

SUMMER SCHOOL  
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**Compatible brackets in Hamiltonian  
mechanics**

After H. P. McKean

Presented by Irina Nenciu - Caltech

## Introduction and general results

Let  $d$  be fixed and let  $C^\infty(\mathbb{R}^{2d})$  be a Lie algebra with the bracket

$$[H_1, H_2] = (\nabla H_1)^t J \nabla H_2 \quad (1)$$

where  $\nabla$  is the gradient and

$$J : \mathbb{R}^{2d} \rightarrow Gl(2d, \mathbb{R})$$

is smooth and skew. In particular, Jacobi's identity is satisfied:

$$[[H_1, H_2], H_3] + [[H_2, H_3], H_1] + [[H_3, H_1], H_2] = 0 \quad (2)$$

for any  $H_1, H_2, H_3 \in C^\infty(\mathbb{R}^{2d})$ .

Note that Jacobi's identity (2) is equivalent to any of the following conditions on  $J$ :

- For all triples  $(i, j, k)$

$$(J\nabla)_i J_{jk} + (J\nabla)_j J_{ki} + (J\nabla)_k J_{ij} = 0; \quad (3)$$

- For all triples  $(i, j, k)$

$$\frac{\partial J_{ij}^{-1}}{\partial x_k} + \frac{\partial J_{jk}^{-1}}{\partial x_i} + \frac{\partial J_{ki}^{-1}}{\partial x_j} = 0; \quad (4)$$

- $\sum_{i < j} J_{ij}^{-1} dx_i \wedge dx_j$  is closed.

Assume that there exist two such brackets

$$\begin{aligned}[H_1, H_2]_J &= (\nabla H_1)^t J \nabla H_2 \\ [H_1, H_2]_K &= (\nabla H_1)^t K \nabla H_2\end{aligned}$$

with  $J$  and  $K$  satisfying the conditions on the previous slides.

**Definition 1** *We say that a Hamiltonian  $H_0$  can be raised if there exists another Hamiltonian  $H_1$  so that*

$$[f, H_0]_K = [f, H_1]_J \quad (5)$$

*for all  $f \in C^\infty(\mathbb{R}^{2d})$ . In this case we say that  $H_0$  raises to  $H_1$  or that  $H_1$  lowers to  $H_0$ , and denote these relations by  $H_0 \uparrow H_1$  and  $H_1 \downarrow H_0$ , respectively.*

We immediately obtain the following

**Lemma 2** *Consider  $J$  and  $K$  as above and assume that there exists a sequence of Hamiltonians so that  $H_0 \uparrow H_1 \uparrow H_2 \uparrow \dots$ . Then*

$$[H_i, H_j]_J = [H_i, H_j]_K = 0 \quad (6)$$

*for all  $i, j \geq 0$ .*

Notice that, for  $i < j$ ,

$$\begin{aligned} [H_i, H_j]_J &= [H_i, H_{j-1}]_K = -[H_{j-1}, H_i]_K \\ &= -[H_{j-1}, H_{i+1}]_J = [H_{i+1}, H_{j-1}]_J. \end{aligned}$$

We have

$$[H_i, H_j]_J = [H_{i+1}, H_{j-1}]_J.$$

Then

- If  $i$  and  $j$  have the same parity, then after a finite number of steps of the above computation we get

$$[H_i, H_j]_J = [H_k, H_k]_J = 0$$

where  $k = \frac{i+j}{2}$ .

- If  $i$  and  $j$  have different parities, then

$$[H_i, H_j]_J = [H_j, H_i]_J = 0.$$

So  $[H_i, H_j]_J = 0$  for all  $i, j$ . This immediately implies that  $[H_i, H_j]_K = 0$  for all  $i, j$ .

**Question:** When is it possible to have such infinite sequences of raising Hamiltonians  $H_0 \uparrow H_1 \uparrow H_2 \uparrow \dots$ ?

**Want** a condition on  $J$  and  $K$  so that

$$H_0 \uparrow H_1 \quad \Rightarrow \quad H_1 \uparrow ?$$

for all  $H_0$ .

This in turn is implied by the condition that  $\uparrow = \downarrow$ . So the next Lemma gives the natural condition to answer our question.

**Lemma 3** *If  $J$ ,  $K$  and  $J + K$  all produce honest brackets (i.e. the brackets they define obey the Jacobi identity), then the class of Hamiltonians that can be raised coincides with the class of those that can be lowered.*

In this case we say that  $J$  and  $K$  are compatible.

**Idea of the proof.** The condition  $\uparrow \Leftrightarrow \downarrow$  is equivalent to

$J^{-1}K\nabla H$  is a gradient

$\Leftrightarrow$

$K^{-1}J\nabla H$  is a gradient

which is further equivalent to the fact that both or neither of the matrices

$$A_+(H) = \left[ \frac{\partial}{\partial x_i} (J^{-1}K\nabla H)_j \right]_{1 \leq i, j \leq 2d}$$

and

$$A_-(H) = \left[ \frac{\partial}{\partial x_i} (K^{-1}J\nabla H)_j \right]_{1 \leq i, j \leq 2d}$$

are symmetric.

This last condition is implied by the relation

$$K(A_-(H) - A_-(H)^t)K + J(A_+(H) - A_+(H)^t)J = 0. \quad (7)$$

Using the fact that  $J$  and  $K$  satisfy the Jacobi identity, one gets that equation (7) is equivalent to the fact that  $J + K$  also satisfies the Jacobi identity.

**Remark.** Let  $L(c) = J - cK$ . Then notice that

$$\begin{aligned} (2)_{L(c)} &= (2)_J + c^2 \cdot (2)_K \\ &\quad - c[(2)_{J+K} - (2)_J - (2)_K] \end{aligned}$$

is a quadratic polynomial in  $c$ .

So  $J$  and  $K$  are compatible iff  $L(c)$  satisfies the Jacobi identity for any  $c$  iff  $L(c)$  satisfies the Jacobi identity for 3 values of  $c$ .

Recall that, if  $L(c)$  is invertible, then it obeys the Jacobi identity iff  $L(c)^{-1}$  defines a closed 2-form. For  $c$  small enough

$$L(c)^{-1} = \sum_{n \geq 0} c^n (J^{-1}K)^n J^{-1}.$$

Hence if  $J$  and  $K$  are compatible, then  $(J^{-1}K)^n J^{-1}$  define closed 2-forms for all  $n \geq 0$ . Conversely, it turns out that if  $(J^{-1}K)^n J^{-1}$  define closed 2-forms for  $n = 0, 1, 2$ , then  $J$  and  $K$  are compatible.

## Magri's result

**Want** to find a second, compatible, bracket for a given completely integrable Hamiltonian system.

**Assume** first that such a bracket exists and find a necessary form for it.

**Hypothesis 1.**  $J$  and  $K$  give two compatible brackets.

Let  $\pi_k = \text{Tr}(J^{-1}K)^k$  for all  $k \geq 1$ . Then one can prove that

$$\pi_1 \uparrow \frac{\pi_2}{2} \uparrow \frac{\pi_3}{3} \uparrow \dots \quad (8)$$

and, consequently, the  $\pi_k$ 's commute.

To prove this, take  $\lambda$  large enough so that  $L = \lambda J - K$  is invertible. Then

$$\begin{aligned}
(L \nabla \log \det L)_i &= \sum_j L_{ij} \operatorname{Tr} \left( L^{-1} \frac{\partial L}{\partial x_j} \right) \\
&= \sum_{j,a,b} L_{ij} L_{ab}^{-1} \frac{\partial L_{ba}}{\partial x_j} \\
&= - \sum_{a,b,j} L_{ab}^{-1} \left( L_{bj} \frac{\partial L_{ai}}{\partial x_j} + L_{aj} \frac{\partial L_{ib}}{\partial x_j} \right) \\
&= - \sum_{a,j} \delta_{aj} \frac{\partial L_{ai}}{\partial x_j} + \sum_{b,j} \delta_{jb} \frac{\partial L_{ib}}{\partial x_j} \\
&= 2 \sum_j \frac{\partial L_{ij}}{\partial x_j}
\end{aligned}$$

for all  $i$ .

**Remark.** The third identity uses the fact that  $L$  satisfies the Jacobi identity for any value of  $\lambda$ .

Now sum over  $i$  and expand in  $\lambda$ . The right-hand side is linear in  $\lambda$ . The left hand side can be expanded for large enough  $\lambda$  as follows:

$$L\nabla\left(2d\log\lambda + \log\det J + \log\det(I - \lambda^{-1}J^{-1}K)\right)$$

$$= L\nabla\log\det J - L\nabla\text{Tr}\log(I - \lambda^{-1}J^{-1}K)$$

$$= L\nabla\log\det J - L\nabla\text{Tr}\left(\sum_{n\geq 1}\lambda^{-n}\frac{(J^{-1}K)^n}{n}\right)$$

$$= L\nabla\log\det J - \sum_{n\geq 0}\lambda^{-n}J\nabla\frac{\text{Tr}(J^{-1}K)^{n+1}}{n+1}$$

$$+ \sum_{n\geq 1}\lambda^{-n}K\nabla\frac{\text{Tr}(J^{-1}K)^n}{n}$$

So all the coefficients of  $\lambda^{-n}$  for  $n \geq 1$  must be zero:

$$J\nabla\left(\frac{\pi_{n+1}}{n+1}\right) = K\nabla\left(\frac{\pi_n}{n}\right)$$

for all  $n \geq 1$ . This is exactly the definition of

$$\frac{\pi_n}{n} \uparrow \frac{\pi_{n+1}}{n+1}.$$

In particular we can conclude from Lemma 2 that

$$[\pi_n, \pi_m]_J = [\pi_n, \pi_m]_K = 0$$

for all  $n, m \geq 1$ .

**Hypothesis 2.** The spectrum of  $J^{-1}K$  is *ample*; that is the sense that the gradients  $\nabla\pi_1, \dots, \nabla\pi_d$  are independent.

In this case, fixing the values of the  $\pi_k$ 's defines a smooth  $d$ -dimensional manifold  $M \subset \mathbb{R}^{2d}$  on which the  $\nabla\pi_k$ 's are normal and the  $J\nabla\pi_k$ 's are tangent for all  $1 \leq k \leq d$ :

$$\langle \nabla\pi_k, J\nabla\pi_l \rangle = (\nabla\pi_k)^t J\nabla\pi_l = [\pi_k, \pi_l]_J = 0.$$

- If  $H \in \uparrow = \downarrow$ , then  $[\pi_k, H]_J = [\pi_k, H]_K = 0$ .
- If the spectrum of  $J^{-1}K$  is ample, then it is double. Let  $\lambda_1, \dots, \lambda_d$  denote the eigenvalues of  $J^{-1}K$ .
- $H \in \uparrow = \downarrow$  if and only if  $H = \sum_j h_j(\lambda_j)$ .

Under

**Hypothesis 1.**  $J$  and  $K$  give two compatible brackets

and

**Hypothesis 2.** The gradients  $\nabla\pi_1, \dots, \nabla\pi_d$  are independent

we have that the functions  $H_k = \frac{\pi_k}{k}$  ( $1 \leq k \leq d$ ) commute and can be completed by functions  $G_k$  ( $1 \leq k \leq d$ ), canonically conjugate relative to  $J$ .

Now we can compute the  $K$ -bracket of the  $H$ 's and  $G$ 's.

$$[H_i, H_j]_K = [H_i, H_{j+1}]_J = 0.$$

$$\begin{aligned} [H_i, G_j]_K &= [H_{i+1}, G_j]_J \\ &= \begin{cases} 1, & \text{if } j = i + 1 \leq d; \\ 0, & \text{if } j \neq i + 1 \leq d; \\ \frac{\partial H_{d+1}}{\partial H_j}, & \text{if } i = d. \end{cases} \end{aligned}$$

Finally, using the Jacobi identity for  $[\cdot, \cdot]_K$  and the fact that  $[H_i, G_j]_J = \delta_{ij}$  is a constant, we get

$$\begin{aligned}
[[G_i, G_j]_K, H_k]_J &= [[G_i, G_j]_K, H_k^-]_K \\
&= -[[G_j, H_k^-]_K, G_i]_K \\
&\quad - [[H_k^-, G_i]_K, G_j]_K \\
&= -[[G_j, H_k]_J, G_i]_K \\
&\quad - [[H_k, G_i]_J, G_j]_K \\
&= 0
\end{aligned}$$

and hence

$$[G_i, G_j]_K = c_{ij}(H_1, \dots, H_d).$$

Since, by Hypothesis 2,  $\nabla H_k$  and  $\nabla G_k$  ( $1 \leq k \leq d$ ) span  $\mathbb{R}^{2d}$ , we find that  $K$  is uniquely determined by Magri's rule and it equals

$$\begin{aligned} & \sum_{i < j} c_{i,j} (J\nabla H_i \otimes J\nabla H_j - J\nabla H_j \otimes J\nabla H_i) \\ & + \sum_{j=1}^d (J\nabla H_j^\dagger \otimes J\nabla G_j - J\nabla G_j \otimes J\nabla H_j^\dagger) \end{aligned}$$

Let us remember that, if  $v, w \in \mathbb{R}^{2d}$ , then

$$v \otimes w : \mathbb{R}^{2d} \rightarrow \mathbb{R}^{2d}$$

by

$$(v \otimes w)u = v\langle w, u \rangle.$$

So Magri's rule is just telling us that

$$K \nabla H_j = J \nabla H_{j+1} \quad \text{for all } j \geq 1 \quad (9)$$

and

$$K \nabla G_j = \begin{cases} J \nabla G_{j-1} + \sum_{i \neq j} J \nabla H_i, & \text{if } j \geq 2; \\ \sum_{i \neq j} J \nabla H_i, & \text{if } j = 1. \end{cases} \quad (10)$$

Yet another way of expressing this is saying that the matrix  $J^{-1} K J^{-1}$  produces the 2-form

$$\sum_{j=1}^d dH_{j+1} \wedge dG_j + \sum_{i < j} c_{ij} dH_i \wedge dH_j.$$

So if we define

$$c_- = \sum_{i < j} c_{ij} dH_i \wedge dH_j$$

we get

**Equivalence 1.1**  $J^{-1}KJ^{-1}$  defines a closed 2-form if and only if  $c_-$  is closed.

Similarly, the form defined by  $(J^{-1}K)^2J^{-1}$  is

$$\sum_{j=1}^d dH_{j+2} \wedge dG_j + c_+$$

where

$$c_+ = \sum_{i < j} (dH_i^+ \wedge dH_j + dH_i \wedge dH_j^+).$$

**Equivalence 1.2**  $(J^{-1}K)^2J^{-1}$  defines a closed 2-form if and only if  $c_+$  is closed.

**Remark.** To prove the formulae for  $J^{-1}KJ^{-1}$  and  $(J^{-1}K)^2J^{-1}$  we also use the fact that  $H_j = p_j(H_1, \dots, H_d)$  for all  $j \geq d + 1$ .

So when can one construct a second bracket for a completely integrable system?

**Hypothesis.** Let  $H_i, i \geq 1$  be any family of commuting functions with respect to a bracket  $J$ . Assume that:

- $G_j, 1 \leq j \leq d$  is a family of functions canonically paired to the  $H_i, 1 \leq i \leq d$ ;
- $H_j = p_j(H_1, \dots, H_d)$  for all  $j \geq d + 1$ ;
- $[c_{ij}(H_1, \dots, H_d)]_{1 \leq i, j \leq d}$  is such that both forms  $c_-$  and  $c_+$  defined above are closed.

**Conclusion.** The skew form  $K$  produced by Magri's rule is compatible with  $J$ . Moreover,  $H_1 \uparrow H_2 \uparrow \dots$

## Further remarks.

- Essentially,  $H_i = \frac{\pi_i}{i}$ ;
- The gauge  $c_{ij}$  can be removed by an appropriate choice of the  $G_j$ 's.

**Big Conclusion.** If  $\dot{x} = J\nabla H(x)$  is an integrable flow on  $\mathbb{R}^{2d}$  with  $d$  independent commuting integrals of motion  $H = H_1, \dots, H_d$ , then Magri's formula provides a compatible  $K$  so that

$$H_1 = \pi_1 \uparrow H_2 = \frac{\pi_2}{2} \uparrow \dots$$

Moreover, by wisely choosing the  $H$ 's and the  $G$ 's, one remove the  $c_{ij}$  and make

$$J^{-1}K\nabla\lambda_i = \lambda_i\nabla\lambda_i, \quad J^{-1}K\nabla\mu_i = \lambda_i\nabla\mu_i.$$