon^{\prime}},\left(\Lambda_{+}^{\prime}\right)^{-\epsilon^{\prime}}$.

By (i), $\mathscr{W}_{p_{+}}$intersects $K_{1}=\partial W_{1}$ along $\left(\Lambda_{+}^{\prime}\right)^{2 \epsilon^{\prime}}$ and therefore limits to $\Lambda_{+}^{\prime} \subset$ $A_{2}$ over $e_{+}$; moreover, there is a unique flow line from $p_{+}$to $q_{+}$. Next, $\mathscr{W}_{p_{-}}$ intersects $C_{3}$ along a Legendrian which is isotopic to a positive pushoff of $\Lambda_{-}^{\prime}$, continues inside $\widehat{B}_{a_{\gamma}+t_{0}}$ to a Legendrian isotopic to a positive pushoff of $\Lambda_{-}^{\prime}$ on $K_{1}$, and limits to $\Lambda_{-}^{\prime} \subset A_{2}$ over $e_{+}$. Moreover, there is a unique flow line from $p_{-}$to $q_{-}$. This implies that $S_{\xi}$ is Morse ${ }^{+}$and convex, with Weinstein structures on $R_{ \pm}(S)$ just described.

The index $(n-1)$ and index $n$ critical points cancel in pairs and $R_{ \pm}(S)$ are Weinstein homotopic to the standard ball with a unique critical point.

Remark 8.3.6. The reader might find it instructive to consider the $n=1$ (i.e., $\operatorname{dim} M=3$ ) case, where we have three index 0 critical points "sitting over $e_{+}$".

Now consider the 1-parameter family of pairwise disjoint homotopy spheres $S_{\widetilde{t}}$, $\widetilde{t} \in\left[0, \frac{3 t_{0}}{2}\right]$, that are slight perturbations of $\partial\left(\left[-t_{0}+\widetilde{t}, t_{0}\right] \times \widehat{B}\right)$ and such that the region $G$ bounded by $S_{3 t_{0} / 2}$ is $t$-invariant.
Claim 8.3.7. The homotopy spheres $S_{\tilde{t}}, \tilde{t} \in\left[0, \frac{3 t_{0}}{2}\right]$, can be made simultaneously convex after a small perturbation.

[^2]Proof of Claim 8.3.7. We use the argument of Claim 8.3.5, but this time there are two moments $t_{0}^{\prime}$ and $t_{0}$ with $t_{0}^{\prime}<t_{0}$, where a bifurcation occurs as explained in the next paragraph. There is still a unique flow line from $p_{+}$to $q_{+}$for all $\tilde{t}$.

Next we describe the trajectories of $\mathscr{W}_{p_{-}}$for $\widetilde{t} \in\left[0, \frac{3 t_{0}}{2}\right]$. When $\widetilde{t}=t_{0}^{\prime}$, all the trajectories of $\mathscr{W}_{p_{-}}$reach $C_{3}$ but one continues to the critical point of $\widehat{B}_{t_{0}} \cap W_{3}$; when $\tilde{t}=t_{0}$, a flow line of $\mathscr{W}_{p_{-}}$limits to the critical point of $\widehat{B}_{a_{\gamma}} \cap W_{3}$. We have:
(1) For $\widetilde{t}<t_{0}^{\prime}, \mathscr{W}_{p_{-}} \cap K_{1}$ is Legendrian isotopic to a positive pushoff of $\Lambda_{-}^{\prime}$.
(2) For $\tilde{t}>t_{0}, \mathscr{W}_{p_{-}} \cap K_{1}$ is Legendrian isotopic to a negative pushoff of $\Lambda_{-}^{\prime}$.
(3) For $t_{0}^{\prime}<\tilde{t}<t_{0}, \mathscr{W}_{p_{-}} \cap K_{1}$ is Legendrian isotopic to $\Lambda_{+}^{\prime} \uplus \Lambda_{-}^{\prime}$.
(4) For $\widetilde{t}=t_{0}^{\prime}, t_{0}, \bar{W}_{p_{-}} \cap K_{1}$ is Legendrian isotopic to $\Lambda_{+}^{\prime} \cup \Lambda_{-}^{\prime}$ intersecting at a point and $\mathscr{W}_{p_{-}} \cap K_{1}$ corresponds to $\Lambda_{-}^{\prime} \backslash \Lambda_{+}^{\prime}$.
Here $K_{1}$ is understood to be on $t=t_{0}$. (3) is a consequence of corner-rounding along $C_{3}$, which has the effect of introducing a slight negative Reeb flow along the corner as we go from the bottom sheet to the top. (4) is the limiting configuration of (1)-(3). In all the cases, there is a unique flow line from $p_{-}$to $q_{-}$, although there may be trajectories from $p_{-}$to $q_{+}$for $t_{0}^{\prime}<\tilde{t}<t_{0}$.

The claim then follows from the usual Elimination Lemma (see [CE12, Proposition 12.22]) and a trick from [Hua13, Lemma 3.3]: By a $C^{0}$-small perturbation one can simultaneously eliminate the pairs $\left(p_{+}, q_{+}\right)$and ( $p_{-}, q_{-}$) on $S_{\breve{t}}$ for all $\widetilde{t} \in\left[0, \frac{3 t_{0}}{2}\right]$ (since the trajectories from $p_{+}$to $q_{+}$and $p_{-}$to $q_{-}$vary continuously with respect to $\widetilde{t}$ ), which in turn implies that all the $S_{\tilde{t}}$ are convex.

Finally we observe that the bypass attachment to $\Sigma$ along $\left(\Lambda_{ \pm} ; D_{ \pm}\right)$restricts to the trivial bypass attachment to $S$ in the sense of [HH, Definition 6.1.1]. It follows from [HH, Proposition 8.3.2] that the contact structure on $\left(I_{t_{0}} \times \widehat{B}\right) \backslash G$ given by a trivial bypass attachment is standard. By Claim 8.3.7, $\xi$ on $\left(I_{t_{0}} \times \widehat{B}\right) \backslash G$ is standard, hence is equivalent to a bypass attachment. This finishes the proof of the proposition.
8.4. Bypass attachment in terms of partial open book decompositions. The goal of this subsection is to summarize the main constructions and results from [HH, Section 8]. The reader is referred to the original paper for details.

### 8.4.1. Partial open book decompositions.

Definition 8.4.1. Given a Weinstein domain $S$, a cornered Weinstein subdomain $W \subset S$ is a (possibly empty) codimension 0 submanifold with corners which satisfies the following properties:
(CW1) There exists a decomposition $\partial W=\partial_{\text {in }} W \cup \partial_{\text {out }} W$ such that
(1) $\partial_{\text {in }} W$ and $\partial_{\text {out }} W$ are compact manifolds with smooth boundary that intersect along their boundaries;
(2) $\partial\left(\partial_{\mathrm{in}} W\right)=\partial\left(\partial_{\mathrm{out}} W\right)$ is the codimension 1 corner of $\partial W$; and
(3) $\partial_{\text {out }} W=W \cap \partial S$ and is a proper subset of each component of $\partial S$.
(CW2) The Liouville vector field $X_{\lambda}$ on $S$ is inward-pointing along $\partial_{\mathrm{in}} W$ and outward-pointing near $\partial_{\text {out }} W$.
(CW3) $S \backslash W$ is a Weinstein domain after smoothing.
A particularly useful class of cornered Weinstein subdomains consists of regular neighborhoods of Lagrangian cocore disks in $S$.

Definition 8.4.2. A Weinstein partial open book decomposition (POBD) is a pair $\left([0,1] \times S, \phi: W_{1} \xrightarrow{\sim} W_{0}\right)$, where:
(1) $[0,1]_{t} \times S$ is a generalized contact handlebody (with $S_{t}:=\{t\} \times S$ Weinstein);
(2) $W_{0} \subset S_{0}$ and $W_{1} \subset S_{1}$ are cornered Weinstein subdomains; and
(3) $\phi$ is a partial monodromy map which preserves the Liouville forms and restricts to the identity map $\partial_{\text {out }} W_{1} \xrightarrow{\sim} \partial_{\text {out }} W_{0}$.

By taking the quotient of $[0,1] \times S$ by the map $\phi$ and filling in $D^{2} \times \partial S$ as in the closed case (note that this yields a concave sutured contact manifold and we need to round the concave sutures; see [CGHH11, Section 4.2] for details) we obtain a compact contact manifold with Weinstein convex boundary.
8.4.2. Bypass attachment in terms of POBDs. We follow the recipe from [HH, Section 8.3] to interpret a bypass attachment in terms of POBDs.

Lemma 8.4.3 ([HH], Proposition 8.3.1). Let $(M, \xi, \Gamma)$ be a compact contact manifold with Weinstein convex boundary associated to the Weinstein POBD $([0,1] \times$ $\left.S, \phi: W_{1} \rightarrow W_{0}\right)$. If $\left(M^{b}, \xi^{b}, \Gamma^{b}\right)$ is the contact manifold obtained by attaching a bypass to $M$ along $\partial M$ with data $\left(\Lambda_{ \pm} ; D_{ \pm}\right)$, where $D_{ \pm} \subset R_{ \pm}$are regular in the sense of [EGL18], then $\left(M^{b}, \xi^{b}, \Gamma^{b}\right)$ is associated to the Weinstein POBD $\left([0,1]_{t} \times S^{b}, \phi^{b}: W_{1}^{b} \rightarrow W_{0}^{b}\right)$ satisfying:
(PO1) $S^{b}$ is obtained from $S$ by attaching a Weinstein handle along $\Lambda_{-} \uplus \Lambda_{+}$; this is done independently of $t$;
(PO2) $W_{1}^{\mathrm{b}}=W_{1} \sqcup N_{\epsilon / 2}\left(D_{+}^{-\epsilon}\right) \subset S^{b}$, where $\epsilon>0$ is small and $N_{\epsilon / 2}\left(D_{+}^{-\epsilon}\right)$ is a cornered Weinstein subdomain which is a standard $\epsilon / 2$-neighborhood of the Lagrangian disk $D_{+}^{-\epsilon}$;
(PO3) $\phi^{b}: W_{1}^{b} \rightarrow W_{0}^{b}$ is determined by specifying the Lagrangian disk $\phi^{b}\left(D_{+}^{-\epsilon}\right) \subset$ $S_{0}^{b}$ : Start with the Lagrangian disk $D_{-}^{\epsilon} \subset S_{0}^{b} \backslash W_{0}^{b}$ with Legendrian boundary $\Lambda_{-}^{\epsilon} \subset \partial S_{0}^{b}$. We can slide $\Lambda_{-}^{\epsilon}$ in the negative Reeb direction across the Weinstein handle along $\Lambda_{-} \uplus \Lambda_{+}$so it precisely matches $\Lambda_{+}^{-\epsilon}=\partial D_{+}^{-\epsilon}$. The sliding is induced by a Weinstein isotopy $\tau_{s}: S_{0}^{b} \xrightarrow{\sim} S_{0}^{b}, s \in[0,1]$, with $\tau_{0}=$ id. Then we set $\phi^{b}\left(D_{+}^{-\epsilon}\right)=\tau_{1}\left(D_{-}^{\epsilon}\right)$.

Note that Lemma 8.4.3 is a direct consequence of the interpretation of a bypass attachment as a smoothly canceling pair of contact handle attachments. If $\operatorname{dim} M=2 n+1$, then the handles have indices $n$ and $n+1$.
8.4.3. Contact Morse functions and vector fields. It is helpful, although not technically necessary in this paper, to interpret the contact handle attachments in terms of contact Morse functions. This is the contact-topological analog of the correspondence between handle decompositions and Morse function presentations of a
given smooth manifold. See Sackel [Sac] for a more thorough discussion of contact Morse functions and a careful construction of contact handles. Sackel also works out the dictionary between contact Morse functions and open book decompositions and our proofs of Corollaries 1.3.1 and 1.3.2 follow the same line of reasoning.

Recall that a vector field $v$ on $(M, \xi)$ is a contact Morse vector field if $v$ is gradient-like for some Morse function $f: M \rightarrow \mathbb{R}$ and the flow of $v$ preserves $\xi$, i.e., $\mathscr{L}_{v} \alpha=g \alpha$ where $\xi=\operatorname{ker} \alpha$ and $g \in C^{\infty}(M)$. The Morse function $f$ is called a contact Morse function. The zeros of $v$ are precisely the critical points of $f$ and hence it makes sense to refer to the (Morse) indices of the zeros of $v$.

If $(\Sigma \times[0,1], \xi)$ is the contact manifold corresponding to a bypass attachment as above, then there exists a contact Morse vector field $v$ on $\Sigma \times[0,1]$ satisfying the following properties:
(BM1) $v$ is inward-pointing along $\Sigma_{0}$ and outward-pointing along $\Sigma_{1}$.
(BM2) $v$ has exactly two zeros - $p$ of index $n$ and $q$ of index $n+1$ - which are connected by a unique flow line of $v$.
(BM3) For any $x \in \Sigma \times[0,1]$, the flow line of $v$ passing through $x$ either converges to a zero of $v$ or leaves $\Sigma \times[0,1]$ in both forward and backward time.
(BM4) The unstable manifold of $p$ intersects $\Sigma_{0}$ along the Legendrian $\Lambda_{-} \uplus \Lambda_{+} \subset$ $\Gamma_{0}$, and the stable manifold of $q$ intersects $\Sigma_{1}$ along the Legendrian $\Lambda_{+}^{-\epsilon} \subset$ $\Gamma_{1}$, viewed as the boundary of $R_{+}^{1}$.
Note, however, that in general $\Sigma_{0}, \Sigma_{1}$ are not regular level sets of $f$ since contact vector fields are not stable under rescaling by positive functions.

## 9. $C^{0}$-APPROXIMATION BY CONVEX HYPERSURFACES

In this section we complete the proofs of Theorems 7.4.2 and 1.2.5. The main technical ingredient is the higher-dimensional plug constructed in Section 7. In fact our proofs are basically the same as those for the 3 -dimensional case discussed in Section 4.

Proof of Theorem 7.4.2. The proof is by induction on the dimension of $M$.
Fix a Riemannian metric on $M$. Given a closed hypersurface $\Sigma \subset(M, \xi)$, we may assume that the singularities of $\Sigma_{\xi}$ are isolated and Morse after a $C^{\infty}$-small perturbation. By Theorem 4.2.3 a finite barricade $B_{I}=\left\{B_{i}=\left[0, s_{i}\right] \times Y_{i}\right\}_{i \in I}$ exists for $\Sigma_{\xi}$, where each $Y_{i}$ is a standard $(2 n-1)$-dimensional Darboux ball.

Choose $\epsilon^{\prime}>0$ much smaller than the sizes of the $Y_{i}$. We then replace each $B_{i}$ with a $Y_{i}$-shaped plug with parameter $\epsilon^{\prime}$ constructed in Section 7; note that the construction of the $Y_{i}$-shaped plug uses the inductive step of Theorem 7.4.2 for $\operatorname{dim}(M)-2$. Let $\Sigma^{\vee}$ be the resulting hypersurface. A trajectory that passes near $Y_{i}$ is either trapped by $Y_{i}$ or has holonomy at most $\epsilon^{\prime}$. Since $\epsilon^{\prime}$ is small, the trajectory that continues is close to the original trajectory on $\Sigma_{\xi}$ and will be trapped by some other $Y_{j}$ by the positioning of the barricade. Hence $\Sigma_{\xi}^{\vee}$ satisfies Conditions (M1)(M3) of Proposition 2.2.6. A further $C^{\infty}$-small perturbation of $\Sigma^{\vee}$ will make it convex by Proposition 2.3.3.

The $\epsilon$-Morse ${ }^{+}$property is guaranteed if all the $B_{i}, i \in I$, have diameters $<\frac{\epsilon}{3}$ and all the trajectories of $\Sigma_{\xi} \backslash\left(\cup_{i \in I} B_{i}\right)$ have lengths $<\frac{\epsilon}{3}$.

Finally, the second statement (for $\Sigma$ already Weinstein convex) follows from observing that all singularities of all the $Y_{i}$-shaped plugs can be "turned on simultaneously"; see Remark 3.2.2 which also holds for higher dimensions.

Proof of Theorem 1.2.5. Let $(\Sigma \times[0,1], \xi)$ be a contact manifold such that the hypersurfaces $\Sigma_{i}, i=0,1$, are Weinstein convex. The proof from $\S 4.5$ carries over almost verbatim with "surface" replaced by "hypersurface". We use Lemma 4.2.4 to construct refinements of barricades so that we can install/uninstall two sets of barricades.

Suppose for simplicity that the barricade consists of one flow box $B=\left[0, s_{0}\right] \times$ $D^{2 n-1}$. Even though the folded hypersurface $\Sigma^{\vee}$ can be made convex, the intermediate hypersurfaces appearing in the procedure of installing and uninstalling a $D^{2 n-1}$-shaped plug need not be Morse. To remedy this defect, we take a cover $D^{2 n-1}=\cup_{1 \leq i \leq K} U_{i}$ by a finite number of balls of small diameter $\frac{1}{N}$ for which there exists a partition

$$
\mathscr{P}:\{1, \ldots, K\} \rightarrow\{1, \ldots, 2 n\}
$$

such that $U_{i} \cap U_{j}=\varnothing$ if $\mathscr{P}(i)=\mathscr{P}(j)$. Now we choose $2 n$ pairwise distinct values in $\left(0, s_{0}\right)$ and position $U_{i}$ along $\left[0, s_{0}\right]$ such that all the $U_{i}$ with the same $\mathscr{P}_{-}$ value have the same $s$-value. When we install/uninstall all the $U_{i}$-shaped plugs, the interior discrepancy for $B$ goes to zero as $N \rightarrow \infty$ by the analog of Lemma 4.4.1 and the fact that the $\epsilon$ that appears in Theorem 7.5.1 is one order of magnitude smaller than the sizes of $U_{i}$.

Finally, proceeding as in $\S 4.5$, we can foliate $\Sigma \times[0,1]$ by hypersurfaces of the form $\Sigma_{t}$, which we may assume are 1-Morse by the smallness of the interior discrepancy and Step 1 of $\S 4.5$. The only obstruction to convexity occurs when $\left(\Sigma_{t}\right)_{\xi}$ fails to be Morse ${ }^{+}$, which occurs at isolated moments by the description of a standard neighborhood of a non-Morse ${ }^{+}$folded Weinstein hypersurface in §8.2. The theorem then follows from Proposition 8.3.2.

## 10. The existence of (partial) open book decompositions

The goal of this section is to prove Corollary 1.3.1, a stronger/more precise version of Corollary 1.3.2, and Corollary 1.3.3. The proofs are, again, essentially the same as the proofs in the 3 -dimensional case; see [Gir02] for the absolute case and [HKM09] for the relative case.

Proof of Corollary 1.3.1. Let $(M, \xi)$ be a closed connected contact manifold of dimension $2 n+1$. Choose a generic self-indexing Morse function $f: M \rightarrow$ $\mathbb{R}$. Then the regular level set $\Sigma:=f^{-1}\left(n+\frac{1}{2}\right)$ is a smooth hypersurface which divides $M$ into two connected components $M \backslash \Sigma=Y_{0} \cup Y_{1}$. It follows that $Y_{i}, i=0,1$, deformation retracts (along $\pm \nabla f$ ) to the skeleton $\operatorname{Sk}\left(Y_{i}\right)$, which is a finite $n$-dimensional CW-complex.

Writing $Y$ for either of $Y_{0}$ or $Y_{1}$, we now construct $N(\operatorname{Sk}(Y))$ as a compact contact handlebody. There exists a neighborhood of the 0 -cells of $Y$ that can be
written as a contact handlebody $H_{0}=[-1,1] \times W_{0}$, where $W_{0}$ is Weinstein. Arguing by induction, assume that a neighborhood of the $k$-skeleton of $\mathrm{Sk}(Y)$ can be realized as a contact handlebody $H_{k}=[-1,1] \times W_{k}$, where $W_{k}$ is Weinstein and $\Gamma_{k}=\{0\} \times \partial W_{k}$ is the dividing set of $\partial H_{k}$. We explain how to attach the $(k+1)$-handles to $\partial H_{k}$, where $k+1 \leq n$, in a contact manner. Write $K$ for the core disk of a $(k+1)$-handle. Then $\operatorname{dim} \partial K=k$ and by dimension reasons $\partial K \subset \partial H_{k}$, after possible perturbation, can be isotoped into $\Gamma_{k}$ using the Liouville flow on $W_{k}$. Let $v$ (resp. $w$ ) be a nonvanishing ( $T M$-valued) vector field along $\Gamma_{k}$ that is tangent to $\xi$, transverse to $\partial H_{k}$ (resp. transverse to $\Gamma_{k}$ and tangent to $\partial H_{k}$ ), and is symplectically orthogonal to $\left.\xi\right|_{\Gamma_{k}}$.

Next we would like to isotop $\partial K$ to an isotropic submanifold $\partial K^{\prime}$ in $\Gamma_{k}$ (it may be Legendrian if $k+1=n$ ) and then isotop $K$ to an isotropic submanifold $K^{\prime} \subset Y \backslash \operatorname{int} H_{k}$ with boundary $\partial K^{\prime}$, using Gromov's $h$-principle [Gro86, p. 339] for isotropic submanifolds in a contact manifold. To this end we show that:

Claim 10.0.1. There is a formal homotopy $\phi_{t}: T K \rightarrow T M, t \in[0,1]$, such that:
(1) $\phi_{0}$ is the derivative of the inclusion map;
(2) $\phi_{t}$ is an injective bundle map for all $t \in[0,1]$;
(3) the fibers of $\phi_{1}(T K)$ are $\xi$-isotropic;
(4) the fibers of $\phi_{t}(T K)$ for all $t \in[0,1]$ along $\partial K$ have the form $\mathbb{R}\langle v\rangle$ times a plane of $T \Gamma_{k}$; and
(5) when $t=1$, the planes in (4) are $\xi^{\prime}:=\xi \cap T \Gamma_{k}$-isotropic.

Proof of Claim 10.0.1. We will explain the Legendrian (i.e., $k+1=n$ ) case, which is the hardest case. Since $K$ is a disk, it is clearly formally Legendrian inside its disk neighborhood $N(K)$. The key point is to make $\partial K$ formally isotropic as well.

Let $\tau$ be a trivialization of $\left.\xi\right|_{N(K)}$. Projecting out the Reeb direction and using the trivialization $\tau$, the embedding $K \hookrightarrow N(K)$ can be converted into the map $\phi_{0}^{b}: K \rightarrow G(n, 2 n)$, where $G(n, 2 n)$ is the Grassmannian of $n$-planes in $\mathbb{R}^{2 n}$. Since $K$ is a disk, $\phi_{0}^{b}$ is homotopic to $\phi_{1 / 2}^{b}: K \rightarrow \mathscr{L}_{n}$, where $\mathscr{L}_{n} \subset G(n, 2 n)$ is the Lagrangian Grassmannian. Next, we would like to further homotop $\phi_{1 / 2}^{b}$ to $\phi_{1}^{b}: K \rightarrow \mathscr{L}_{n}$ such that $v(x) \in \phi_{1}^{b}(x)$ for all $x \in \partial K$. At this point we note that over $\partial K$ the trivial complex bundle $\xi$ satisfies

$$
\xi \simeq \mathbb{R}\langle v, w\rangle \oplus \xi^{\prime}
$$

and that the classification of complex vector bundles over $\partial K \simeq S^{n-1}$ is given by $\pi_{n-2} U \simeq 0$ or $\mathbb{Z}$. In the former case, $\xi^{\prime}$ is a trivial complex vector bundle; in the latter case, $\xi^{\prime}$ is classified by its $(n-1) / 2$ nd Chern class, but then $\mathbb{C} \oplus \xi^{\prime}$ is trivial, so the Chern class must vanish and $\xi^{\prime}$ must be trivial. Hence we can view the desired map $\phi_{1}^{b}$ as restricting to $\partial K \rightarrow \mathscr{L}_{n-1} \subset \mathscr{L}_{n}$, corresponding to a standard inclusion $\mathbb{R}^{2 n-2} \hookrightarrow \mathbb{R}^{2 n}$.

We now claim that $\pi_{n-1} \mathscr{L}_{n-1} \rightarrow \pi_{n-1} \mathscr{L}_{n}$ is surjective. Using the fact that $\mathscr{L}_{n}=U(n) / O(n)$, we have:


Using the homotopy exact sequences for $U(n) / U(n-1)=S^{2 n-1}$ and $O(n) / O(n-$ $1)=S^{n-1}$, it follows that $a, c$ are surjective and $d$ is injective. The claim then follows from the five lemma.

The claim then allows us to homotop $\phi_{1 / 2}^{b}$ to $\phi_{1}^{b}$ such that, when we view $\phi_{t}^{b}$, $t \in[0,1]$, as family of maps $\phi_{t}: T K \rightarrow T M, t \in[0,1]$, Conditions (1)-(3) and (5) hold and (4) holds for $t=0,1$. It remains to modify $\phi_{t}$ so that (4) holds for all $t \in[0,1]$ : Observe that $Z:=\sqcup_{t \in[0,1]} \phi_{t}\left(\left.T K\right|_{\partial K}\right)$ has dimension $2 n$. Since $\phi_{t}$ may be taken to be generic, $Z \pitchfork \mathbb{R}\langle w\rangle$ and projecting to $\mathbb{R}\langle v\rangle \oplus T \Gamma_{k}$ along $\mathbb{R}\langle w\rangle$ is an isomorphism for each $\phi_{t}\left(T_{x} K\right), t \in[0,1], x \in \partial K$. Hence (4) is also satisfied and Claim 10.0.1 follows.

Hence a neighborhood of the $(k+1)$-skeleton of $\operatorname{Sk}(Y)$ can be realized as a contact handlebody $H_{k+1}=[-1,1] \times W_{k+1}$ and $N\left(\operatorname{Sk}\left(Y_{0}\right)\right) \cup N\left(\operatorname{Sk}\left(Y_{1}\right)\right)$ can be realized as compact contact handlebodies with sutured convex boundary and its complement in $M$ has sutured concave boundary.

Now identify $M \backslash\left(N\left(\operatorname{Sk}\left(Y_{0}\right)\right) \cup N\left(\operatorname{Sk}\left(Y_{1}\right)\right)\right)$ with $\Sigma \times[0,1]$ such that if we write $\Sigma_{t}:=\Sigma \times\{t\}$, then $\Sigma_{i}=\partial N\left(\operatorname{Sk}\left(Y_{i}\right)\right), i=0,1$, are convex with dividing sets corresponding to the sutures. By Theorem 1.2.5, $\left.\xi\right|_{\Sigma \times[0,1]}$ is given by a finite sequence of bypass attachments, which can be further turned into a sequence of modifications of the trivial Weinstein POBD of

$$
N\left(\operatorname{Sk}\left(Y_{0}\right)\right)=[-1,1] \times W_{0, n}
$$

according to Lemma 8.4.3 (here $W_{0, n}$ is $W_{n}$ for $Y_{0}$ ). In this way we obtain a 1-Weinstein POBD of $M \backslash N\left(\operatorname{Sk}\left(Y_{1}\right)\right)$, viewed as a contact manifold with sutured concave boundary. It remains to fill in $N\left(\operatorname{Sk}\left(Y_{1}\right)\right)$ in the obvious manner to get a compatible 1-Weinstein OBD. (Alternatively, we can attach all the contact $n$-handles of the bypass attachments to $N\left(\operatorname{Sk}\left(Y_{0}\right)\right)$, which gives a generalized contact handlebody, and all the contact $(n+1)$-handles of the bypass attachments to $N\left(\operatorname{Sk}\left(Y_{1}\right)\right)$, which also gives a generalized contact handlebody. Hence we can view the OBD as consisting of two halves and each half is a generalized contact handlebody foliated by 1 -Weinstein domains.)

Next we turn to the relative case, i.e., to contact manifolds with boundary. The following theorem is a more precise version of Corollary 1.3.2 (note the boundary condition of Corollary 1.3.2 was vaguely stated).

Theorem 10.0.2. If $(M, \xi)$ is a compact contact manifold with sutured concave boundary and $R_{ \pm}(\partial M)$ are Weinstein, then there exists a compatible 1 -Weinstein POBD where each of $R_{+}(\partial M)$ and $R_{-}(\partial M)$ extends to a page.

Proof. Choose a generic self-indexing Morse function $f: M \rightarrow\left[\frac{1}{2}, \infty\right)$ such that $f \equiv \frac{1}{2}$ on $\partial M$. In other words, $f$ has no index 0 critical points. As in the absolute case, consider the hypersurface $\Sigma:=f^{-1}\left(n+\frac{1}{2}\right)$ which divides $M$ into two components $Y_{i}, i=0,1$, such that $Y_{0}$ contains all the critical points of index at most $n$ and $Y_{1}$ contains all the critical points of index at least $n+1$. By the handle attachment discussion in the proof of Corollary 1.3.1, we can turn the critical points in $Y_{0}$ into isotropic handles attached to $\partial M$ along the suture which we still denote by $N\left(\mathrm{Sk}\left(Y_{0}\right)\right)$ although it is no longer a contact handlebody, and the critical points in $Y_{1}$ into the handle decomposition of a contact handlebody $N\left(\mathrm{Sk}\left(Y_{1}\right)\right)$ with sutured convex boundary, as in the closed case. The rest of the proof proceeds as in the closed case.

Proof of Corollary 1.3.3. We modify the proof of Corollary 1.3.1 as follows: In the decomposition of $M$ into contact handlebodies $Y_{0}$ and $Y_{1}$, we first take a standard contact neighborhood $N(\Lambda)$ of the Legendrian $\Lambda$ and a corresponding contact Morse function $f$ on $N(\Lambda)$ that is constant on $\partial N(\Lambda)$. We then extend $f$ arbitrarily to a Morse function which is self-indexing on $M \backslash N(\Lambda)$, realize the $k$-handles with $k \leq n$ as contact handles, and attach them to $N(\Lambda)$ to obtain $Y_{0}$. The rest of the proof is the same as that of Corollary 1.3.1.

## 11. APPLICATIONS TO CONTACT SUBMANIFOLDS

In this section, we apply the techniques developed so far to prove Corollaries 1.3.6, 1.3.7, and 1.3.9.

Recall that Ibort, Martínez-Torres, and Presas [IMTP00] constructed contact submanifolds $Y$ of $(M, \xi)$ as the zero loci of "approximately holomorphic" sections of a complex line or vector bundle over $M$. Our strategy for constructing contact submanifolds is rather different and relies on the following key observation: if $\Sigma \subset M$ is a hypersurface which contains a codimension 2 submanifold $Y$ such that the characteristic foliation $\Sigma_{\xi}$ is transverse to $Y$, then $Y \subset(M, \xi)$ is a contact submanifold. ${ }^{4}$
11.1. Models for pushing across mushrooms. We first explain the idea of the proofs of Corollaries 1.3 .6 and 1.3 .7 in the $n=1$ case. We apply the operation of "pushing across a mushroom", as given in Figure 11.1.1. (Also refer to Figure 3.2.1.) There are three variants of this operation: The initial step is given by Figure 11.1.1(a), where the curve $Y$ (in purple) is pushed across the shaded region (in light blue). The result is a curve $Y^{\prime}$ which intersects the characteristic foliation $Z_{\xi}$ of the mushroom $Z$ with opposite sign on the blue interval between the dots. (The purple portion represents where $Y^{\prime}$ intersects $Z_{\xi}$ with the sign that we want to reverse and the blue portion represents where $Y^{\prime}$ intersects $Z_{\xi}$ with the desired sign.) Now assuming that there is another mushroom placed just to the right of the first one (i.e., with smaller $t$-coordinate) but with larger $s$-coordinate, we can push $Y^{\prime}$ across that mushroom (or more precisely the shaded region given in Figure 11.1.1(b)) to obtain $Y^{\prime \prime}$ with a "larger" blue portion. Finally, after several steps

[^3]of type (b), we complete the "sign reversal" by pushing across the shaded region given by Figure 11.1.1(c).


Figure 11.1.1. Pushing across mushrooms.

We now discuss the general case. We start with the ambient contact manifold

$$
\left(M=\mathbb{R}_{z, s, t}^{3} \times W, \xi=\operatorname{ker}\left(d z+e^{s}(d t+\lambda)\right)\right)
$$

as before, where $(W, \lambda)$ is a Weinstein manifold and the hypersurface before perturbation is $\Sigma=\{z=0\}$. Let $H=\left[0, t_{0}\right] \times W_{\tau_{0}}^{c}$ (refer to Eq. (5.1.1)) be a contact handlebody, $Z$ a mushroom with contact handlebody profile $H$, and $H_{\text {in }}$ the subset of $\mathbb{R}_{t} \times W$ given by (5.5.2) and satisfying Proposition 5.5.2. By Definition 5.5.1, $Z$ and $\Sigma$ agree outside of a small neighborhood $N\left(\left[0, s_{0}\right] \times\left[0, t_{0}\right] \times W_{\tau_{0}}^{c}\right)$ of $\left[0, s_{0}\right] \times\left[0, t_{0}\right] \times W_{\tau_{0}}^{c}$ whose size depends on the smoothing parameters.

By taking an isotopy $\mathcal{I}$ from $Z$ to $\Sigma$, which is supported outside of

$$
\Sigma \backslash N\left(\left[0, s_{0}\right] \times\left[0, t_{0}\right] \times W_{\tau_{0}}^{c}\right)
$$

and assumed to satisfy (I1)-(I3) below, and carrying the germ of the contact structure with it, we will assume in this section that $Z=\{z=0\}=\mathbb{R}_{s} \times \mathbb{R}_{t} \times W$. The isotoped contact structure will still be called $\xi$.

Let $\rho=\rho_{W}: W \backslash \operatorname{Sk}(W) \rightarrow \mathbb{R}$ be the function such that $\rho\left(\partial W^{c}\right)=0$ and $d \rho\left(X_{\lambda}\right)=1$, and let $\tilde{\tau}=\tilde{\tau}_{W}: W \rightarrow \mathbb{R}_{\geq 0}$, where $\tilde{\tau}=e^{\rho}$ on $W \backslash \operatorname{Sk}(W)$ and $\tilde{\tau}=0$ on $\operatorname{Sk}(W)$. The function $\tilde{\tau}$ is a variant of $\tau$ which appears in Section 5.4. Consider the map

$$
\begin{gather*}
\Phi=\Phi_{W}: Z=\mathbb{R}_{s} \times \mathbb{R}_{t} \times W \rightarrow \mathbb{R}_{\geq 0} \times \mathbb{R}^{2}  \tag{11.1.1}\\
(s, t, x) \mapsto(\tilde{\tau}(x), s, t)
\end{gather*}
$$

In view of Eq. (5.2.1) and Eq. (5.4.1), we may choose the isotopy $\mathcal{I}$ from $Z$ to $\Sigma$ so that $Z_{\xi}=Z_{\xi}^{A}+Z_{\xi}^{B}$, where:
(I1) $Z_{\xi}^{A}(s, t, x)$ is the component in the $\mathbb{R}_{s, t}^{2}$-direction and depends only on $(\tilde{\tau}(x), s, t)$;
(I2) $Z_{\xi}^{A}(s, t, x)=Z_{\xi}^{3}(s, t)$ for $\tilde{\tau}(x) \leq 1 / 2$, where $Z_{\xi}^{3}$ is the characteristic foliation of the 3 -dimensional mushroom; and
(I3) $Z_{\xi}^{B}(s, t, x)$ has the form $f(\tilde{\tau}(x), s, t) X_{\lambda}(x)+g(s, t, x) R_{\eta}$ away from $\mathbb{R}^{2} \times$ $\operatorname{Sk}(W)$, where $R_{\eta}$ is the Reeb vector field of $\left.\lambda\right|_{\partial W^{c}}$.

Claim 11.1.1. Away from $\mathbb{R}^{2} \times \operatorname{Sk}(W), \Phi_{*}\left(Z_{\xi}\right)$ is a well-defined (i.e., singlevalued) vector field on $\mathbb{R}_{>0} \times \mathbb{R}^{2}$.

Proof. This follows immediately from (I1)-(I3).
Given a codimension 0 submanifold $S$ of $\{z=s=0\}$, let $U(S) \subset \Sigma$ be the set of points $p$ for which there is a forward smooth flow line of $Z_{\xi}$ from some $q \in S$ to $p$ and let $\bar{U}(S)$ be its closure. We orient $S$ using the normal orientation $\partial_{s}$ which agrees with $Z_{\xi}$.

Remark 11.1.2. For our submanifolds $S$ of interest, Claim 11.1.1 reduces the problem of determining $\bar{U}(S)$, pushing $S$ across $\bar{U}(S)$ to $S^{\prime}$, and determining the set of points where $S^{\prime}$ is positively/negative transverse to $Z_{\xi}$, to a 3-dimensional one on $\mathbb{R}_{\geq 0} \times \mathbb{R}^{2}$.

Recall that the zeros of $Z_{\xi}$ are $e_{-}^{x}, e_{+}^{x}, h_{-}^{x}, h_{+}^{x}$, where $e_{-}, e_{+}, h_{-}, h_{+}$are zeros of the 3 -dimensional mushroom $Z^{3}$ and $\lambda(x)=0$. They are all contained in $Z^{\prime}=Z^{3} \times W^{c}$; see $\S 5.3$ and Figure 3.2.1.

Lemma 11.1.3. $\bar{U}\left(H_{\mathrm{in}}\right)$ is a compact manifold diffeomorphic to $H_{\mathrm{in}} \times[0,1]$, after smoothing $Z_{\xi}$ on a sufficiently small neighborhood of $E_{-}:=\left\{e_{-}\right\} \times \operatorname{Sk}(W)$. Here $H_{\text {in }}=H_{\text {in }} \times\{0\}$.

The smoothing is generally necessary even in the $n=1$ case, as the shaded region of Figure 11.1.1(a) will usually have a corner at $e_{-}$. Note that $E_{-}$is the closure of the union of all the unstable trajectories of $e_{-}^{x}$ as we range over all the $x$ satisfying $\lambda(x)=0$.

Proof. This is immediate from Proposition 5.5.2.
To prepare for the next lemma let us identify $\Phi\left(\bar{U}\left(H_{\text {in }}\right)\right)$ with $V \times[0,1]$, where $V=\Phi\left(H_{\text {in }} \times\{0\}\right)$ is a half-disk $\left\{z \in \mathbb{C}||z| \leq 1, \operatorname{Im} z \geq 0\}\right.$. Let $V^{\delta}$ be the slight enlargement $\{z \in \mathbb{C}||z| \leq 1+\delta, \operatorname{Im} z \geq 0\}$ of $V$ with $\delta>0$. Then the back face $V^{\delta} \times\{1\}$ and the bottom face $\left(V^{\delta} \times[0,1]\right) \cap\{\tilde{\tau}=0\}$ are as shown in Figure 11.1.2.

The following lemma explains how to modify three types of $S$. Figures 11.1.2(a), (b), (c) correspond to (a), (b), (c) in the lemma.

## Lemma 11.1.4.

(a) If $S$ is a slight enlargement of $H_{\text {in }}$ such that $S \backslash H_{\text {in }}$ is a collar neighborhood of $S$, then there is a submanifold $S^{\prime}$ obtained by isotoping $S$ across $\bar{U}\left(H_{\mathrm{in}}\right)$ rel $\partial S$ in $Z$ such that $S^{\prime}$ decomposes into $S_{+}^{\prime} \cup S_{-}^{\prime}$ and:
(i) $Z_{\xi}$ is positively (resp. negatively) transverse to $S^{\prime}$ along $\operatorname{int}\left(S_{+}^{\prime}\right)$ (resp. $\operatorname{int}\left(S_{-}^{\prime}\right)$ ).
(ii) $S_{+}^{\prime}$ is a pushoff of $\left(\partial H_{\mathrm{in}}\right) \times[0,1]$ and $S_{-}^{\prime}$ is a pushoff of $H_{\mathrm{in}} \times\{1\}$, using an outward-pointing vector field to $H_{\mathrm{in}} \times[0,1]$ after rounding the corner along $\partial H_{\mathrm{in}} \times\{1\}$.
(iii) Along the folding locus $S_{+}^{\prime} \cap S_{-}^{\prime}, Z_{\xi}$ points into $S_{+}^{\prime}$.


Figure 11.1.2. The image $V \times[0,1]=\Phi\left(\bar{U}\left(H_{\text {in }}\right)\right)$ of the sink $\bar{U}\left(H_{\text {in }}\right)$, after smoothing. Shown are the back face $V^{\delta} \times\{1\}$ and the bottom face $\left(V^{\delta} \times[0,1]\right) \cap\{\tilde{\tau}=0\}$. The black dots represent the images of $h_{+}, e_{-}, h_{+}$times $\operatorname{Sk}(W)$. The purple and blue arcs are the intersections of $\Phi\left(S^{\prime}\right)$ with the back and bottom faces and are subsets of $\Phi\left(S_{+}^{\prime}\right)$ and $\Phi\left(S_{-}^{\prime}\right)$, respectively. Figure 11.1.2(a), (b), (c) represent the higher-dimensional analogs of Figure 11.1.1(a), (b), (c).
(b) Let $W_{+}^{\prime}$ (resp. $W_{-}^{\prime}$ ) be a possibly empty Weinstein subdomain of $W^{c} \times\left\{t_{0}\right\}$ (resp. $W^{c} \times\{0\}$ ). If $S$ is obtained from $H_{\mathrm{in}}$ by perturbing $\partial H_{\mathrm{in}}$ to $\partial S$ so that $\partial S \pitchfork \partial H_{\text {in }}$ and either $(\alpha) \partial S \backslash H_{\text {in }}$ is a submanifold parallel to $W_{+}^{\prime}$ or ( $\beta$ ) $S \cap \partial H_{\mathrm{in}}=W_{-}^{\prime}$, then there is a submanifold $S^{\prime}$ obtained by isotoping $S$ across $\bar{U}\left(S \cap H_{\mathrm{in}}\right)$ rel $\partial S$ in $Z$ such that $S^{\prime}$ decomposes into $S_{+}^{\prime} \cup S_{-}^{\prime}$ and:
(i) $Z_{\xi}$ is positively (resp. negatively) transverse to $S^{\prime}$ along $\operatorname{int}\left(S_{+}^{\prime}\right)$ (resp. $\operatorname{int}\left(S_{-}^{\prime}\right)$ ).
(ii) Let $\partial_{-} \bar{U}$ be the closure of the subset of $\partial \bar{U}\left(S \cap H_{\text {in }}\right)$ for which there is a smooth flow line from some $y \in S \cap \partial H_{\text {in }}$ and let $\partial_{+} \bar{U}$ be the closure of $\partial \bar{U}\left(S \cap H_{\mathrm{in}}\right) \backslash \partial_{-} \bar{U}$. Then $S_{+}^{\prime}$ (resp. $S_{-}^{\prime}$ ) is a pushoff of $\partial_{+} \bar{U}\left(\right.$ resp. $\left.\partial_{-} \bar{U}\right)$.
(iii) Along the folding locus $S_{+}^{\prime} \cap S_{-}^{\prime}, Z_{\xi}$ points into $S_{+}^{\prime}$.
(c) If $S$ is a slight retraction of $H_{\text {in }}$ such that $H_{\text {in }} \backslash S$ is a collar neighborhood of $H_{\mathrm{in}}$, then there is a submanifold $S^{\prime}$ obtained by isotoping $S$ across $\bar{U}\left(S^{\text {sh }}\right)$ rel $\partial S$ in $Z$ such that $Z_{\xi}$ is negatively transverse to $S^{\prime}$. Here $S^{\text {sh }}$ is a slight retraction of $S$.

The smoothing of $Z_{\xi}$ can be done inside a sufficiently small neighborhood $N\left(E_{-}\right)$of $E_{-}$, and hence the resulting $S^{\prime}$ can be made to avoid $N\left(E_{-}\right)$. Also note that in Figure 11.1.1 $S_{+}^{\prime}$ (resp. $S_{-}^{\prime}$ ) corresponds to the purple (resp. blue) portion of $Y^{\prime}$.

Remark 11.1.5. The directions of the folding locus given by (a)(iii) and (b)(iii) are essential when iterating the moves given by (a)-(c).
Proof. (a) The submanifold $S^{\prime}$ is constructed using the map $\Phi$ as follows: Let $\mathscr{S}$ be the union of the side $\mathscr{S}_{+}=\partial \Phi(S) \times\left[0,1+\epsilon^{\prime}\right]$ and the back $\mathscr{S}_{-}=\Phi(S) \times\left\{1+\epsilon^{\prime}\right\}$
of $\partial\left(\Phi(S) \times\left[0,1+\epsilon^{\prime}\right]\right)$, with $\epsilon^{\prime}>0$ small. Then $S^{\prime}$ is the preimage of the set $\Phi\left(S^{\prime}\right)$ obtained by rounding the corner and perturbing $\mathscr{S}$ so that:

- $\Phi\left(S^{\prime}\right) \cap\{\tilde{\tau} \in[0, \epsilon]\}$ with $\epsilon>0$ small agrees with $Y^{\prime} \times\{\tilde{\tau} \in[0, \epsilon]\}$, where $Y^{\prime}$ is as given in Figure 11.1.1(a).
- $\Phi\left(S_{+}^{\prime}\right)$ is a perturbation of $\mathscr{S}_{+}$, is positively transverse to $\Phi_{*}\left(Z_{\xi}\right)$, and intersects the back and bottom faces in the purple arcs given in Figure 11.1.2(a).
- $\Phi\left(S_{-}^{\prime}\right)$ is a perturbation of $\mathscr{S}_{-}$and is negatively transverse to $\Phi_{*}\left(Z_{\xi}\right)$.
(i)-(iii) are immediate from the construction.
(c) is similar to (a). In this case $\mathscr{S}$ is the union of the side $\partial \Phi(S) \times\left[0,1+\epsilon^{\prime}\right]$ and the back $\Phi(S) \times\left\{1+\epsilon^{\prime}\right\}$ of $\partial\left(\Phi(S) \times\left[0,1+\epsilon^{\prime}\right]\right)$, with $\epsilon^{\prime}>0$ small, and $\Phi\left(S^{\prime}\right)$ satisfies:
- $\Phi\left(S^{\prime}\right) \cap\{\tilde{\tau} \in[0, \epsilon]\}$ with $\epsilon>0$ small agrees with $Y^{\prime} \times\{\tilde{\tau} \in[0, \epsilon]\}$, where $Y^{\prime}$ is as given in Figure 11.1.1(c).
- $\Phi\left(S^{\prime}\right)$ is negatively transverse to $\Phi_{*}\left(Z_{\xi}\right)$ and intersects the back and bottom faces in the purple arcs given in Figure 11.1.2(c).
(b) is similar to (a) and (c) in the case where ( $\beta$ ) holds and $W_{-}^{\prime} \subset W^{c}$ is a retraction of $W^{c}$ along the Liouville flow. In this case $\mathscr{S}$ is again the union of the side $\partial \Phi(S) \times\left[0,1+\epsilon^{\prime}\right]$ and the back $\Phi(S) \times\left\{1+\epsilon^{\prime}\right\}$ of $\partial\left(\Phi(S) \times\left[0,1+\epsilon^{\prime}\right]\right)$, with $\epsilon^{\prime}>0$ small. $\mathscr{S}$ decomposes into $\mathscr{S}_{+}$and $\mathscr{S}_{-}$along the seam consisting of the blue dot on the back face of Figure 11.1.2(b) times $\left[0,1+\epsilon^{\prime}\right]$, followed by the extension along $\partial \Phi(S) \times\left\{1+\epsilon^{\prime}\right\}$ to the point with $\tilde{\tau}=0$ and $t \leq 0$ (which is analogous to the blue dot on $Y^{\prime \prime}$ in Figure 11.1.1(b)). (i)-(iii) are immediate from the construction.

For $(\beta)$ in general, i.e., when $W_{-}^{\prime} \subset W^{c}$ is an arbitrary Weinstein subdomain, we first write $S$ as the union of $S_{1}$ of type (c) and $S_{2}$ a subset of

$$
\{s=0\} \times\left\{t \in\left[-\epsilon^{\prime \prime}, \epsilon^{\prime \prime}\right]\right\} \times W_{-}^{\prime},
$$

such that $S_{1} \cup S_{2}$ has smooth boundary and $\epsilon^{\prime \prime}>0$ is small. (In particular, $S_{2}$ has a rounded corner at $\left\{-\epsilon^{\prime \prime}\right\} \times \partial W_{-}^{\prime}$ and a cusp corner at $\left\{\epsilon^{\prime \prime}\right\} \times \partial W_{-}^{\prime}$.) We then apply (c) to push $S$ across $S_{1}$. Next, let $\rho_{W_{-}^{\prime}}: W_{-}^{\prime} \backslash \operatorname{Sk}\left(W_{-}^{\prime}\right) \rightarrow \mathbb{R}$ be the function such that $\rho_{W_{-}^{\prime}}\left(\partial W_{-}^{\prime}\right)=0$ and $d \rho_{W_{-}^{\prime}}\left(X_{\lambda}\right)=1$, and $\tilde{\tau}_{W_{-}^{\prime}}: W_{-}^{\prime} \rightarrow \mathbb{R}_{\geq 0}$ satisfy $\tilde{\tau}_{W_{-}^{\prime}}=e^{\rho_{W_{-}^{\prime}}}$ on $W_{-}^{\prime} \backslash \operatorname{Sk}\left(W_{-}^{\prime}\right)$ and $\tilde{\tau}_{W_{-}^{\prime}}=0$ on $\operatorname{Sk}\left(W_{-}^{\prime}\right)$, and

$$
\begin{aligned}
& \Phi_{W_{-}^{\prime}}: \mathbb{R} \times\left[-\epsilon^{\prime \prime}, \epsilon^{\prime \prime}\right] \times W_{-}^{\prime} \rightarrow \mathbb{R}_{\geq 0} \times \mathbb{R} \times\left[-\epsilon^{\prime \prime}, \epsilon^{\prime \prime}\right], \\
&(s, t, x) \mapsto\left(\tilde{\tau}_{W_{-}^{\prime}}, s, t\right),
\end{aligned}
$$

as in Eq. (11.1.1). We then push $S_{2}$ - viewed as a subset of $S$ from the previous paragraph - across using $\Phi_{W^{\prime}}$. The seam with respect to the vector field $\left(\Phi_{W_{-}^{\prime}}\right)_{*}\left(\left.Z_{\xi}\right|_{\mathbb{R} \times\left[-\epsilon^{\prime \prime}, \epsilon^{\prime \prime}\right] \times W_{-}^{\prime}}\right)$ is analogous to that of the previous paragraph, where the blue dot on the back face of Figure 11.1.2(b) times $\left[0,1+\epsilon^{\prime}\right]$ can be viewed as the rounded corner near $\Phi_{W_{-}^{\prime}}\left(\left\{\left(0,-\epsilon^{\prime \prime}\right)\right\} \times \partial W_{-}^{\prime}\right) \times\left[0,1+\epsilon^{\prime}\right]$.

The proof of $(\alpha)$ follows the same general reasoning and is omitted.
11.2. Existence $h$-principles. We use the models from $\S 11.1$ (and specifically Lemma 11.1.4(a), (b), (c)) to prove Corollaries 1.3.6 and 1.3.7 as well as the existence $h$-principle for contact structures in this subsection.

Let $Y$ be a closed codimension 2 submanifold of $(M, \xi)$. For the moment assume that $Y \subset(M, \xi)$ has a trivial normal bundle. Any even-codimensional submanifold with trivial normal bundle is almost contact by [BCS14, Lemma 2.17].
Proof of Corollary 1.3.7 assuming trivial normal bundle. Let $Y \subset(M, \xi)$ be a contact submanifold. By the trivial normal bundle condition there exists a hypersurface $\Sigma:=Y \times[-1,1]_{s} \subset M$ such that $Y=Y_{0}$ and $\Sigma_{\xi}=\partial_{s}$, where $Y_{s}:=Y \times\{s\}$.

By Corollary 1.3.1, $Y$ admits a strongly Weinstein $\operatorname{OBD}(B, \pi)$ and an adapted contact form $\alpha$. Referring to Appendix B, by Lemma B.2.1, Lemma B.3.1 and Proposition B.0.1, for $\delta>0$ small there exists a strongly Weinstein and damped OBD $\left(B^{\prime}, \pi^{\prime}\right)$ of $\left(M, \alpha^{\prime}\right)$, where $\alpha^{\prime}$ is $C^{1}$-close to $\alpha$ with $\alpha^{\prime}$ strongly adapted and damped and $A\left(B^{\prime}, \pi^{\prime}, \alpha^{\prime}\right)<\delta$.

Fix a large integer $N>0$. For $\delta>0$ small, we can cover $Y$ by a finite number $H_{1}, \ldots, H_{N}$ of contact handlebodies over the Weinstein pages (arranged in order of increasing $\theta$ ) of total thickness at most $\delta$. We then install mushrooms with profiles $H_{1}, \ldots, H_{N}$, arranged as in $\S 7.3$ and Figure 7.3.1, on a small neighborhood of $Y \times\left\{\frac{1}{2}\right\}$. The mushroom corresponding to $H_{N}$ will be of type (a) from Lemma 11.1.4, those corresponding to $H_{N-1}, \ldots, H_{2}$ will be of type (b), and the last one corresponding to $H_{1}$ will be of type (c). Pushing $Y$ across the mushrooms and applying Lemma 11.1.4 in order yields a contact submanifold with orientation reversed.

Proof of Corollary 1.3.6 assuming trivial normal bundle. As before, by the trivial normal bundle condition there exists a hypersurface $\Sigma:=Y \times[-1,1]_{s} \subset(M, \xi)$ such that $Y=Y \times\{0\}$, but now $\Sigma_{\xi}$ is not necessarily transverse to $Y$.

By Proposition A.3.2, we can assume that $Y$ is $\Sigma_{\xi}$-folded with folding locus $C \subset Y$. We then have the decomposition $Y=Y_{+} \cup_{C} Y_{-}$such that $\Sigma_{\xi}$ is positively transverse to $\operatorname{int}\left(Y_{+}\right)$and negatively transverse to $\operatorname{int}\left(Y_{-}\right)$. Using Theorem 1.2.3, we perturb $C$ into a convex hypersurface. We then decompose $C=C_{+} \sqcup C_{-}$, where $C_{+}$(resp. $C_{-}$) is the set of points where $\Sigma_{\xi}$ points out of (resp. into) $Y_{-}$.

Next we perturb $Y$ on a neighborhood of $C$ and relative to $C$ and convert $C$ from a fold to a cusp, keeping the same notation $C, Y_{+}, Y_{-}$. For more details see Appendix A. Note that $Y_{-}$is the "lower sheet" near the cusp $C_{+}$and the "upper sheet" near the cusp $C_{-}$.

Since we can deal with the connected components of $Y_{-}$one at a time, assume $Y_{-}$is connected.
(i) Suppose for the moment that $\partial Y_{-}=C_{+}$or $\partial Y_{-}=C_{-}$. Assuming $\partial Y_{-}=$ $C_{+}$(the case $\partial Y_{-}=C_{-}$is similar), there exist a collar neighborhood $N_{\epsilon}\left(Y_{-}\right)=$ $Y_{-} \times[-\epsilon, 0]_{s}$ of $Y_{-}=Y_{-} \times\{0\}$ and a piecewise smooth approximation $Y^{\star}$ of $Y$ such that:

- $\left.\Sigma_{\xi}\right|_{N_{\epsilon}\left(Y_{-}\right)}=\partial_{s} ;$
- $Y^{\star} \cap N_{\epsilon}\left(Y_{-}\right)=\left(Y_{-} \times\left\{-\frac{\epsilon}{2}\right\}\right) \cup\left(C_{+} \times\left[-\frac{\epsilon}{2}, 0\right]\right) \cup\left(N_{\epsilon}\left(C_{+}\right) \times\{0\}\right)$, where $N_{\epsilon}\left(C_{+}\right) \subset Y_{-}$is a collar neighborhood of the boundary; and
- $Y^{\star} \backslash N_{\epsilon}\left(Y_{-}\right)$is positively transverse to $\Sigma_{\xi}$.

Informally, $Y^{\star}$ is sutured with respect to $\Sigma_{\xi}$, where the suture is $C_{+} \times\left[-\frac{\epsilon}{2}, 0\right]$.
We now apply the proof of Corollary 1.3.7 above, with POBDs in place of OBDs. By Theorem 10.0.2 $Y^{\star}$ admits a compatible POBD. In particular we have the following data:
(1) a contact 1 -form $\alpha$ on $Y^{\star}$, the binding $B$, and Weinstein pages $S_{\theta}, \theta \in$ $[0,2 \pi]$;
(2) an identification $\sim$ of Weinstein cobordisms $S_{0}^{\prime} \subset S_{0}$ and $S_{2 \pi}^{\prime} \subset S_{2 \pi}$ such that $\partial S_{0}^{\prime} \supset \partial S_{0}$ and $\partial S_{2 \pi}^{\prime} \supset \partial S_{2 \pi}$; and
(3) an identification of $Y^{\star} \backslash B$ with $\sqcup_{\theta \in[0,2 \pi]} S_{\theta} / \sim$.

Applying the relative version of Proposition B.0.1, for $\delta>0$ small there exists a strongly Weinstein POBD of $Y^{\star}$, where the new contact form $\alpha^{\prime}$ is $C^{1}$-close to the old one, $\alpha^{\prime}$ is strongly adapted and damped, and $\sup _{\gamma \in \mathscr{R}\left(\alpha^{\prime}\right)} A(\gamma)<\delta$ (defined as in Definition B.1.1). We can cover $Y$ by a finite number $H_{1}, \ldots, H_{N}$ of contact handlebodies over the new Weinstein pages $S_{\theta}^{\prime}$ (arranged in order of increasing $\theta$ ) of total thickness at most $\delta$; here $H_{1}$ is a contact handlebody over $S_{0}^{\prime}$ and $H_{N}$ is a contact handlebody over $S_{2 \pi}^{\prime}$. We then install mushrooms with profiles $H_{1}, \ldots, H_{N}$ as before. The mushrooms corresponding to $H_{N}, \ldots, H_{2}$ will be of type (b) from Lemma 11.1.4 and the last one corresponding to $H_{1}$ will be of type (c). We then push $Y^{\star}$ across the mushrooms and replace $Y^{\star} \cap N_{\epsilon}\left(Y_{-}\right)$ by a hypersurface which is positively transverse to $\Sigma_{\xi}^{\vee}$ and has the same boundary.
(ii) Next suppose that $\partial Y_{-}$has components in $C_{+}$and $C_{-}$. In this case we perturb $Y$ to create extra folds (or equivalently cusps) near $C_{-}$as in Figure 11.2.1 so that the new $Y$ is $Y_{-}^{\prime} \cup A_{+} \cup A_{-} \cup Y_{+}^{\prime}$, where $A_{ \pm} \simeq C_{-} \times[-1,1]$.


Figure 11.2.1. Modification near $C_{-}$. The actual folds are obtained by multiplying by $C_{-}$and the dots represent the folding loci. The arrows indicate $\Sigma_{\xi}$.

Note that after the modification each of the negative regions $Y_{-}^{\prime}$ and $A_{-}$satisfies the conditions of (i). Hence we can now apply (i) to conclude the proof.

Completion of proofs of Corollaries 1.3.6 and 1.3.7. We explain how to remove the requirement that $Y \subset M$ have trivial normal bundle. Let $D \subset Y$ be a closed ball around some point in $Y$. Then $Y \backslash D$ is an open manifold. If $Y$ is an almost contact manifold, then so is $Y \backslash D$. By Gromov's $h$-principles for open contact manifolds (see [EM02, 10.3.2]) and open isocontact embeddings (see [EM02, 12.3.1]), we may assume that $Y \backslash D \subset(M, \xi)$ is an open contact submanifold. Hence the problem reduces to an extension problem over the ball $D$, which clearly has a trivial
normal bundle in $M$. Observe that the proof of Corollary 1.3.6 is essentially a relative extension problem. Corollary 1.3.6 then immediately implies Corollary 1.3.7 since if $Y$ is contact then $-Y$ is almost contact (just reverse the orientation of the line field transverse to $T Y \cap \xi$ ).

Proof of Corollary 1.3.9. We refer the reader to Uebele [Ueb16, Section 4.2] and the references therein for a useful summary of almost contact structures and their homotopy classes.

Let $M^{2 n-1}$ be a closed manifold with an almost contact structure $\zeta=\left(H_{\zeta}, J_{\zeta}\right)$, where $H_{\zeta}$ is a hyperplane field and $J_{\zeta}$ is an almost complex structure on $H_{\zeta}$. We view $\zeta$ as a section $\mathfrak{s}_{\zeta}$ of the fiber bundle $P \rightarrow M$ with fiber $S O(2 n-1) / U(n-1)$ associated to $T M$. Let $D \subset M$ be a closed ball in $M$. By Gromov's $h$-principle (see [EM02]) there exists a homotopy of $\mathfrak{s} \zeta_{\zeta}$ to $\mathfrak{s}_{\zeta^{\prime}}$ (as sections of $P \rightarrow M$ ) such that $\left.\zeta^{\prime}\right|_{M \backslash i n t(D)}$ is a genuine contact structure.

Since $T D$ is trivial, we may view $\left.\mathfrak{s}_{\zeta^{\prime}}\right|_{S^{2 n-2}=\partial D}$ as a map

$$
S^{2 n-2} \rightarrow S O(2 n-1) / U(n-1)
$$

the map is homotopically trivial since it extends over $D$. The set of homotopy classes of extensions of $\left.\mathfrak{s}_{\zeta^{\prime}}\right|_{S^{2 n-2}}$ to $D$ is given by $\pi_{2 n-1}(S O(2 n-1) / U(n-1))$. Now

$$
S O(2 n-1) / U(n-1) \simeq S O(2 n) / U(n)
$$

(see [Ueb16, Section 4.2]), so their homotopy groups agree.
Next we stabilize $\left.\zeta^{\prime}\right|_{D}=\left.\left(H_{\zeta^{\prime}}, J_{\zeta^{\prime}}\right)\right|_{D}$ to the almost contact structure

$$
\eta=\left.\left(H_{\zeta^{\prime}} \times \mathbb{R}^{2}, J_{\zeta^{\prime}} \times j\right)\right|_{D \times \mathbb{R}^{2}}
$$

where $j$ is the standard almost complex structure on $\mathbb{R}^{2}$. Again by Gromov's $h$ principle there exists a formal homotopy on $D \times \mathbb{R}^{2}$ from $\eta$ to a genuine contact structure $\eta^{\prime}$ relative to $N(\partial D) \times\{0\}$, where $N(\partial D) \subset D$ is a neighborhood of $\partial D$. We then apply Corollary 1.3 .6 to obtain a contact submanifold of $\left(D \times \mathbb{R}^{2}, \eta^{\prime}\right)$ which is $C^{0}$-close to $D \times\{0\}$, isotopic to $D \times\{0\}$, and agrees with $D \times\{0\}$ on $N(\partial D) \times\{0\}$. The induced contact structure $\xi$ on $D$ is stably homotopic to $\left.\zeta^{\prime}\right|_{D}$ relative to $\partial D$.

It remains to find contact structures in each homotopy class of almost contact structures on $S^{2 n-1}$ to connect sum with so that we can adjust the homotopy class of the almost contact structure.

When $n \equiv 1$ or $3 \bmod 4$ (i.e., the manifold has dimension $4 k+1$ where $k \in$ $\mathbb{Z}^{+}$), we claim that each homotopy class can be represented by connected sums of canonical contact structures $\xi_{a_{0}, \ldots, a_{n}}$ on Brieskorn spheres $\Sigma\left(a_{0}, \ldots, a_{n}\right)$ that are diffeomorphic to standard spheres. In our case we have $\pi_{2 n-1}(S O(2 n) / U(n)) \simeq$ $\mathbb{Z} / d$ by Massey [Mas61]. By Morita [Mor75] there is a map $a c$ that assigns the value

$$
a c\left(\xi_{a_{0}, \ldots, a_{n}}\right)=\frac{1}{2} \prod_{j=1}^{n}\left(a_{j}-1\right) \in \mathbb{Z} / d
$$

In particular, we can take $\left(a_{0}, \ldots, a_{n}\right)=(7,2, \ldots, 2)$ or $(9,2, \ldots, 2)$ (these were studied by Ustilovsky [Ust99b] and are diffeomorphic to the standard spheres).

Since

$$
a c\left(\xi_{(7,2, \ldots, 2)}\right)=3 \quad \text { and } \quad a c\left(\xi_{(9,2, \ldots, 2)}\right)=4
$$

are relatively prime they generate $\mathbb{Z} / d$.

## Appendix A. Wrinkled and folded embeddings

The technique of wrinkled maps and wrinkled embeddings, developed by Eliashberg and Mishachev in the series of papers [EM97, EM98, EM00, EM09], is extremely powerful in dealing with homotopy problems of smooth maps between manifolds. The goal of this appendix is to give a brief overview of their theory and prove a technical result, Proposition A.3.2, which is only used in Section 11.

This appendix is organized as follows. First we review several fundamental definitions and results in the theory of wrinkled embeddings following Eliashberg and Mishachev. Then we use the wrinkling technique to put a generic hypersurface in a "good" position with respect to a nonvanishing vector field.
A.1. Wrinkled and cuspidal embeddings. In this subsection we review the main results of [EM09].

Let $f: \Sigma \rightarrow M$ be a smooth map between smooth manifolds. In this subsection we assume that $\operatorname{dim} \Sigma=\operatorname{dim} M-1=k$, unless otherwise specified. All the results in this subsection are also valid whenever $\operatorname{dim} \Sigma<\operatorname{dim} M$. If $\operatorname{dim} \Sigma \leq$ $\operatorname{dim} M$, then the singular set $\operatorname{Sing}(f)$ of $f: \Sigma \rightarrow M$ is the set of points in $\Sigma$ where $d f$ is not injective.

Definition A.1.1 (Wrinkled embedding). A smooth map $f: \Sigma \rightarrow M$ is $a$ wrinkled embedding if:
(WE1) $f$ is a topological embedding.
(WE2) $\operatorname{Sing}(f)$ is diffeomorphic to a disjoint union of spheres $S_{i} \cong S^{k-1}$, each of which bounds a $k$-disk in $\Sigma$. Each such $S_{i}$ is called a wrinkle of $f$.
(WE3) The map $f$ can be written locally near each wrinkle as:

$$
\begin{aligned}
& \mathscr{O} p_{\mathbb{R}^{k}}\left(S^{k-1}\right) \rightarrow \mathbb{R}^{k+1}, \\
&(y, z) \mapsto\left(y, z^{3}+3\left(|y|^{2}-1\right) z, \int_{0}^{z}\left(z^{2}+|y|^{2}-1\right)^{2} d z\right),
\end{aligned}
$$

where $(y, z)=\left(y_{1}, \ldots, y_{k-1}, z\right)$ denotes the Cartesian coordinates on $\mathbb{R}^{k}$ such that $S^{k-1}=\left\{|y|^{2}+z^{2}=1\right\}$ is the unit sphere.
Let $f: \Sigma \rightarrow M$ be a wrinkled embedding. Consider a wrinkle $S \cong S^{k-1}$ of $f$ given in the local model specified by (WE3). Let $S^{\prime}:=\{z=0\} \subset S$ be the equator of $S$. By identifying the wrinkled map $f$ with its image in $M$ which we also denote by $\Sigma$, we say that $\Sigma$ has cusp singularities along $S \backslash S^{\prime}$ and unfurled swallowtail singularities along $S^{\prime}$. See Figure A.1.1.

Remark A.1.2. Although a wrinkled embedding $f: \Sigma \rightarrow M$ is in general not a smooth embedding, it follows from (WE3) that the image $f(\Sigma)$ has a well-defined $k$-dimensional tangent plane everywhere. We shall denote by $G d f: \Sigma \rightarrow \operatorname{Gr}_{k}(M)$


Figure A.1.1. Left: cusp singularity; Right: unfurled swallowtail singularity.
the corresponding "Gauss map", where $\pi: \operatorname{Gr}_{k}(M) \rightarrow M$ is the tangent $k$-plane bundle on $M$.

According to [EM09], the significance of wrinkled embeddings is that they satisfy an $h$-principle with respect to tangential rotations.

Definition A.1.3 (Tangential rotation). Given a smooth embedding $f: \Sigma \rightarrow M, a$ tangential rotation is a smooth homotopy $G_{t}: \Sigma \rightarrow \operatorname{Gr}_{k}(M), t \in[0,1]$, such that $G_{0}=G d f$ and $f=\pi \circ G_{t}$.

The following theorem was proved by Eliashberg and Mishachev in [EM09, Theorem 2.2]. Although we are only interested in codimension 1 submanifolds $\Sigma \subset M$, the theorem holds for embedded submanifolds of any codimension.
Theorem A.1.4 (Wrinkled approximation of a tangential rotation). Let $G_{t}: \Sigma \rightarrow$ $\operatorname{Gr}_{k}(M)$ be a tangential rotation of a smooth embedding $f: \Sigma \rightarrow M$. Then there exists a homotopy of wrinkled embeddings $f_{t}: \Sigma \rightarrow M$ with $f_{0}=f$ such that $G d f_{t}: \Sigma \rightarrow \operatorname{Gr}_{k}(M)$ is arbitrarily $C^{0}$-close to $G_{t}$. If the rotation $G_{t}$ is fixed on a closed set $K \subset \Sigma$, then the homotopy $f_{t}$ can also be chosen to be fixed on $K$.

Here a homotopy of wrinkled embeddings allows birth-death type singularities.
Remark A.1.5. The wrinkles that appear in Theorem A.1.4 can be made arbitrarily small.

It turns out that the unfurled swallowtail singularities in a wrinkle can be eliminated by a $C^{0}$-small operation called Whitney surgery. Whitney surgery involves first choosing an embedded $(k-1)$-disk $D$ in the wrinkled $\Sigma$ such that $\partial D=S^{\prime}$ for some wrinkle $S \subset \Sigma$ and the interior of $D$ is disjoint from the wrinkles. (The existence of such a disk $D$ is immediate.) Then one removes the unfurled swallowtail singularities along $S^{\prime}$ and adds a family of zigzags along $D$ as in Figure A.1.2. The formal treatment of Whitney surgery can be found in [EM09, §2.10].

The resulting hypersurface has only cusp singularities along spheres and we formalize it in the following definition.
Definition A.1.6 (Cuspidal embedding). A smooth map $f: \Sigma \rightarrow M$ is a spherically cuspidal embedding (or simply a cuspidal embedding) if the following hold:


Figure A.1.2. Left: before the Whitney surgery; Right: after the Whitney surgery. The vertical sides are identified in these pictures.
(CE1) $f$ is a topological embedding.
(CE2) $\operatorname{Sing}(f)$ is a finite disjoint union of smoothly embedded spheres $S_{i} \cong$ $S^{k-1}$, called cusp edges, in $\Sigma$.
(CE3) The map $f$ can be written locally on a tubular neighborhood $S_{i} \times(-\epsilon, \epsilon)$ of $S_{i}$ in $\Sigma$ as:

$$
\begin{gathered}
S^{k-1} \times(-\epsilon, \epsilon) \rightarrow S^{k-1} \times \mathbb{R}^{2} \\
(y, z) \mapsto\left(y, z^{2}, z^{3}\right)
\end{gathered}
$$

Remark A.1.7. As in the case of wrinkled embeddings, the image of a cuspidal embedding $f: \Sigma \rightarrow M$ also has well-defined tangent planes everywhere. We denote by $G d f: \Sigma \rightarrow \operatorname{Gr}_{k}(M)$ the corresponding Gauss map.

Remark A.1.8. Our cuspidal embeddings are called folded embeddings in [EM09], where the cusp edges are not necessarily diffeomorphic to the sphere. The reason we use the terminology "cuspidal embedding" is that a "folded embedding" means something else in this paper. See Definition A.3.1.

The following result follows immediately from Theorem A.1.4 and the Whitney surgery on wrinkles discussed above.

Theorem A.1.9 (Cuspidal approximation of tangential rotation). Let $G_{t}: \Sigma \rightarrow$ $\operatorname{Gr}_{k}(M)$ be a tangential rotation of a smooth embedding $f: \Sigma \rightarrow M$. Then there exists a homotopy of cuspidal embeddings $f_{t}: \Sigma \rightarrow M$ with $f_{0}=f$ such that $G d f_{t}: \Sigma \rightarrow \operatorname{Gr}_{k}(M)$ is arbitrarily $C^{0}$-close to $G_{t}$. If the rotation $G_{t}$ is fixed on a closed set $K \subset \Sigma$, then the homotopy $f_{t}$ can be chosen to be fixed on $K$.

We conclude this subsection with a smoothing operation which turns a cuspidal embedding into a smooth embedding. Suppose $f: \Sigma \rightarrow M$ is a cuspidal embedding with cusp edges $S_{i}$. Using the local model near cusps given by (CE3), the smoothing operation amounts to replacing each fiber $\left\{\left(y_{0}, z^{2}, z^{3}\right) \mid z \in(-\epsilon, \epsilon)\right\}$ at $y_{0} \in S_{i}$ by $\left\{\left(y_{0}, z^{2}, z^{\nu(z)}\right) \mid z \in(-\epsilon, \epsilon)\right\}$. Here $\nu:(-\epsilon, \epsilon) \rightarrow[1,3]$ is a smooth function which equals 1 near 0 , equals 3 near $\pm \epsilon$, is nondecreasing on $[0, \epsilon)$, and is nonincreasing on $(-\epsilon, 0]$. We denote the resulting smooth embedding by $\operatorname{Sm}(f): \Sigma \rightarrow M$ and the image by $\operatorname{Sm}(\Sigma)$.
A.2. Cuspidal embeddings of a disk. In the previous subsection, we saw that any tangential rotation of a smooth embedding can be $C^{0}$-approximated by a homotopy of wrinkled or cuspidal embeddings. However, for our purposes, we also need to change the homotopy class of the tangential distribution, and ask if it can be approximated by cuspidal embeddings. This was done in great generality by Eliashberg and Mishachev in [EM00]. In this subsection we review their work in a special case.

Let $D^{k}$ be the unit disk in $\mathbb{R}^{k}$ and let $f: D^{k} \rightarrow \mathbb{R}^{k} \times \mathbb{R}_{s}$ be a smooth embedding such that $f$ is positively transverse to $\partial_{s}$ on a neighborhood of $\partial D^{k}$. We identify $D^{k}$ with its image in $\mathbb{R}^{k+1}$ and coorient $D^{k}$ by declaring that $\partial_{s}$ is positively transverse to it near $\partial D^{k}$. Using the Euclidean metric on $\mathbb{R}^{k+1}$, let $\mathbf{n}$ be the positive unit normal vector field along $D^{k}$.

Remark A.2.1. Since $f$ is a codimension 1 embedding, specifying an oriented hyperplane distribution along $D^{k}$ is equivalent to specifying a nonvanishing vector field along $D^{k}$.

Let $C \subset \operatorname{int}\left(D^{k}\right)$ be an embedded codimension 1 submanifold which divides $D^{k}$ into two parts $D^{k} \backslash C=D_{+} \sqcup D_{-}$such that $\partial D^{k} \subset D_{+}$and the sign switches when we cross $C$. Identify a small tubular neighborhood $N(C) \subset D^{k}$ of $C$ with $C \times[-\epsilon, \epsilon]$. Choose a decomposition $C=C_{+} \cup C_{-}$and define a vector field $v$ on $N(C)$ such that $v$ points into $D_{ \pm}$along $C_{ \pm}$.

Given $\mathbf{n}$ and $C=C_{+} \cup C_{-}$as above, we define a nonvanishing vector field $\mathbf{n}(C)$ along $D^{k}$ as follows:

- $\mathbf{n}(C)=\mathbf{n}$ along $D_{+} \backslash N(C)$.
- $\mathbf{n}(C)=-\mathbf{n}$ along $D_{-} \backslash N(C)$.
- Along each fiber $\{y\} \times[-\epsilon, \epsilon] \subset C \times[-\epsilon, \epsilon]=N(C), \mathbf{n}(C)$ rotates counterclockwise from $\mathbf{n}$ to $-\mathbf{n}$ in the oriented 2-plane spanned by $(\mathbf{n}, v)$.

Roughly speaking, $C_{+}$becomes a convex suture and $C_{-}$becomes a concave suture with respect to $\mathbf{n}\left(C_{+}, C_{-}\right)$.

We state the following result [EM00, Theorem 1.7], adapted to our special case; see also [Eli72].

Theorem A.2.2. Suppose the manifolds $C_{+}$and $C_{-}$are nonempty and the vector field $\mathbf{n}(C)$ is homotopic to $\partial_{s}$ rel $\partial D^{k}$. Then there exists a cuspidal embedding $f^{\prime}: D^{k} \rightarrow \mathbb{R}^{k} \times \mathbb{R}$ that is everywhere transverse to $\partial_{s}$, such that $f^{\prime}=f$ near $\partial D^{k}$ and $\operatorname{Sm}\left(f^{\prime}\right)$ is $C^{0}$-small isotopic to $f$ rel $\partial D^{k}$.

Remark A.2.3. In fact a stronger result is given in [EM00], i.e., one can further arrange so that the cusp edges of $f^{\prime}$ coincide with $C$. This fact, however, is not needed in this paper.

Corollary A.2.4. Given any smooth embedding $f: D^{k} \rightarrow \mathbb{R}^{k} \times \mathbb{R}$ which is positively transverse to $\partial_{s}$ on a neighborhood of $\partial D^{k}$, there exists a cuspidal embedding $f^{\prime}: D^{k} \rightarrow \mathbb{R}^{k} \times \mathbb{R}$ that is everywhere transverse to $\partial_{s}$, such that $f^{\prime}=f$ near $\partial D^{k}$ and $\operatorname{Sm}\left(f^{\prime}\right)$ is $C^{0}$-small isotopic to $f$ rel $\partial D^{k}$.

Sketch of proof. Let $f: D^{k} \rightarrow \mathbb{R}^{k} \times \mathbb{R}$ be a smooth embedding which is positively transverse to $\partial_{s}$ on a neighborhood of $\partial D^{k}$.

We claim there exists $C=C_{+} \cup C_{-}$such that $\mathbf{n}(C)$ is homotopic to $\partial_{s}$ rel $\partial D^{k}$, where $C_{ \pm}$can be taken to be spherical boundaries of small neighborhoods of points in $D^{k}$. For generic $D^{k}$, there exists a finite set of points $\left\{x_{i}\right\}_{i \in I}$ in $D^{k}$ where $\mathbf{n}=-\partial_{s}$. Let $D_{i} \subset D^{k}$ be a small disk neighborhood of $x_{i}$ and let $S_{i}=\partial D_{i}$. We then take $C=\cup_{i} \partial S_{i}$ and construct $\mathbf{n}(C)$ as appropriate (depending on the local degree of the Gauss map on $D_{i}$ ) so that $\mathbf{n}(C)$ is homotopic to $\partial_{s}$. (It is not hard to add extra components of $C_{ \pm}$if necessary without affecting the homotopy type of $\mathbf{n}(C)$.

Corollary A.2.4 then follows from Theorem A.2.2.
A.3. Folding hypersurfaces. Using the techniques reviewed in §A. 1 and §A.2, we show in this subsection how to "fold" a generic hypersurface with respect to a nonvanishing vector field.

Let $\Sigma \subset M$ be a closed cooriented hypersurface and $v$ a nonvanishing vector field defined on a neighborhood of $\Sigma$. In general it is not possible to find a $C^{0}$-small isotopy $\phi_{t}: M \xrightarrow{\sim} M$ with $\phi_{0}=\mathrm{id}_{M}$ such that $\Sigma_{1}$ is everywhere transverse to $v$, where $\Sigma_{t}:=\phi_{t}(\Sigma)$. However, if we allow $\Sigma_{t}$ to have cusp singularities (here we are implicitly allowing birth-death type singularities), then there exists a cuspidal embedding $\Sigma_{1} \subset M$ which is everywhere transverse to $v$, and whose smoothing $\operatorname{Sm}\left(\Sigma_{1}\right)$ is $C^{0}$-small isotopic to $\Sigma$ - this is the content of Proposition A.3.2.

The smoothing $\operatorname{Sm}\left(\Sigma_{1}\right)$ is a $v$-folded hypersurface in the following sense:
Definition A.3.1 ( $v$-folded hypersurface). Let $\Sigma \subset M$ be a closed, cooriented hypersurface. If $v$ is a nonvanishing vector field defined on a tubular neighborhood of $\Sigma$, then $\Sigma$ is $v$-folded if there exists a codimension 1 submanifold $C(\Sigma) \subset \Sigma$ such that:
(1) $\Sigma \backslash C(\Sigma)=\Sigma_{+} \sqcup \Sigma_{-}$, where $v$ is positively (resp. negatively) transverse to $\Sigma_{+}$(resp. $\Sigma_{-}$) with respect to the coorientation of $\Sigma$ and the sign switches when we cross $C(\Sigma)$.
(2) For each connected component $C$ of $C(\Sigma)$, there exists an orientationpreserving diffeomorphism from $C \times \mathbb{R}_{x_{1}, x_{2}}^{2}$ to a tubular neighborhood $U$ of $C$ in $M$ such that $\Sigma \cap U$ is identified with $C \times\left\{x_{1}=x_{2}^{2}\right\}, C$ is identified with $C \times\{0\}$, and $\left.v\right|_{U}$ is identified with $\partial_{x_{2}}$.
The submanifold $C(\Sigma)$ is called the $v$-seam (or the seam if $v$ is understood) of $\Sigma$. Then $C(\Sigma)=C_{+}(\Sigma) \cup C_{-}(\Sigma)$, where a component $C$ of $C(\Sigma)$ belongs to $C_{+}(\Sigma)$ (resp. $C_{-}(\Sigma)$ ) if, in the local model described in (2) above, $\Sigma_{+} \cap U$ is identified with $C \times\left\{x_{1}=x_{2}^{2}, x_{2}>0\right\}$ (resp. $\Sigma_{+} \cap U$ is identified with $C \times\left\{x_{1}=x_{2}^{2}, x_{2}<0\right\}$ ).
Proposition A.3.2. Given a closed cooriented hypersurface $\Sigma \subset M$ and a nonvanishing vector field $v$ defined on a tubular neighborhood of $\Sigma$, there exists a $C^{0}$-small isotopy $\phi_{t}: M \xrightarrow{\sim} M$ with $\phi_{0}=\mathrm{id}_{M}$ such that $\Sigma_{1}=\phi_{1}(\Sigma)$ is v-folded.
Proof. Fix a Riemannian metric on $M$ such that $v$ has unit length. Let $\mathbf{n}$ be the positive unit normal vector field along $\Sigma$. We note that since $v$ and $\mathbf{n}$ may not be homotopic, some extra steps are needed to apply the wrinkling $h$-principle.

As in the proof of Corollary A.2.4, for generic $\Sigma$, there exists a finite set of points $\left\{x_{i}\right\}_{i \in I}$ in $\Sigma$ where $v=-\mathbf{n}$. Let $D_{i} \subset \Sigma$ be a small disk neighborhood of $x_{i}$ and let $S_{i}=\partial D_{i}$. Choose nested tubular neighborhoods $S_{i} \subset N_{\epsilon}\left(S_{i}\right) \subset N_{2 \epsilon}\left(S_{i}\right)$ of $S_{i}$ in $\Sigma$ such that $x_{i} \notin N_{2 \epsilon}\left(S_{i}\right)$.

It is not hard to see that there exists a homotopy $\mathbf{n}_{t}, t \in[0,1]$, of nonvanishing vector fields along $\Sigma$ with $\mathbf{n}_{0}=\mathbf{n}$ such that:
(1) $\mathbf{n}_{t}=\mathbf{n}$ on $\cup_{i \in I}\left(D_{i} \backslash N_{2 \epsilon}\left(S_{i}\right)\right)$;
(2) $\mathbf{n}_{1}=v$ on the complement of $\cup_{i \in I}\left(D_{i} \backslash N_{\epsilon}\left(S_{i}\right)\right)$.

Now we apply Theorem A.1.9 to the tangential rotation induced by $\mathbf{n}_{t}$ to obtain a $C^{0}$-approximation of $\Sigma$ by a cuspidal hypersurface $\Sigma^{\prime \prime}$, whose smoothing is $v$ folded on the complement of $\cup_{i \in I}\left(D_{i} \backslash N_{\epsilon}\left(S_{i}\right)\right)$ and such that the $v$-seam is disjoint from $S_{i}$ for all $i \in I$. To see this, first apply Theorem A.1.4 to obtain a $C^{0}{ }_{-}$ approximation $\Sigma^{\prime}$ of $\Sigma$ by a wrinkled hypersurface whose Gauss map is close to (the orthogonal complement of) $\mathbf{n}_{1}$. Since the wrinkles can be made arbitrarily small by Remark A.1.5, we can $C^{0}$-small isotop $S_{i} \subset \Sigma$ so that they avoid all the wrinkles. The Whitney surgery of the wrinkles can also be made disjoint from $S_{i}$, so we obtain a cuspidal $C^{0}$-approximation $\Sigma^{\prime \prime}$ of $\Sigma^{\prime}$ such that the cusp edges are disjoint from $S_{i}$. The smoothing of $\Sigma^{\prime \prime}$ converts the cusp edges to the $v$-seam.

Finally we apply Corollary A.2.4 to each $D_{i}$ to conclude the proof.

## Appendix B. Quantitative stabilizations of open book decompositions in general

The goal of this appendix is to explain (see Section B. 1 for the definitions of the terms involved) and prove the following generalization of Lemma 6.2.1:
Proposition B.0.1 (Quantitative stabilization). Let $\alpha$ be a contact form on $M^{2 n+1}$ which is strongly adapted to the strongly Weinstein $\operatorname{OBD}(B, \pi)$ and damped. Choose $\delta>0$ small. Suppose $\left.\alpha\right|_{B}$, the restriction to the binding $B$, is strongly adapted to the strongly Weinstein $\operatorname{OBD}\left(B_{1}, \pi_{1}\right)$, damped, and with action

$$
A\left(B_{1}, \pi_{1},\left.\alpha\right|_{B}\right)<\delta .
$$

Then there exists a strongly Weinstein OBD $\left(B^{\prime}, \pi^{\prime}\right)$ of $\left(M, \alpha^{\prime}\right)$ with a strongly adapted contact form $\alpha^{\prime} C^{1}$-close to $\alpha$ which is damped and with action

$$
A\left(B^{\prime}, \pi^{\prime}, \alpha^{\prime}\right)<\delta .
$$

While Lemma 6.2.1 suffices for the purposes of constructing a plug, we need the much more technical Proposition B.0.1 for proving Corollaries 1.3.6, 1.3.7, and 1.3.9.
B.1. Definitions. In this subsection we define the terms involved in Proposition B.0.1. In what follows let $(B, \pi)$ be an $\alpha$-compatible strongly Weinstein OBD on $M$. Its pages are denoted by $S_{\theta}=\pi^{-1}\left(e^{i \theta}\right)$.
Definition B.1.1. The action $A(B, \pi, \alpha)$ is $\sup _{\gamma \in \mathscr{R}(\alpha)} A(\gamma)$, where $\mathscr{R}(\alpha)$ is the set of Reeb chords $\gamma$ in $M \backslash S_{0}$ whose closures have endpoints on $S_{0}$ and $A(\gamma)$ is the action $\int_{\gamma} \alpha$.

Suppose $M$ admits a decomposition $M=N_{\epsilon}(B) \cup T_{\phi}$, where $N_{\epsilon}(B)$ and $T_{\phi}$ are glued along their boundary and:
(SA1) $N_{\epsilon}(B)=D^{2}(\epsilon) \times B$, where $\epsilon>0$ is small, $D^{2}(\epsilon)$ is the open disk $\{(r, \theta) \mid r<\epsilon\}$, and $(r, \theta)$ are polar coordinates, is a tubular neighborhood of the binding $\{0\} \times B$, with contact form

$$
\begin{equation*}
\left.\alpha\right|_{N_{\epsilon}(B)}=\left(1-c_{1} r^{2}\right) \lambda+c_{2} r^{2} d \theta \tag{B.1.1}
\end{equation*}
$$

where $c_{1}, c_{2}>0$ are constants.
(SA2) On $N_{\epsilon}(B), S_{\theta_{0}}$ restricts to $\left\{((r, \theta), x) \mid \theta=\theta_{0}\right\}$ for all $e^{i \theta_{0}} \in S^{1}$.
(SA3) $T_{\phi}$ is the mapping torus of $\left(S^{\star}, \phi\right)$, where $S^{\star}=S_{0}^{\star}$ is a truncated page $S_{0} \cap\left(M \backslash N_{\epsilon}(B)\right)$ and $\left.\phi\right|_{\partial S_{0}^{\star}}$ is a positive-time flow of the Reeb vector field $R_{\lambda}$.
On $N_{\epsilon}(B)$ the Reeb vector field is given by

$$
\begin{equation*}
R_{\alpha}=R_{\lambda}+\left(c_{1} / c_{2}\right) \partial_{\theta} . \tag{B.1.2}
\end{equation*}
$$

Hence all the Reeb chords $\gamma$ in $N_{\epsilon}(B) \backslash S_{0}$ have the same action $A(\gamma)=2 \pi c_{2} / c_{1}$.
Definition B.1.2. A contact form $\alpha$ adapted to the strongly Weinstein $(B, \pi)$ is strongly adapted to $(B, \pi)$ if there exists $\epsilon>0$ such that (SA1)-(SA3) hold.

Given a strongly adapted contact form which takes the form of (B.1.1) near the binding, one can easily rescale the $r$-coordinate such that $c_{2}=1$, which we assume to be the case from now on.

Definition B.1.3. A strongly adapted contact form $\alpha$ adapted to $(B, \pi)$ is damped if its total infinitesimal variation (see Definition 6.1.2) satisfies $V<11 / 10 .{ }^{5}$

Given a damped contact form $\alpha$ compatible with the $\operatorname{OBD}(B, \pi)$, one can assume that $\rho$ from Eq. (6.1.1) is almost constant.
Definition B.1.4. Given a Riemannian manifold $M$, two contact forms $\alpha_{0}, \alpha_{1}$ are $C^{1}$-close if both $\left\|\alpha_{0}-\alpha_{1}\right\|_{C^{0}}$ and $\left\|d \alpha_{0}-d \alpha_{1}\right\|_{C^{0}}$, measured with respect to the Riemannian metric, are sufficiently small.

Our notion of $C^{1}$-closeness given in Definition B.1.4 is strictly weaker than the usual notion of $C^{1}$-closeness since we take the exterior derivatives rather than all the partial derivatives.

## B.2. Perturbing to a strongly adapted contact form.

Lemma B.2.1. Let $(B, \pi)$ be a strongly Weinstein $O B D$ of $\left(M^{2 n+1}, \xi\right)$ with an adapted contact form $\alpha$. Then there exists a contact form $\alpha^{\prime}$ which is $C^{1}$-close to $\alpha$ such that $\alpha^{\prime}$ is strongly adapted to $(B, \pi)$.

Proof. By assumption $B \subset(M, \xi)$ is a contact submanifold and there exists a tubular neighborhood $N_{\epsilon}(B)=D^{2}(\epsilon) \times B$ such that $\left.\pi\right|_{N_{\epsilon}(B)}$ is the projection onto

[^4]the $D^{2}(\epsilon)$-factor. It follows that $B_{p}:=\{p\} \times B$ is contact for any $p \in D^{2}(\epsilon)$ if $\epsilon$ is sufficiently small.

The characteristic foliation $\mathscr{F}_{\theta}$ on $S_{\theta} \cap N_{\epsilon}(B)$ has a nonzero $\partial_{r}$-component for any $e^{i \theta} \in S^{1}$. Hence through any point $x \in B=B_{0}$ there exists a 2 -disk $D_{x}$ such that $D_{x} \cap S_{\theta}$ is tangent to $\mathscr{F}_{\theta}$ for all $e^{i \theta} \in S^{1}$, and the family of disks $D_{x}, x \in B$, varies smoothly in $B$. Using the disks $D_{x}$ we can reparametrize $N_{\epsilon}(B)$ such that the characteristic foliation on $S_{\theta} \cap N_{\epsilon}(B)$ is parallel to $\partial_{r}$ and $\left.\alpha\right|_{S_{\theta} \cap N_{\epsilon}(B)}$ has the form $F \lambda$, where $\lambda=\left.\alpha\right|_{B}$ and $F$ only depends on $r$. This implies that

$$
\begin{equation*}
\left.\alpha\right|_{N_{\epsilon}(B)}=F \lambda+G d \theta, \tag{B.2.1}
\end{equation*}
$$

where $F, G: N_{\epsilon}(B) \rightarrow \mathbb{R}_{\geq 0}$ satisfy (i) $\left.F\right|_{B}=1,\left.G\right|_{B}=0$, (ii) $F$ only depends on $r$, and (iii) $\lim _{r \rightarrow 0} \frac{\partial_{r} G}{r}>0$. Here (iii) follows from the contact condition $\alpha \wedge(d \alpha)^{n}>0$.

We compute

$$
\left.d \alpha\right|_{N_{\epsilon}(B)}=\partial_{r} F d r \wedge \lambda+F d \lambda+d_{B} G \wedge d \theta+\partial_{r} G d r \wedge d \theta
$$

and claim that $\left.R_{\alpha}\right|_{N_{\epsilon}(B)}$ has the form $\phi\left(\frac{\partial_{r} G}{r} R_{\lambda}-\frac{\partial_{r} F}{r} \partial_{\theta}+X\right)$, where $\phi$ is a positive function and $X \in \operatorname{Span}\left(\partial_{r}\right.$, ker $\left.\lambda\right)$. Indeed we calculate

$$
\begin{aligned}
\left.i_{R_{\lambda}} d \alpha\right|_{N_{\epsilon}(B)} & =-\partial_{r} F d r+d_{B} G\left(R_{\lambda}\right) d \theta \\
\left.i_{\partial_{\theta}} d \alpha\right|_{N_{\epsilon}(B)} & =-d_{B} G-\partial_{r} G d r \\
\left.i_{\partial_{r}} d \alpha\right|_{N_{\epsilon}(B)} & =\partial_{r} F \lambda+\partial_{r} G d \theta \\
\left.i_{Y} d \alpha\right|_{N_{\epsilon}(B)} & =F i_{Y} d \lambda+d_{B} G(Y) d \theta
\end{aligned}
$$

where $Y \in \operatorname{ker} \lambda$. The vanishing of coefficient of the $d r$ term implies that $\left.R_{\alpha}\right|_{N_{\epsilon}(B)}$ is parallel to $\frac{\partial_{r} G}{r} R_{\lambda}-\frac{\partial_{r} F}{r} \partial_{\theta}+C_{0} \partial_{r}+Y$ for some $C_{0}$ and $Y$. Since $\alpha$ is adapted by assumption, $R_{\alpha}$ must have a positive $\partial_{\theta}$-component for $r>0$. It follows that $\partial_{r} F<0$ for $r>0$.

Finally, to obtain the strongly adapted $\alpha^{\prime}$, it suffices to pick $0<\epsilon^{\prime} \ll \epsilon$ and write $\left.\alpha^{\prime}\right|_{N_{\epsilon}(B)}=F^{\prime} \lambda+G^{\prime} d \theta$ such that

- $F^{\prime}=\left(1+c_{0}\right)-c_{1} r^{2}$ and $G^{\prime}=\frac{1}{2} r^{2}$ for $r<\epsilon^{\prime}$ and some $c_{0}, c_{1}>0$;
- $F^{\prime}=F$ and $G^{\prime}=G$ for $r$ close to $\epsilon$;
- $\partial_{r} F^{\prime}<0$ and $\partial_{r} G^{\prime}>0$ for all $0<r<\epsilon$.

It is straightforward to check that $\alpha^{\prime}$ is $C^{1}$-close to $\alpha$, is contact, and is strongly adapted.

A key feature of $C^{1}$-close contact forms is the following:
Lemma B.2.2. If $\alpha, \alpha^{\prime}$ are two $C^{1}$-close contact forms on $M$, then there exists a $C^{1}$-diffeomorphism $\phi: M \xrightarrow{\sim} M$ isotopic to the identity and a function $f \in$ $C^{\infty}(M)$ which is $C^{0}$-close to the constant function 1 such that $\phi^{*}\left(\alpha^{\prime}\right)=f \alpha$.

The proof is a standard application of the proof of Gray's theorem and is omitted.
B.3. Damped OBD. In order to construct a damped contact form adapted to an OBD, we work with abstract OBDs, whose data consists of a Weinstein domain $(S, \eta)$ and an exact symplectomorphism $\phi: S \xrightarrow{\sim} S$ such that $\phi=$ id near $\partial S$ and $\phi^{*}(\eta)=\eta+d F$ for some $F \in C^{\infty}(S)$ which vanishes near $\partial S$.

We construct the compatible contact structure as follows: Let $\mathbb{R}_{t} \times S$ be the contactization of $(S, \eta)$ with contact form $\alpha=d t+\eta$. Choose a constant $C>0$ such that $F+C>0$. Let $T_{\phi, C}$ be the mapping torus

$$
T_{\phi, C}:=\{(t, x) \in \mathbb{R} \times S \mid 0 \leq t \leq F(x)+C\} /(0, \phi(x)) \sim(F(x)+C, x) .
$$

Then $\partial T_{\phi, C}=\mathbb{R} / C \mathbb{Z} \times \partial S$. We extend $\alpha$ to $N(B)=D^{2} \times B$, where $D^{2}=$ $\{(r, \theta) \mid r \leq 1\}$ (in polar coordinates), by

$$
\begin{equation*}
\left.\alpha\right|_{N(B)}=f(r) \lambda+g(r) d \theta \tag{B.3.1}
\end{equation*}
$$

where $\lambda$ is a contact form on $B$ and $f, g$ satisfy the following conditions:
(NB1) There exists $\epsilon>0$ small such that $\left.\alpha\right|_{N_{\epsilon}(B)}$ satisfies (SA1)-(SA3);
(NB2) The contact condition $\left(f^{\prime}, g^{\prime}\right) \cdot(-g, f)>0$ holds for $r>0$;
(NB3) $f(r) \lambda=\eta$ and $g(r) \equiv \frac{C}{2 \pi}$ for $r$ close to 1 . In particular $\theta=\frac{2 \pi t}{C}$ along $\{r=1\}$.
The resulting contact manifold $N(B) \cup T_{\phi, C}$ will be denoted by $M_{(S, \phi), C}$.
The mapping torus $T_{\phi, C}$ admits a foliation by

$$
S_{t}:=\operatorname{graph}\left(h_{t}\right), t \in[0, C]
$$

where $h_{t}: S \rightarrow \mathbb{R}$ varies smoothly with respect to $t$ and satisfies the following:
(1) $h_{0}=0, h_{C}=F+C$;
(2) $h_{t_{0}}(x)<h_{t_{1}}(x)$ for any $t_{0}<t_{1}$ and $x \in S$;
(3) $h_{t} \equiv t$ near $\partial S$.

Then $S_{t}, t \in[0, C]$, are the truncated pages of a compatible OBD of $M_{(S, \phi)}$; see Figure B.3.1. Each $S_{t}$ is naturally a Liouville domain with Liouville form $\eta+d h_{t}$.

We will impose a strongly Weinstein $O B D$ condition which guarantees that there exist $S_{t}, t \in[0, C]$, that are all Weinstein.


Figure B.3.1. Foliation of $T_{\phi, C}$ by the pages $S_{t}$.

Lemma B.3.1. Suppose $\alpha$ is a contact form for $(M, \xi)$ which is strongly adapted to the strongly Weinstein $O B D(B, \pi)$. Then $\alpha$ can be isotoped to the contact form $\alpha^{\prime}$ which is damped in addition.
Proof. Let $((S, \eta), \phi)$ be an abstract OBD representing $(B, \pi)$. We first discuss the mapping torus part of the open book. Starting with the fundamental domain

$$
T_{\dot{\phi}, C}^{\bullet}=\{(t, x) \in \mathbb{R} \times S \mid 0 \leq t \leq F(x)+C\} \subset \mathbb{R} \times S
$$

foliated by Weinstein domains $S_{t}=\operatorname{graph}\left(h_{t}\right), t \in[0, C]$, the trick is to thicken $T_{\phi, C}$ by inserting $\left[0, C^{\prime}\right] \times S$ with $C^{\prime} \gg 0$. More precisely, take $k \gg 0$ and consider a new angular variable $\tau \in[0, k C]$. Then the graph of

$$
H_{\tau}:=h_{\tau / k}+\frac{k-1}{k} \tau, \quad \tau \in[0, k C],
$$

defines a foliation on the mapping torus

$$
T_{\phi, k C}=\{(t, x) \in \mathbb{R} \times S \mid 0 \leq t \leq F+k C\} /(0, \phi(x)) \sim(F(x)+k C, x) .
$$

Since the Liouville form on each $S_{\tau}^{k}:=\operatorname{graph}\left(H_{\tau}\right)$ coincides with that on $S_{\tau / k}$, all the pages $S_{\tau}^{k}$ are Weinstein by assumption. Now observe that

$$
\dot{H}_{\tau}=\frac{1}{k} \dot{h}_{\tau / k}+\frac{k-1}{k} \rightarrow 1
$$

uniformly as $k \rightarrow \infty$, where the dot means $\tau$-derivative. Hence we can choose a large $k$ such that $T_{\phi, k C}$ is damped with respect to the contact form $\left.\alpha_{k}\right|_{T_{\phi, k C}}:=$ $d \tau+\eta$.

Next we extend $\alpha_{k}$ to $N(B)$. Using $\left.\alpha\right|_{N(B)}=f(r) \lambda+g(r) d \theta$ as in Eq. (B.3.1) and satisfying (NB1)-(NB3), we can write

$$
\alpha_{k}=f_{k}(r) \lambda+g_{k}(r) d \theta,
$$

near $\partial D^{2} \times B$, where $f_{k}(r)=f(r)$ and $g_{k}(r)=k g(r)=\frac{k C}{2 \pi}$ for $r$ close to 1 . Define $g_{k}:=k g$ for $r \in[0,1]$, where we are assuming without loss of generality that $g$ is constant for $r \in\left[\frac{1}{2}, 1\right]$ and is strictly increasing for $r \in\left[0, \frac{1}{2}\right]$. We extend $f_{k}$ to $r \in[0,1]$ as a strictly decreasing function in three steps as follows. Fix $\epsilon>0$ small. First extend $f_{k}$ to $r \in\left[\frac{1}{2}, 1\right]$ arbitrarily; then to $r \in\left[\frac{1}{2}-\epsilon, \frac{1}{2}\right]$ such that $f_{k}^{\prime} \leq-c g_{k}^{\prime}$, with equality near $r=\frac{1}{2}-\epsilon$, where we choose a constant $c>0$ such that $f_{k} \ll c g_{k}(1)$ for all $r \in\left[\frac{1}{2}-\epsilon, 1\right]$; and finally to $r \in\left[0, \frac{1}{2}-\epsilon\right]$ such that $f_{k}^{\prime}=-c g_{k}^{\prime}$ holds. The reader might find it helpful to note that the curve $\left\{\left(f_{k}(r), g_{k}(r)\right) \mid r \in[0,1]\right\}$ - for $k \gg 0$ and viewed from sufficiently far away - is close to the line segment connecting $\left(c g_{k}(1), 0\right)$ and $\left(0, g_{k}(1)\right)$.

We claim that the total infinitesimal variation on $N(B)$ with respect to $\left.\alpha_{k}\right|_{N(B)}$ is less than $11 / 10$ for $\epsilon$ sufficiently small. Indeed, the Reeb vector field is given by

$$
\left.R_{\alpha_{k}}\right|_{N(B)}=\frac{g_{k}^{\prime} R_{\lambda}-f_{k}^{\prime} \theta_{\theta}}{f_{k} g_{k}^{\prime}-f_{k}^{\prime} g_{k}} .
$$

The coefficient of $\partial_{\theta}$ is equal to $\frac{1}{g_{k}(1)}$ on $r \geq \frac{1}{2}$ and is equal to $\frac{c g_{k}^{\prime}}{f_{k} g_{k}^{\prime}+c g_{k}^{\prime} g_{k}}=$ $\frac{c}{f_{k}+c g_{k}} \approx \frac{1}{g_{k}(1)}$ on $r \leq \frac{1}{2}-\epsilon$. On $\frac{1}{2}-\epsilon \leq r \leq \frac{1}{2}$, we can estimate $\left|f_{k} g_{k}^{\prime} / f_{k}^{\prime}\right| \leq$ $\left|f_{k} / c\right| \ll g_{k}(1)$, which implies that the coefficient $\frac{-f_{k}^{\prime}}{f_{k} g_{k}^{\prime}-f_{k}^{\prime} g_{k}}$ of $\partial_{\theta}$ is close to $\frac{1}{g_{k}(1)}$. The claim then follows.

The contact structure $\operatorname{ker} \alpha^{\prime}:=\operatorname{ker} \alpha_{k}$ is isotopic to $\xi$ by varying the parameter $k$. Moreover $\alpha^{\prime}$ is clearly strongly adapted by construction.
B.4. Quantitative stabilization of OBD. The goal of this subsection is to prove Proposition B.0.1.

Let $(B, \pi)$ be a compatible strongly Weinstein OBD of $(M, \xi)$ and $\alpha$ be a strongly adapted contact form. Let $N_{\epsilon}(B) \cong D^{2}(\epsilon) \times B$ be the $\epsilon$-neighborhood of $B$ such that $\left.\alpha\right|_{N_{\epsilon}(B)}$ satisfies (SA1)-(SA3). Here $\epsilon>0$ is a small constant subject to conditions specified later in the proof.

We construct a map $s: M \rightarrow \mathbb{C}$ as follows: First define

$$
\left.s\right|_{N_{\epsilon}(B)}: N_{\epsilon}(B) \rightarrow D^{2}(\epsilon) \subset \mathbb{C}
$$

as the projection to the first factor $D^{2}(\epsilon) \subset \mathbb{C}$, and then uniquely extend $s$ continuously to all of $M$ by requiring it to be constant on each $S_{\theta} \backslash N_{\epsilon}(B), \epsilon e^{i \theta} \in \partial D^{2}(\epsilon)$. Hence $\pi=s /|s|$ on $M \backslash B$. Strictly speaking, $s$ is only piecewise smooth, but in what follows we will pretend that $s$ is smooth since a smoothing can easily be constructed. By definition $B=s^{-1}(0)$ and is transversely cut out.

We then consider the map $s^{k}: M \rightarrow \mathbb{C}$ for $k \in \mathbb{Z}_{>0}$. Since 0 is not a regular value of $s^{k}$ for any $k>1$, we need to add a small perturbation term coming from $B$ to $s^{k}$.

By assumption $B$ also admits a compatible strongly Weinstein $\operatorname{OBD}\left(B_{1}, \pi_{1}\right)$ such that $\lambda:=\left.\alpha\right|_{B}$ is strongly adapted. As in (SA1), let $N_{\epsilon_{1}}\left(B_{1}\right)=D_{r_{1}, \theta_{1}}^{2}\left(\epsilon_{1}\right) \times$ $B_{1}$ with $\epsilon_{1}>0$ be a neighborhood of $B_{1}$ on which

$$
\begin{equation*}
\left.\lambda\right|_{N_{\epsilon_{1}}\left(B_{1}\right)}=\left(1-c_{1,1} r_{1}^{2}\right) \lambda_{1}+r_{1}^{2} d \theta_{1} \tag{B.4.1}
\end{equation*}
$$

and let $s_{1}: B \rightarrow \mathbb{C}$ be the associated map defined in the same manner as $s$ with $N_{\epsilon}(B)$ replaced by $N_{\epsilon_{1}}\left(B_{1}\right)$. We will be using the convention that the subscript 1 (e.g., $B_{1}, s_{1}, r_{1}, c_{1,1}$ ) refers to subsets etc. of $B$ that are analogous to those of $M$ (e.g., $B, s, r, c_{1}$ ).

Next let $0<\epsilon^{\prime} \ll \epsilon$ and let $\rho:[0, \epsilon] \rightarrow \mathbb{R}_{\geq 0}$ be a $C^{\infty}$-small nonincreasing bump function such that $\rho(0)>0, \rho$ is constant on $\left[0, \epsilon^{\prime} / 2\right]$ and $\rho$ is supported on $\left[0, \epsilon^{\prime}\right]$. See Figure B.4.1 We then define the map


Figure B.4.1. The graph of $\rho$.

$$
\begin{equation*}
s_{(k)}:=s^{k}-\rho(r) s_{1}: M \rightarrow \mathbb{C} \tag{B.4.2}
\end{equation*}
$$

where $s_{1}$ is first extended to $N_{\epsilon}(B)$ by precomposing with the projection onto the second factor $B$, and then the cutoff function $\rho(r)$ guarantees that $\rho(r) s_{1}$ is globally defined on $M$.

We analyze the $\operatorname{OBD}\left(B_{(k)}, \pi_{(k)}\right)$ given by $s_{(k)}$ and the corresponding Reeb dynamics in steps. Steps 1 and 2 give topological descriptions of the binding $B_{(k)}$ and the page $S_{(k)}$ and the remaining steps describe the compatibility with a suitably $C^{1}$-small perturbed $\alpha$.

Note that besides the trivial case of $\operatorname{dim} M=1$, the case $\operatorname{dim} M=3$ is slightly different from and substantially easier than the higher-dimensional cases since $B_{1}=\varnothing$. We will point out such differences in the proof when applicable.
STEP 1. The binding $B_{(k)}=s_{(k)}^{-1}(0)$.
We can write

$$
s(x)= \begin{cases}r(x) e^{i \theta(x)}, & \text { if } x \in N_{\epsilon}(B), \\ \epsilon e^{i \theta(x)}, & \text { if } x \in M \backslash N_{\epsilon}(B),\end{cases}
$$

where $e^{i \theta(x)}=\pi(x)$ for $x \notin B$. More concisely, we write $s=r e^{i \theta}$ with the understanding that $r(x)=\epsilon$ for $x \in M \backslash N_{\epsilon}(B)$. Similarly we write $s_{1}=r_{1} e^{i \theta_{1}}$ on $B$, where $r_{1}(y)=\epsilon_{1}$ for $y \in B \backslash N_{\epsilon_{1}}\left(B_{1}\right)$. Then

$$
\begin{gather*}
s_{(k)}=r^{k} e^{i k \theta}-\rho(r) r_{1} e^{i \theta_{1}}  \tag{B.4.3}\\
B_{(k)}=\left\{s_{(k)}=r^{k} e^{i k \theta}-\rho(r) r_{1} e^{i \theta_{1}}=0\right\} .
\end{gather*}
$$

Observe that, since $r^{k} / \rho(r)$ is strictly increasing, there exists a unique $a \in$ $\left(0, \epsilon^{\prime}\right)$ such that $a^{k}=\rho(a) \epsilon_{1}$. We may assume that $\rho$ is sufficiently small such that $a \ll \epsilon^{\prime} / 2$ and hence $\rho$ is equal to the constant $a^{k} / \epsilon_{1}$ on $\left[0, \epsilon^{\prime} / 2\right]$.

We have the following description of $B_{(k)}$ :

## Claim B.4.1.

(1) $B_{(k)} \subset\{r \leq a\}$.
(2) $B_{(k)}$ is a $k$-fold branched cover of $B$ with branch locus $B_{1} \subset\{r=0\}$.
(3) $B_{(k)} \cap\{r=a\}$ is a $k$-fold cover of $B \backslash N_{\epsilon_{1}}\left(B_{1}\right)$.
(4) $B_{(k)} \cap\{r<a\}$ is graphical over $D^{2}(a):=\{(r, \theta) \mid r<a\}$ times $B_{1}$.

Proof. (1) This follows from (i) $r_{1} \leq \epsilon_{1}$ on $B$ and (ii) $B_{(k)} \cap\{\rho(r)=0\}=\varnothing$.
(2) Clearly $B_{(k)} \cap B=B_{1} \subset\{r=0\}$. Since $\rho(r) \neq 0$ on $B_{(k)}$ we can write $B_{(k)}=\left\{r_{1}=r^{k} / \rho(r), e^{i \theta_{1}}=e^{i k \theta}\right\}$. This implies that, for each point in $s_{1}^{-1}\left(r_{1}, \theta_{1}\right)$ with $r_{1}>0$, there exist $k$ distinct values of $(r, \theta)$ for which $s_{(k)}=0$.
(3) is immediate from (2).
(4) Given a point in $B_{(k)} \cap\{r<a\}$, its ( $r_{1}, \theta_{1}$ )-coordinates are determined by its $(r, \theta)$-coordinates: we have $\theta_{1}=k \theta(\bmod 2 \pi)$ and $r_{1}=r^{k} / \rho(r)$.

Note that when $\operatorname{dim} M=3$, we have $\operatorname{dim} B=1$ and $r_{1} \equiv \epsilon_{1}$. Eq. (B.4.3) implies that the new binding $B_{(k)}$ is the closure of a $k$-strand braid around $B$.
STEP 2. The page $S_{(k)}$.

It is an easy verification that the map

$$
\begin{gather*}
\pi_{(k)}=s_{(k)} /\left|s_{(k)}\right|: M \backslash B_{(k)} \rightarrow S^{1}  \tag{B.4.5}\\
\pi_{(k)}=\frac{r^{k} e^{i k \theta}-\rho(r) r_{1} e^{i \theta_{1}}}{\left|r^{k} e^{i k \theta}-\rho(r) r_{1} e^{i \theta_{1}}\right|}
\end{gather*}
$$

is a submersion, and hence induces a smooth fibration.
We analyze the page $S_{(k)}=\overline{\pi_{(k)}^{-1}(1)}$, i.e., examine the solution set to:

$$
\begin{equation*}
r^{k} e^{i k \theta}-\rho(r) r_{1} e^{i \theta_{1}} \in \mathbb{R}_{\geq 0} \tag{B.4.6}
\end{equation*}
$$

First consider $P:=S_{(k)} \cap\{r>a\}$. In this case we always have $r^{k}>\rho(r) r_{1}$ since $r_{1} \leq \epsilon_{1}$. We claim that

$$
P \cong \cup_{0 \leq j<k}\left(S_{\theta=2 j \pi / k} \cap\{r>a\}\right)
$$

where the right-hand side is the disjoint union of $k$ copies of the page $S$ and $\cong$ means the left-hand side can be viewed as a graph over the right-hand side. Indeed, if $r \geq \epsilon^{\prime}$, i.e., $\rho(r)=0$, then Eq. (B.4.6) holds precisely when $e^{i k \theta}=1$. Hence

$$
S_{(k)} \cap\left\{r \geq \epsilon^{\prime}\right\}=\cup_{0 \leq j<k}\left(S_{\theta=2 j \pi / k} \cap\left\{r \geq \epsilon^{\prime}\right\}\right) .^{6}
$$

On the other hand, if $a<r \leq \epsilon^{\prime}$, then referring to the right-hand side of Figure B.4.2, we have for any fixed $r, r_{1}, \theta_{1}$, there exists a unique $e^{i k \theta}$ such that Eq. (B.4.6) holds; the set of allowed values of $e^{i k \theta}$ is drawn in blue. Hence the components of $S_{(k)} \cap\left\{a<r \leq \epsilon^{\prime}\right\}$ can be viewed as graphs over $S_{\theta=2 j \pi / k} \cap\{a<$ $\left.r \leq \epsilon^{\prime}\right\}$, where $\theta$ is viewed as a function of $r, r_{1}, \theta_{1}$. This completes the proof of the claim.

Next consider $Q:=S_{(k)} \cap\{r<a\}$. Write $Q=Q_{1} \cup Q_{2}$, where

$$
Q_{1}:=Q \cap\left\{r^{k}<\rho(r) r_{1}\right\} \quad \text { and } \quad Q_{2}:=Q \cap\left\{r^{k} \geq \rho(r) r_{1}\right\}
$$

We first examine $Q_{1}$. Referring to the left-hand side of Figure B.4.2, any fixed $(r, \theta)$ determines a unique $\theta_{1} \in\left(\frac{\pi}{2}, \frac{3 \pi}{2}\right)$ which solves Eq. (B.4.6). Hence

$$
\begin{equation*}
Q_{1} \cong \cup_{r<a}\left(S_{1, \theta_{1}=\pi} \cap\left\{r_{1}>r^{k} / \rho(r)\right\}\right) \tag{B.4.7}
\end{equation*}
$$

where $S_{1, \theta_{1}=\pi}$ is the page of $\left(B_{1}, \pi_{1}\right)$ at $\theta_{1}=\pi$, and $\theta_{1}$ is viewed as a function of $r, r_{1}, \theta$; also the right-hand side of Eq. (B.4.7) deformation retracts to $D^{2}(a)_{(r, \theta)} \times$ $\left(S_{1, \theta_{1}=\pi} \cap\left\{r_{1}=\epsilon_{1}\right\}\right)$. Now we turn to $Q_{2}$. For each $r<a$, the set of points in $M$ such that $r^{k} \geq \rho(r) r_{1}$ is $N_{r^{k} / \rho(r)}\left(B_{1}\right)$, which is a disk times $B_{1}$ and deformation retracts to $B_{1}$. Hence

$$
Q_{2} \cong \cup_{r<a, 0 \leq j<k}\left(\left\{\frac{2 j \pi}{k}\right\} \times\{r\} \times N_{r^{k} / \rho(r)}\left(B_{1}\right)\right)
$$

where $\theta$ is viewed as a function of $r, r_{1}, \theta_{1}$.
Finally consider $R:=S_{(k)} \cap\{r=a\}$, which is the union $R_{1} \cup R_{2}$, where

$$
\begin{aligned}
R_{1} & :=\left\{r=a, r_{1}=\epsilon, e^{i \theta_{1}}=e^{i k \theta}\right\} \subset B_{(k)} \\
R_{2} & \cong \cup_{0 \leq j<k}\left(\left\{\frac{2 j \pi}{k}\right\} \times\{a\} \times N_{\epsilon_{1}}\left(B_{1}\right)\right)
\end{aligned}
$$

[^5]

Figure B.4.2. The circles have radii $r^{k}$ and $\rho(r) r_{1}$ with coordinates $e^{i k \theta}$ and $e^{i \theta_{1}}$, respectively. On the left we have $r^{k}<\rho(r) r_{1}$ and on the right $r^{k}>\rho(r) r_{1}$. The blue regions indicate the values of $e^{i k \theta}$ for which there exist $e^{i \theta_{1}}$ such that Eq. (B.4.6) holds.
where $\theta$ is a function of $r_{1}, \theta_{1}$.
Step 3. Description of modified contact form and Reeb vector field.
One technical difficulty is that the Reeb vector field given by Eq. (B.1.2) is almost never tangent to $B_{(k)}$ unless the Reeb flow of $R_{\lambda}$ on $B \backslash B_{1}$ takes pages to pages. (Note that this can easily be achieved when $\operatorname{dim} M=3$ since $\operatorname{dim} B=1$ and $B_{1}=\varnothing$ in this case.) The main task of this step therefore is to "synchronize" the Reeb flows on $M \backslash B$ and $B \backslash B_{1}$. In particular, we apply a $C^{1}$-small perturbation to $\left.\alpha\right|_{N_{\epsilon}(B)}$ given by (B.1.1) into the form given by (B.4.11).

Fix $\delta>0$ small as in the assumption of the proposition. Recall $\lambda=\left.\alpha\right|_{B}$, which is strongly adapted to $\left(B_{1}, \pi_{1}\right)$ and satisfies Eq. (B.4.1) on $N_{\epsilon_{1}}\left(B_{1}\right)$. We first choose a splitting of the exact sequence

$$
\begin{equation*}
0 \rightarrow T S_{1}^{\circ} \rightarrow T\left(B \backslash B_{1}\right) \rightarrow T_{S_{1}^{\circ}}\left(B \backslash B_{1}\right) \rightarrow 0 \tag{B.4.8}
\end{equation*}
$$

where $S_{1}^{\circ}$ denotes the interior of the page: For $r_{1} \leq \epsilon_{1} / 2$, the splitting is given by the product structure $N_{\epsilon_{1} / 2}\left(B_{1}\right)=D^{2}\left(\epsilon_{1} / 2\right) \times B_{1}$. On $B \backslash N_{\epsilon_{1}}\left(B_{1}\right)$, we choose the splitting $T\left(B \backslash N_{\epsilon_{1}}\left(B_{1}\right)\right)=\mathbb{R}\left\langle R_{\lambda}\right\rangle \oplus T S_{1}^{\star}$, where $S_{1}^{\star}=S_{1} \backslash N_{\epsilon_{1}}\left(B_{1}\right)$ is the truncated page from before. Finally we interpolate between the two splittings on $\left\{\epsilon_{1} / 2 \leq r_{1} \leq \epsilon_{1}\right\}$, e.g., one can use the linear interpolation with suitable smoothing at the endpoints.

We then define the vector field $\partial_{\theta_{1}}$ on $B \backslash B_{1}$ via the splitting (B.4.8) and recalling Eq. (B.1.2) (applied to $\lambda$ instead of $\alpha$ ) we write:

$$
\begin{equation*}
R_{\lambda}=f\left(r_{1}\right) R_{\lambda_{1}}+h \partial_{\theta_{1}}, \tag{B.4.9}
\end{equation*}
$$

where $f\left(r_{1}\right)$ is supported in $\left\{r_{1}<\epsilon_{1}\right\}, f\left(r_{1}\right) \equiv 1$ for $r_{1} \leq \epsilon_{1} / 2, h: B \rightarrow \mathbb{R}_{>0}$ is constant on $\left\{r_{1}<\epsilon_{1}\right\}$, and the functions $f$ and $h$ depend on the splitting (B.4.8). Here $h$ can be taken to be almost constant since ( $B_{1}, \pi_{1}$ ) is damped.

Since $\alpha$ is strongly adapted to ( $B, \pi$ ), there exists a small tubular neighborhood $N_{\epsilon}(B)=D^{2}(\epsilon) \times B$ of $B$ such that $\left.\alpha\right|_{N_{\epsilon}(B)}=\left(1-c_{1} r^{2}\right) \lambda+r^{2} d \theta$. There exists
$0<\epsilon^{\prime} \ll \epsilon$ and a perturbation of $\alpha$ (with the same name) such that

$$
\begin{equation*}
\left.\alpha\right|_{N_{\epsilon}(B)}=F(x, r) \lambda+r^{2} d \theta, \tag{B.4.10}
\end{equation*}
$$

where $F: N_{\epsilon}(B) \rightarrow \mathbb{R}_{>0}$ depends on $x \in B$ and $r$ and satisfies:
(i) $F(x, r)=1-c_{1} r^{2}$ near $r=\epsilon$;
(ii) $F(x, r)=1-g r^{2}$ for $r \leq \epsilon^{\prime}$, where $g: B \rightarrow \mathbb{R}_{>0}$ is $h / k$ and $k$ is a large positive integer such that $A(B, \pi, \alpha) / k<\delta ;{ }^{7}$
(iii) $\frac{\partial F}{\partial r}<0$ for any $x \in B$ and $r>0$.

Here (iii) guarantees that $\left.\alpha\right|_{N_{\epsilon}(B)}$ is contact and is adapted to $(B, \pi)$. (A direct computation shows that $R_{\alpha}$ is parallel to $R_{\lambda}-\frac{\partial_{r} F}{2 r} \partial_{\theta}+v$, where $v \in \eta=\operatorname{ker} \lambda$, and hence is positively transverse to the pages.) We note that

$$
\begin{equation*}
\left.\alpha\right|_{N_{\epsilon^{\prime}}(B)}=\left(1-g r^{2}\right) \lambda+r^{2} d \theta . \tag{B.4.11}
\end{equation*}
$$

Using Eq. (B.4.9) we compute:

$$
\begin{gathered}
\left.d \alpha\right|_{N_{\epsilon^{\prime}}(B)}=\left(1-g r^{2}\right) d \lambda-r^{2} d g \wedge \lambda-2 g r d r \wedge \lambda+2 r d r \wedge d \theta, \\
\left.R_{\alpha}\right|_{N_{\epsilon^{\prime}}(B)}=R_{\lambda}+g \partial_{\theta}+v=f\left(r_{1}\right) R_{\lambda_{1}}+h \partial_{\theta_{1}}+g \partial_{\theta}+v,
\end{gathered}
$$

where $v$ is the unique vector field tangent to $\eta=\operatorname{ker} \lambda$ solving the equation

$$
\begin{equation*}
\left(1-g r^{2}\right) i_{v} d \lambda+r^{2} d_{\eta} g=0 . \tag{B.4.12}
\end{equation*}
$$

Here $d_{\eta} g=d g-d g\left(R_{\lambda}\right) \lambda$. Note that if Eq. (B.4.12) holds, then $d g(v)=0$ and we have

$$
\left(1-g r^{2}\right) i_{v} d \lambda-r^{2}(d g(v)) \lambda+r^{2} d_{\eta} g=0 .
$$

Since both $g$ and $d_{\eta} g$ are bounded, $|v|$ is small as long as $r$ is sufficiently small. Using the splitting $T\left(B \backslash B_{1}\right)=\mathbb{R}\left\langle\partial_{\theta_{1}}\right\rangle \oplus T S_{1}$, we can write

$$
\begin{equation*}
v=\mu \partial_{\theta_{1}}+\tilde{v}, \quad \tilde{v} \in T S_{1} . \tag{B.4.13}
\end{equation*}
$$

We then obtain:

$$
\begin{equation*}
\left.R_{\alpha}\right|_{N_{\epsilon^{\prime}}(B)}=f\left(r_{1}\right) R_{\lambda_{1}}+\tilde{v}+\tilde{h} \partial_{\theta_{1}}+g \partial_{\theta}, \quad \text { where } \quad \tilde{h}=h+\mu . \tag{B.4.14}
\end{equation*}
$$

The $C^{0}$-norm $\|\tilde{h}-k g\|=\|\mu\|=O\left(r^{2}\right)$ for small $r$ by Eq. (B.4.12); we assume that $\|\tilde{h}\| \gg\|\mu\|$ since we may take $\epsilon^{\prime}>0$ to be small.

Moreover, note that $\tilde{h}=k g$ and $\tilde{v}=0$ on $\left\{r_{1}<\epsilon_{1}\right\}$, since $h$ and $g=h / k$ are constant there and $v=0$ by Eq. (B.4.12). Hence $R_{\alpha}=f\left(r_{1}\right) R_{\lambda_{1}}+h \partial_{\theta_{1}}+\frac{h}{k} \partial_{\theta}$. One easily computes that $d s_{(k)}\left(R_{\alpha}\right)=i h s_{(k)}$. Hence $R_{\alpha}$ is tangent to $B_{(k)}$ and transverse to $S_{(k)}$ on $N_{\epsilon^{\prime}}(B) \cap\left\{r_{1}<\epsilon_{1}\right\}$.

Unfortunately, since $\tilde{h} \neq k g$ in general, $\left.R_{\alpha}\right|_{N_{\epsilon^{\prime}}(B)}$ is not everywhere tangent to $B_{(k)}$. In fact, the first-order PDE $h=k g-\mu$, whose solution would have solved the problem, has no solutions in $g$ for general $h$. We will deal with this technical issue in the next few steps.

As a motivation for the above construction, we state the following, which is proved in Step 5:

[^6]Claim B.4.2. There exists a small tubular neighborhood of the stabilized binding $B_{(k)}$, away from which we have

$$
\begin{equation*}
A\left(B_{(k)}, \pi_{(k)}, \alpha\right) \leq \max \left(A\left(B_{1}, \pi_{1}, \lambda\right), A(B, \pi, \alpha) / k\right)<\delta \tag{B.4.15}
\end{equation*}
$$

STEP $4^{-}$. At this point we summarize the order in which we choose the constants. We are given $\delta>0$ and $\epsilon_{1}>0$. We choose $k>0$ such that $A(B, \pi, \alpha) / k<\delta$. We then choose $\epsilon>0$, followed by $\epsilon^{\prime}>0$. Finally we choose a small $\rho(r)$ so that $a>0$ satisfying $a^{k}=\rho(a) \epsilon_{1}$ is much smaller than $\epsilon^{\prime}$. Recall that $\rho$ is constant on $0 \leq r \leq \epsilon^{\prime} / 2$.

Step 4. The stabilized binding $B_{(k)}$ is contact.
The goal of this step is to show that $B_{(k)} \subset(M, \xi)$, as constructed in Step 1, is a contact submanifold. It follows immediately that a small tubular neighborhood of $B_{(k)}$ is foliated by contact submanifolds since the contact condition is open. We then prove Claim B.4.3, which estimates the size of such a neighborhood.

In the following we calculate modulo error terms of order $O\left(a^{2}\right)$; recall that $B_{(k)} \subset\{r \leq a\}$ by Claim B.4.1(1). For $\alpha$ satisfying Eq. (B.4.11) we obtain:

$$
\begin{gather*}
\alpha=\lambda+O\left(a^{2}\right), \quad d \alpha=d \lambda-2 g r d r \lambda+2 r d r d \theta+O\left(a^{2}\right)  \tag{B.4.16}\\
\alpha \wedge d \alpha^{n-1}=\lambda \wedge\left(d \lambda^{n-1}+(n-1) 2 r d r d \theta d \lambda^{n-2}\right)+O\left(a^{2}\right) \tag{B.4.17}
\end{gather*}
$$

First consider $B_{(k)} \cap\{r=a\}$, which is a $k$-fold cover of $B \backslash N_{\epsilon_{1}}\left(B_{1}\right)$ by Claim B.4.1(3). Since $r$ is constant, Eq. (B.4.17) becomes

$$
\begin{equation*}
\left.\alpha \wedge d \alpha^{n-1}\right|_{B_{(k)} \cap\{r=a\}}=\lambda \wedge d \lambda^{n-1}+O\left(a^{2}\right) \tag{B.4.18}
\end{equation*}
$$

which implies that $B_{(k)} \cap\{r=a\}$ is contact.
Next consider $B_{(k)} \cap\{r<a\}$, which by Claim B.4.1(4) is graphical over $D^{2}(a) \times B_{1}$. By expanding Eq. (B.4.17) using $\lambda=\left(1-c_{1,1} r_{1}^{2}\right) \lambda_{1}+r_{1}^{2} d \theta_{1}$ and using the fact that $r_{1} d r_{1} d \theta_{1}$ can be written as a nonnegative function of $r$ times $d r d \theta$ away from $r_{1}=0$,

$$
\begin{equation*}
\left.\alpha \wedge d \alpha^{n-1}\right|_{B_{(k)} \cap\{r<a\}}=\phi(r) r d r d \theta \wedge \lambda_{1} \wedge d \lambda_{1}^{n-2}+O\left(a^{2}\right) \tag{B.4.19}
\end{equation*}
$$

where $\phi(r) \geq 2(n-1)\left(1-c_{1,1} \epsilon^{2}\right)^{n-1}$. Hence $B_{(k)} \cap\{r<a\}$ is contact.
Claim B.4.3. There exists a small constant $c>0$ which depends on $\alpha$ but not on $a$ and $\rho$, for which the tubular neighborhood $N_{c a^{k}}\left(B_{(k)}\right):=s_{(k)}^{-1}\left(\left\{|w| \leq c a^{k}\right\}\right)$ of $B_{(k)}$ is foliated by contact submanifolds $B_{w}:=\left\{s_{(k)}=w\right\}$.

Proof. Let $|w| \leq c a^{k}$. We consider $B_{w} \cap\left\{r_{1}=\epsilon_{1}\right\}$. Since $\rho=a^{k} / \epsilon_{1}, w=$ $r^{k} e^{i k \theta}-a^{k} e^{i \theta_{1}}$. Then $r$, viewed as a function $B_{w} \cap\left\{r_{1}=\epsilon_{1}\right\} \rightarrow \mathbb{R}$, is close to $r=a$, with error that goes to zero as $c \rightarrow 0$, where $c$ is independent of $a$ and $\rho$. The middle term on the right-hand side of Eq. (B.4.17) then goes to zero as $c \rightarrow 0$. Hence $B_{w} \cap\left\{r_{1}=\epsilon_{1}\right\},|w| \leq c a^{k}$, is contact provided $c>0$ is small.

Next we consider $B_{w} \cap\left\{r_{1}<\epsilon_{1}\right\},|w| \leq c a^{k}$, which we write as a graph:

$$
r_{1} e^{i \theta_{1}}=(1 / \rho(r))\left(r^{k} e^{i k \theta}-w\right)
$$

By writing $r_{1} d r_{1} d \theta_{1}$ in terms of $d r d \theta$ as in Eq. (B.4.19) and observing that $\rho=$ $a^{k} / \epsilon_{1}$, we see that $B_{w} \cap\left\{r_{1}<\epsilon_{1}\right\}$ is contact.

STEP 5. Transversality away from $B_{(k)}$.
Lemma B.4.4. For the small constant $c>0$ from Claim B.4.3, $R_{\alpha}$ is transverse to $S_{(k)}^{\circ}$ on $M \backslash N_{c a^{k}}\left(B_{(k)}\right)$, provided $a>0$ is sufficiently small.

Proof. We fix $S_{(k)}$ to be the page at angle 0 . All the other pages can be treated in the same manner.

First consider the restriction of $S_{(k)}$ to $M \backslash N_{\epsilon^{\prime}}(B)$. Outside of $N_{\epsilon}(B)$ the contact form $\alpha$ is the original one and $R_{\alpha}$ is transverse to $S_{(k)}$ since $\rho=0$. On $N_{\epsilon}(B) \backslash N_{\epsilon^{\prime}}(B)$, we still have $\rho=0$ and by (iii) following Eq. (B.4.10), $\alpha$ is adapted to $(B, \pi)$ and $R_{\alpha}$ is transverse to $S_{(k)}$. Recall that $B_{(k)} \subset\{r \leq a\}$, where $a<\epsilon^{\prime}$.

Next we restrict to $N_{\epsilon^{\prime}}(B)$, on which $R_{\alpha}$ is given by Eq. (B.4.14). Observe that $R_{\alpha}$ is transverse to $S_{(k)}$ when $d s_{(k)}\left(R_{\alpha}\right)$ has positive $i \mathbb{R}$-component. When $r_{1}<\epsilon_{1}$, by the paragraph after Eq. (B.4.14) in Step 3, $R_{\alpha}$ is tangent to $B_{(k)}$ and transverse to $S_{(k)}^{\circ}$ on this region.

It remains to consider $N_{\epsilon^{\prime}}(B) \cap\left\{r_{1}=\epsilon_{1}\right\}$. Since $f\left(r_{1}\right)=0$ on $r_{1}=\epsilon_{1}$ by definition (see the line after Eq. (B.4.9)), $R_{\alpha}=\tilde{v}+\tilde{h} \partial_{\theta_{1}}+\frac{\tilde{\tilde{-}}-\mu}{k} \partial_{\theta}$ by Eq. (B.4.14); here $\tilde{v} \in T S_{1}$ by Eq. (B.4.13). We compute

$$
\begin{equation*}
d s_{(k)}\left(R_{\alpha}\right)=i\left(k g r^{k} e^{i k \theta}-\tilde{h} \rho(r) r_{1} e^{i \theta_{1}}\right)=i \tilde{h} s_{(k)}-i \mu r^{k} e^{i k \theta} \tag{B.4.20}
\end{equation*}
$$

By the paragraph after Eq. (B.4.4), $\rho=a^{k} / \epsilon_{1}$; hence $s_{(k)}=r^{k} e^{i k \theta}-a^{k} e^{i \theta_{1}}$ by Eq. (B.4.3). By the definition of $N_{c a^{k}}\left(B_{(k)}\right)$, for $x \in S_{(k)} \backslash N_{c a^{k}}\left(B_{(k)}\right)$ we have $\left|s_{(k)}(x)\right| \geq c a^{k}$. Recall that $\tilde{h}$ is bounded below by a positive constant and $\|\mu\|=$ $O\left(r^{2}\right)$. If $r(x) \leq C a$, where $C>0$ is a large constant, then $\left\|\tilde{h} s_{(k)}\right\| \geq c^{\prime} c a^{k}$ for some $c^{\prime}>0$ and $\left\|\mu r^{k} e^{i k \theta}\right\| \leq c^{\prime \prime} C^{k+2} a^{k+2}$ for some $c^{\prime \prime}>0$. If $a$ is sufficiently small, then $\left\|\tilde{h} s_{(k)}\right\|>\left\|\mu r^{k} e^{i k \theta}\right\|$, and hence $d s_{(k)}\left(R_{\alpha}\right)$ is dominated by the first term $i \tilde{h} s_{(k)}$. If $r(x)>C a$, then $s_{(k)}$ is dominated by the $r^{k} e^{i k \theta}$ term whose size is bounded below by $C^{k} a^{k}$. This in turn implies that $d s_{(k)}\left(R_{\alpha}\right)$ is dominated by the first term $i \tilde{h} s_{(k)}$ since we may take $\|\tilde{h}\| \gg\|\mu\|$. Hence $R_{\alpha}$ is transverse to $S_{(k)} \backslash N_{c a^{k}}\left(B_{(k)}\right)$ on $N_{\epsilon^{\prime}}(B) \cap\left\{r_{1}=\epsilon_{1}\right\}$.

Claim B.4.2 now readily follows from the observation that the maximal action of Reeb chords of $\left(B_{(k)}, \pi_{(k)}\right)$ in $N_{\epsilon}(B) \backslash N_{c a}\left(B_{(k)}\right)$ is approximately equal to $A\left(B_{1}, \pi_{1}, \lambda\right)$.
STEP 6. Transversality near the binding $B_{(k)}$.
In this step we modify $\alpha$ on $N_{c a^{k}}\left(B_{(k)}\right)$ so that the Reeb vector field becomes compatible with the $\operatorname{OBD}\left(B_{(k)}, \pi_{(k)}\right)$. Since $R_{\alpha}$ is already transverse to the pages and tangent to the binding for $r_{1}<\epsilon_{1}$ by the proof of Lemma B.4.4, we assume $r_{1}=\epsilon_{1}$ throughout this step.

Let $B_{(k)}^{\star}:=B_{(k)} \cap\{r=a\}$ be the $k$-fold cyclic cover of $B \backslash N_{\epsilon_{1}}\left(B_{1}\right)$ and let $S_{1}^{\star}:=S_{1} \backslash N_{\epsilon_{1}}\left(B_{1}\right)$ be a truncated page of $\left(B_{1}, \pi_{1}\right)$. (It is instructive to keep in mind the $\operatorname{dim} M=3$ case, i.e., when $S_{1}=S_{1}^{\star}$ is a point.) We identify

$$
B_{(k)}^{\star} \simeq\left\{(z, w) \in \mathbb{R} \times S_{1}^{\star} \mid 0 \leq z \leq C(w)\right\} / \sim,
$$

where $C: S_{1}^{\star} \rightarrow \mathbb{R}_{>0}$ is a smooth function such that $C(w)$ is roughly linear in $k$ for each $w \in S_{1}^{\star}$ and $(C(w), w) \sim(0, \phi(w))$ for some diffeomorphism $\phi$ of $S_{1}^{\star}$. Moreover we may assume that $\left.\alpha\right|_{B_{(k)}^{\star}}=d z+\sigma$ for a Liouville form $\sigma$ on $S_{1}^{\star}$. We then identify

$$
\begin{equation*}
N_{c a^{k}}\left(B_{(k)}^{\star}\right) \simeq D_{x, y}^{2} \times B_{(k)}^{\star} \simeq D_{x, y}^{2} \times\{0 \leq z \leq C(w)\} / \sim, \tag{B.4.21}
\end{equation*}
$$

so that:

- $R_{\alpha}$ is parallel to $\partial_{z}+W+O\left(a^{2}\right)$ on $N_{c a^{k}}\left(B_{(k)}^{\star}\right)$, where $W$ is tangent to $S_{1}^{\star}$ and also tangent to $\partial S_{1}^{\star}$, by Eq. (B.4.14), and
- $S_{(k)} \cap\left(\partial D^{2} \times B_{(k)}^{\star}\right)$ is a fiber $\Pi^{-1}(p), p \in \partial D^{2}$, of the projection $\Pi$ : $N_{c a^{k}}\left(B_{(k)}^{\star}\right) \rightarrow D^{2}$.
By Claim B.4.3, the fibers of $\Pi$ are contact submanifolds and, after a diffeomorphism of $N_{c a^{k}}\left(B_{(k)}^{\star}\right)$ which is isotopic to the identity and takes fibers to fibers, $\alpha$ (for the rest of the step we elide $\left.\right|_{N_{c a}^{k}\left(B_{(k)}^{\star}\right)}$ ) from the notation) can be written as

$$
\begin{equation*}
\alpha=F(d z+\sigma)+\frac{1}{2}(x d y-y d x), \tag{B.4.22}
\end{equation*}
$$

where $F$ is a positive function on $N_{c a^{k}}\left(B_{(k)}^{\star}\right)$ such that $\|F-1\|_{C^{1}}=O\left(a^{2}\right)$ and we are taking $D^{2}$ to be a disk of small radius depending on $a$.

Writing $R_{\alpha} \| X:=\partial_{z}+a \partial_{x}+b \partial_{y}+W$, where $W$ is tangent to $S_{1}^{\star}$, we have (B.4.23)
$a=(1+\sigma(W)) F_{y}, \quad b=-(1+\sigma(W)) F_{x}, \quad F i_{W} d \sigma=(1+\sigma(W))\left(d_{S_{1}^{\star}} F-F_{z} \sigma\right)$,
where $d_{S_{1}^{\star}}$ is the differential in the $S_{1}^{\star}$-direction. Indeed we verify that $i_{X} d \alpha=0$ :

$$
\begin{aligned}
d \alpha & =\left(F_{x} d x+F_{y} d y+F_{z} d z+d_{S_{1}^{\star}} F\right) \wedge(d z+\sigma)+F d \sigma+d x d y, \\
i_{\partial_{z}} d \alpha & =-\left(F_{x} d x+F_{y} d y+d_{S_{1}^{\star}} F\right)+F_{z} \sigma, \\
i_{a \partial_{x}+b \partial_{y}} d \alpha & =\left(a F_{x}+b F_{y}\right)(d z+\sigma)+(a d y-b d x), \\
i_{W} d \alpha & =d_{S_{1}^{\star}} F(W)(d z+\sigma)-\sigma(W)\left(F_{x} d x+F_{y} d y+F_{z} d z+d_{S_{1}^{\star}} F\right)+F i_{W} d \sigma .
\end{aligned}
$$

Setting $a=(1+\sigma(W)) F_{y}$ and $b=-(1+\sigma(W)) F_{x}$, we can cancel all the $d x, d y$ terms in $i_{X} d \alpha$ as well as $\left(a F_{x}+b F_{y}\right)(d z+\sigma)$. Setting $F i_{W} d \sigma=(1+$ $\sigma(W))\left(d_{S_{1}^{\star}} F-F_{z} \sigma\right)$, we have $d_{S_{1}^{\star}} F(W)=F_{z} \sigma(W)$, and the remaining terms of $i_{X} d \alpha$ can be canceled.

Observe that as $a \rightarrow 0,\|d F\|_{C^{0}}$ becomes small and hence $|\sigma(W)| \ll 1$. Also the component $Y:=\partial_{z}+(1+\sigma(W)) F_{y} \partial_{x}-(1+\sigma(W)) F_{x} \partial_{y}$ of the vector field $X$ is positively transverse to $\Pi^{-1}(\{\theta=$ const $\})$ near $\{\theta=$ const $\} \cap \partial D^{2}$, where we are using polar coordinates $(r, \theta)$ on $D^{2}$. More precisely, by the proof of Lemma B.4.4, $d s_{(k)}\left(R_{\alpha}\right)$ is dominated by the $i \tilde{h} s_{(k)}$ term; hence $Y$ is close to the vector field $Y^{\prime}$ on $\partial D^{2} \times B_{(k)}^{\star}$ that roughly winds $+k$ times around $\partial D^{2}$ while
going once around the $z$-direction. Finally we replace $F$ by the function $G$ such that:
(a) $F=G$ near $\partial N_{c a^{k}}\left(B_{(k)}^{\star}\right)$ and $\|F-G\|_{C^{1}}$ is small on $N_{c a^{k}}\left(B_{(k)}^{\star}\right)$;
(b) for each $z, w, G(x, y, z, w)$ has the form $G(0,0, z, w)+C_{0} r^{2}$ near $r=0$, where $C_{0}$ is a negative constant;
(c) $\frac{\partial G}{\partial r}<0$ for $r>0$, which guarantees transversality to the pages $S_{(k)}$ on $N_{c a^{k}}\left(B_{(k)}^{\star}\right) ;$
(d) $\alpha$ is still contact with $F$ replaced by $G$.

In view of the above description of the pages $S_{(k)} \cap N_{c a^{k}}\left(B_{(k)}^{\star}\right)$, the new contact form $C^{1}$-approximates $\left.\alpha\right|_{N_{c a^{k}}\left(B_{(k)}^{\star}\right)}$ and is strongly adapted.

STEP 7. Weinstein structure on the page $S_{(k)}$.
In this step we describe the Weinstein structure on $S_{(k)}$, which we fix to be the page at angle 0 . All the other pages can be treated in a similar manner. We will decompose $S_{(k)}=T_{1} \cup T_{2} \cup T_{3}$ into three pieces and study the characteristic foliation on each piece separately. Note that the decomposition of $S_{(k)}$ in this step will be different from, but based on, the one from Step 2.

First let $T_{1}:=S_{(k)} \cap\left\{r \geq \epsilon^{\prime}\right\}=\cup_{0 \leq j<k}\left(S_{\theta=2 j \pi / k} \cap\left\{r \geq \epsilon^{\prime}\right\}\right)$. The characteristic foliation on $T_{1}$ is Morse since $T_{1}$ agrees with restrictions of the Weinstein pages on the region $r \geq \epsilon$ and the modification given by Eq. (B.4.10) does not change the dynamical properties of the characteristic foliation on the region $\epsilon^{\prime} \leq r \leq \epsilon$.

Next let $T_{3}:=S_{(k)} \cap\left(D_{r, \theta}^{2}\left(\epsilon^{\prime}\right) \times N_{\epsilon_{1}}\left(B_{1}\right)\right)$ and use the contact form $\alpha$ before the modification done in Step 6, since the topological conjugacy type of the characteristic foliation on $T_{3}$ is the same for both contact forms. On $N_{\epsilon^{\prime}}(B)=D_{r, \theta}^{2}\left(\epsilon^{\prime}\right) \times B$ we use the contact form $\alpha=\lambda+\frac{r^{2}}{1-g r^{2}} d \theta$ (see Eq. (B.4.11)) which is $C^{1}$-close to $\lambda+r^{2} d \theta$ when $\epsilon^{\prime}$ is small. Restricted to $\left(D_{r, \theta}^{2}\left(\epsilon^{\prime}\right) \times N_{\epsilon_{1}}\left(B_{1}\right)\right.$ we have

$$
\alpha=\left(1-c_{1,1} r_{1}^{2}\right) \lambda_{1}+r_{1}^{2} d \theta_{1}+\frac{r^{2}}{1-g r^{2}} d \theta
$$

Let us write $\tilde{r}=r^{k} / \rho(r)$ and $\tilde{\theta}=k \theta$. Using Cartesian coordinates $(\tilde{x}, \tilde{y}),\left(x_{1}, y_{1}\right)$ corresponding to the polar coordinates $(\tilde{r}, \tilde{\theta}),\left(r_{1}, \theta_{1}\right)$, Eq. (B.4.6) becomes $\tilde{x}-$ $x_{1} \in \mathbb{R}_{\geq 0}$ and $\tilde{y}=y_{1}$, and we write

$$
\alpha=\left(1-c_{1,1} r_{1}^{2}\right) \lambda_{1}+\left(x_{1} d \tilde{y}-\tilde{y} d x_{1}\right)+\psi(\tilde{r})(\tilde{x} d \tilde{y}-\tilde{y} d \tilde{x})
$$

where $\psi>0$. We now calculate the characteristic foliation $Y$.

$$
\begin{aligned}
d \alpha= & \left(1-c_{1,1} r_{1}^{2}\right) d \lambda_{1}-2 c_{1,1}\left(x_{1} d x_{1}+\tilde{y} d \tilde{y}\right) \lambda_{1}+2 d x_{1} d \tilde{y}+2 \phi(\tilde{r}) d \tilde{x} d \tilde{y} \\
= & \left(1-c_{1,1} r_{1}^{2}\right) d \lambda_{1}+\frac{2 c_{1,1}}{1-c_{1,1} r_{1}^{2}}\left(x_{1} d x_{1}+\tilde{y} d \tilde{y}\right)\left(\left(x_{1} d \tilde{y}-\tilde{y} d x_{1}\right)+\psi(\tilde{r})(\tilde{x} d \tilde{y}-\tilde{y} d \tilde{x})\right) \\
& \quad+2 d x_{1} d \tilde{y}+2 \phi(\tilde{r}) d \tilde{x} d \tilde{y} \\
= & \left(1-c_{1,1} r_{1}^{2}\right) d \lambda_{1}+\frac{2 c_{1,1}}{1-c_{1,1} r_{1}^{2}}\left(r_{1}^{2} d x_{1} d \tilde{y}+\psi(\tilde{r})\left(x_{1} \tilde{x} d x_{1} d \tilde{y}-x_{1} \tilde{y} d x_{1} d \tilde{x}+\tilde{y}^{2} d \tilde{x} d \tilde{y}\right)\right) \\
\quad & \quad 2 d x_{1} d \tilde{y}+2 \phi(\tilde{r}) d \tilde{x} d \tilde{y}
\end{aligned}
$$

where $\phi>0$ and $\psi(\tilde{r})$ and $\phi(\tilde{r})$ are dominated by a term of the form $C \tilde{r}^{-2+2 / k}$ for $\tilde{r} \neq 0$ small. Let us write $Y=Z+\eta R_{\lambda_{1}}$, where

$$
Z=\frac{2 c_{1,1}}{1-c_{1,1} r_{1}^{2}}\left(-r_{1}^{2} \partial_{\tilde{x}}+\psi(\tilde{r})\left(-x_{1} \tilde{x} \partial_{\tilde{x}}-x_{1} \tilde{y} \partial_{\tilde{y}}+\tilde{y}^{2} \partial_{x_{1}}\right)\right)-2 \partial_{\tilde{x}}+2 \phi(\tilde{r}) \partial_{x_{1}}
$$

Then $i_{Y} d \alpha=0$ and $\alpha(Y)=0$ for a suitable choice of function $\eta$.
We claim the forward flow of $Y$ limits to $T_{3} \cap B_{(k)}=\left\{\tilde{x}=x_{1}, \tilde{y}=y_{1}\right\} \times B_{1}$. Assume $\tilde{r} \neq 0$. The case $\tilde{r}=0$ can be treated separately, and is left to the reader. When $x_{1}<0$, then $Y$ flows to the region $x_{1}>0$ since $r_{1}<\epsilon_{1}$ and the coefficient of $\partial_{x_{1}}$ is positive and bounded below. When $x_{1}>0$, then the coefficient of $\partial_{x_{1}}$ is positive and the coefficient of $\partial_{\tilde{x}}$ is less than -2 ; since $\tilde{x}-x_{1} \geq 0, Y$ eventually flows to $\tilde{x}=x_{1}$. The claim then follows.

Finally we consider $T_{2}:=S_{(k)} \cap\left(D_{r, \theta}^{2}\left(\epsilon^{\prime}\right) \times\left(B \backslash N_{\epsilon_{1}}\left(B_{1}\right)\right)\right)$. Recall the fibration $\pi_{1}: B \backslash N_{\epsilon_{1}}\left(B_{1}\right) \rightarrow S_{\theta_{1}}^{1}$ with fiber $S_{1}^{\star}$ such that $\left.\lambda\right|_{B \backslash N_{\epsilon_{1}}\left(B_{1}\right)}=f_{1} d \theta_{1}+\beta_{1}$, where $f_{1}$ is a positive function on $B \backslash N_{\epsilon_{1}}\left(B_{1}\right)$ and $\beta_{1}$ is a $\theta_{1}$-dependent Liouville form on $S_{1}^{\star}$. Let $Z_{k}$ be the surface in the 3 -manifold $D_{r, \theta}^{2}\left(\epsilon^{\prime}\right) \times S_{\theta_{1}}^{1}$ satisfying $r^{k} e^{i k \theta}-\rho(r) \epsilon_{1} e^{i \theta_{1}} \in \mathbb{R}_{\geq 0}$. Then $T_{2}=\left(\operatorname{id}_{D_{r, \theta}\left(\epsilon^{\prime}\right)}, \pi_{1}\right)^{-1}\left(Z_{k}\right)$. Restrict attention to where $\rho(r)$ is constant and write $\tilde{r}=\frac{r^{k}}{\epsilon_{1} \rho(r)}=\left(\frac{r}{a}\right)^{k}$ and $\tilde{\theta}=k \theta$. Then Eq. (B.4.6) gives $\tilde{r} \sin \tilde{\theta}=\sin \theta_{1}$. We view $\tilde{\theta}$ as a function of $\tilde{r}, \theta_{1}$. Without loss of generality we may replace $\frac{r^{2}}{1-g r^{2}}$ by $r^{2}$ in $\alpha$, since $\epsilon^{\prime}$ can be taken to be small. Since $\frac{\partial \tilde{\theta}}{\partial \theta_{1}}=$ $\frac{1}{\tilde{r}} \frac{\cos \theta_{1}}{\cos \tilde{\theta}}$ and $\frac{\partial \tilde{\theta}}{\partial \tilde{r}}=-\frac{1}{\tilde{r}} \frac{\sin \tilde{\theta}}{\cos \tilde{\theta}}$, we have

$$
\begin{aligned}
\alpha & =f_{1} d \theta_{1}+\beta_{1}+r^{2} d \theta=\left(f_{1}+\frac{r^{2}}{k \tilde{r}} \frac{\cos \theta_{1}}{\cos \tilde{\theta}}\right) d \theta_{1}-\frac{r^{2}}{k \tilde{r}} \frac{\sin \tilde{\theta}}{\cos \tilde{\theta}} d \tilde{r}+\beta_{1}, \\
d \alpha & =d_{S_{1}} f_{1} \wedge d \theta_{1}+d_{S_{1}} \beta_{1}+d \theta_{1} \wedge \dot{\beta}_{1}+\frac{2 r^{2}}{k^{2} \tilde{r}} d \tilde{r} \wedge d \tilde{\theta} \\
& =\left(d_{S_{1}} f_{1}-\dot{\beta}_{1}\right) \wedge d \theta_{1}+d_{S_{1}} \beta_{1}+\frac{2 r^{2}}{k^{2} \tilde{r}^{2}} \frac{\cos \theta_{1}}{\cos \tilde{\theta}} d \tilde{r} \wedge d \theta_{1}
\end{aligned}
$$

where $d_{S_{1}}$ is the exterior derivative in the $S_{1}$-direction and $\dot{\beta}_{1}=\frac{\partial \beta_{1}}{\partial \theta_{1}}$. Writing $Y=X+a \partial_{\tilde{r}}+b \partial_{\theta_{1}}$, where $X \in T S_{1}$, we solve for $Y$ in $i_{Y} d \alpha=-\left(\cos \theta_{1}\right) \alpha$. We compute:

$$
\begin{aligned}
X\lrcorner d \alpha & =i_{X} d_{S_{1}} \beta_{1}+\left(d_{S_{1}} f_{1}(X)-\dot{\beta}_{1}(X)\right) d \theta_{1}, \\
\left.a \partial_{\tilde{r}}\right\lrcorner d \alpha & =a \frac{2 r^{2}}{k^{2} \tilde{r}^{2}} \frac{\cos \theta_{1}}{\cos \tilde{\theta}} d \theta_{1}, \\
\left.b \partial_{\theta_{1}}\right\lrcorner d \alpha & =b\left(\dot{\beta}_{1}-d_{S_{1}} f_{1}\right)-b \frac{2 r^{2}}{k^{2} \tilde{r}^{2}} \frac{\cos \theta_{1}}{\cos \tilde{\theta}} d \tilde{r} .
\end{aligned}
$$

Comparing the coefficients of $d \tilde{r}$ and $d \theta_{1}$ and the $S_{1}$-component, we obtain

$$
\begin{aligned}
b & =-\frac{k}{2} \tilde{r} \sin \tilde{\theta} \\
i_{X} d_{S_{1}} \beta_{1} & =-\left(\cos \theta_{1}\right) \beta_{1}+\frac{k}{2} \tilde{r} \sin \tilde{\theta}\left(\dot{\beta}_{1}-d_{S_{1}} f_{1}\right) \\
a & =\frac{k^{2} \tilde{r}^{2}}{2 r^{2}} \frac{\cos \tilde{\theta}}{\cos \theta_{1}}\left(-\left(\cos \theta_{1}\right) f_{1}+\dot{\beta}_{1}(X)-d_{S_{1}} f_{1}(X)\right)-\frac{k \tilde{r}}{2} \cos \theta_{1}
\end{aligned}
$$

By the $b \partial_{\theta_{1}}$ term in $Y$, unless $\tilde{r}=0$, one either flows to the binding or to $\theta_{1}=0$ or $\pi$. If $\theta_{1}=\pi$, then $\tilde{\theta}=0$ or $\pi$ and $X$ is the Liouville vector field of $\beta_{1}$. One then flows to $\partial S_{1}$ and hence into $T_{3}$. If $\theta_{1}=0$, then $\tilde{\theta}=0$ and $X$ is minus the Liouville
vector field of $\beta_{1}$. One then flows to a zero of $X$ on $S_{1}$; then $a=-\frac{k^{2} \tilde{r}^{2}}{2 r^{2}} f_{1}-\frac{k \tilde{r}}{2}<0$ and $\tilde{r} \rightarrow 1$ along the flow, which means one flows to the binding.

Remark B.4.5. As an aid to understanding $T_{2}$, consider the situation when $\operatorname{dim} M=$ 3. Then $Z_{k} \subset S_{(k)}$ and is Morse. The topological determination of $Z_{k} \subset D^{2} \times S^{1}$ is straightforward and we see that the characteristic foliation on $Z_{k}$ has one index 0 critical point and $k$ index 1 critical points; see Figure B.4.3 for an illustration in the case $k=6$.


Figure B.4.3. The characteristic foliation on $Z_{k}$.

STEP 8. Damping property.
We finally prove the damping property for $\left(B_{(k)}, \pi_{(k)}\right)$.
On the region $M \backslash N_{\epsilon^{\prime}}(B)$, the page $S_{(k)}$ restricts to $\cup_{0 \leq j<k} S_{\theta=2 j \pi / k} \cap\left\{r \geq \epsilon^{\prime}\right\}$. Since $(B, \pi)$ is damped, the actions of the Reeb chords are close to $\frac{1}{k} A(B, \pi, \alpha) \approx$ $\frac{2 \pi}{k c_{1}}$ in view of Eq. (B.1.2) with the normalization $c_{2}=1$.

On the region $N_{\epsilon^{\prime}}(B)-N_{c a}\left(B_{(k)}\right)$, by Eq. (B.4.14), the relevant component of $R_{\alpha}$ is close to $h \partial_{\theta_{1}}+\frac{h}{k} \partial_{\theta}=h\left(\partial_{\theta_{1}}+\frac{1}{k} \partial_{\theta}\right)$ and the actions of the Reeb chords are close to $2 \pi / h$. (We can see this for example from $d s_{(k)}\left(\partial_{\theta_{1}}+\frac{1}{k} \partial_{\theta}\right)=i s_{(k)}$.)

On the region $N_{c a}\left(B_{(k)}\right)$, in view of the modifications (a)-(c) from Step 6, the actions of the Reeb chords are close to $A\left(B_{1}, \pi_{1}, \lambda\right) \approx \frac{2 \pi}{h}$. Taken together, $\left(B^{\prime}, \pi^{\prime}\right)$ is damped with respect to a $C^{1}$-small perturbation of $\alpha^{\prime}$.

This completes the proof of Proposition B.0.1.

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University of California, Los Angeles, Los Angeles, CA 90095<br>Email address: honda@math.ucla.edu<br>URL: http://www.math.ucla.edu/~honda

Somewhere On Earth
Email address: hymath@gmail.com
URL: https://sites.google.com/site/yhuangmath


[^0]:    ${ }^{1}$ Note the latter involves Liouville domains, which is a much harder problem.

[^1]:    ${ }^{2}$ Since Morse-Smale vector fields are considered in [Gir00], there exists a different kind of bifurcation where a pair of periodic orbits appear or disappear. This phenomenon does not occur here since we are dealing with Morse gradient vector fields.

[^2]:    ${ }^{3}$ For the moment we only know that $S$ is a homotopy sphere. Presumably $S$ is a standard sphere, although we will not need this.

[^3]:    ${ }^{4}$ We can view this as the contact analog of symplectic reduction.

[^4]:    ${ }^{5}$ Here the precise constant slightly larger than 1 is not important, but we choose one for definiteness.

[^5]:    ${ }^{6}$ We will be writing $S_{\theta=*}$ to avoid confusion with a page $S_{1}$ of $\left(B_{1}, \pi_{1}\right)$.

[^6]:    ${ }^{7}$ The choice of $k$ depends only on $A(B, \pi, \alpha)$ and, in particular, not on $\epsilon$.

