

## Delta and Heaviside functions

$\delta$ -functions are limiting cases of differentiable/continuous functions.

for example,

$$f(x) = \int_{-\infty}^{+\infty} dx' g(x-x') f(x')$$

$$g(x-x') = \frac{1}{2\pi} \int_{-\infty}^{+\infty} d\lambda e^{i\lambda(x'-x)} \rightarrow \delta(x-x')$$

↳ picks off the value at  $x$ .

$\delta$ -function is infinitely narrow and infinitely high such that  $\int \delta(x) dx = 1$

Derivative of  $\delta$ -function?

$$\int_{-\infty}^{+\infty} f(x') \frac{d\delta(x'-x)}{dx'} dx' = \int_{-\infty}^{+\infty} f(x') \frac{d\delta(x'-x)}{dx'} dx'$$

Integrating by parts,

$$\delta(x'-x) f(x') \Big|_{x'=-\infty}^{x'=\infty} - \int_{-\infty}^{+\infty} \delta'(x'-x) \frac{\partial}{\partial x'} f(x') dx'$$

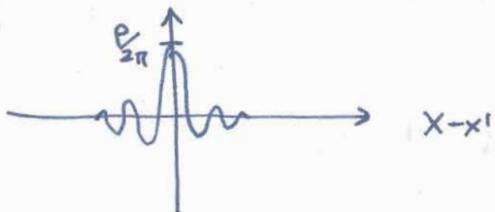
$$= -f'(x)$$

## Representation of $\delta$ -function

$$\frac{1}{2\pi} \int_{-\frac{p_1}{2}}^{\frac{p_1}{2}} d\lambda \cos \lambda(x' - x) = \frac{1}{2\pi} \left. \frac{\sin \lambda(x' - x)}{(x' - x)} \right|_{\lambda = -\frac{p_1}{2}}^{\lambda = \frac{p_1}{2}}$$

$$= \frac{1}{\pi} \frac{\sin \frac{p_1}{2}(x' - x)}{(x' - x)} \xrightarrow[p \rightarrow \infty]{} \delta(x' - x)$$

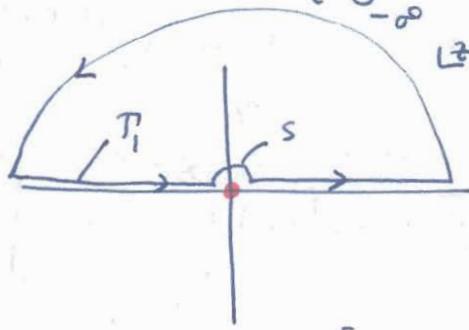
if  $x' \neq x$ ,



check for area:

$$\int_{-\infty}^{+\infty} \frac{1}{\pi} \frac{\sin \frac{p_1}{2} y}{y} dy = \frac{2}{\pi} \int_0^\infty \frac{\sin \frac{p_1 y}{2}}{\frac{p_1 y}{2}} d\left(\frac{p_1 y}{2}\right)$$

note that this integral is  $\text{Im} \left\{ \int_{-\infty}^{+\infty} \frac{e^{iz}}{z} dz \right\}$



$$\oint_C \frac{e^{iz}}{z} dz = 0, \quad \therefore \underbrace{\oint_{T_1} \frac{e^{iz}}{z} dz}_{\text{Indicates the contour } T_1} = - \oint_s \frac{e^{iz}}{z} dz$$

$$\int_{-\infty}^{-\varepsilon} \frac{e^{iz}}{z} dz + \int_{\varepsilon}^{\infty} \frac{e^{iz}}{z} dz$$

$$= - \int_0^\pi d\theta i \varepsilon e^{i\theta} \frac{e^{i\varepsilon e^{i\theta}}}{\varepsilon e^{i\theta}} = i\pi$$

Important  $= \pi$ ,  $\frac{1}{\pi} \int_{-\infty}^{+\infty} \frac{\sin \frac{p_1}{2} y}{y} dy \checkmark = 1$

Another representation:

$$\lim_{p \rightarrow \infty} \frac{P}{\pi(1+p^2x^2)} \longrightarrow \delta(x)$$

for  $x \neq 0$ ,  $\sim \frac{1}{\pi p x^2} \xrightarrow[p \rightarrow \infty]{} 0$

for  $x=0$ ,  $\sim \frac{P}{\pi} \xrightarrow[p \rightarrow \infty]{} \infty$

yet another:  $\delta_p(x) = \begin{cases} \frac{P}{2 \cosh^2 px} & \rightarrow 0 \text{ if } x \neq 0 \\ \rightarrow \infty \text{ if } x=0 \end{cases}$

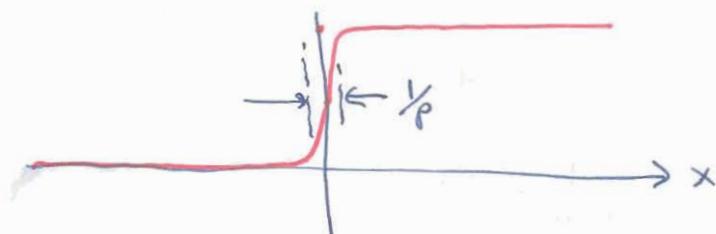
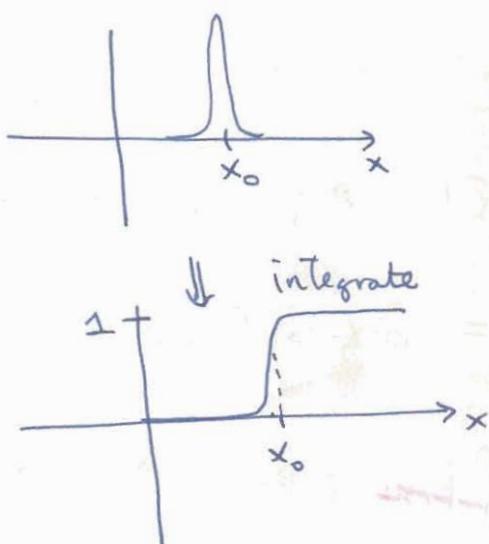
what is the integral of a  $\delta(x)$ ?

$$\int_{-\infty}^x \frac{P}{2 \cosh^2 px} dx = \frac{1}{2} \int_{-\infty}^{px} d\zeta \frac{1}{\cosh^2 \zeta}$$

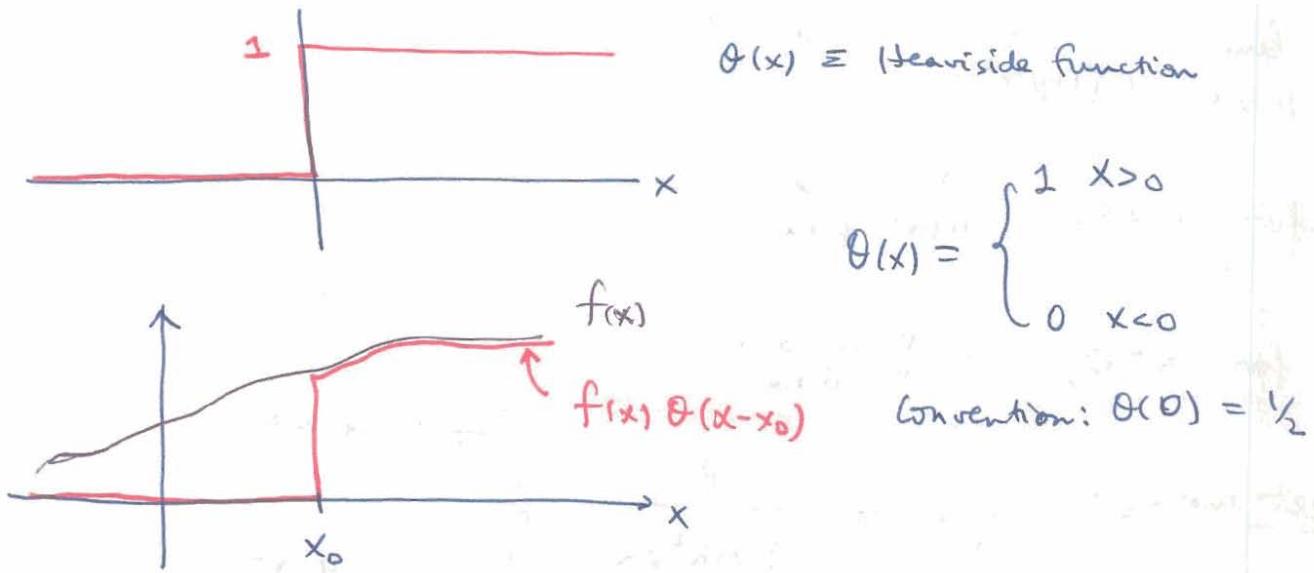
$$= 2 \int_{-\infty}^{px} \frac{d\zeta}{(e^\zeta + e^{-\zeta})^2}, \quad e^\zeta = y, \quad d\zeta = \frac{dy}{y}$$

$$= \int_0^{e^{px}} \frac{dy}{(y^2+1)^2} \Rightarrow \int_0^{e^{2px}} \frac{dy}{(y+1)^2}$$

$$= \int_1^{e^{2px}+1} \frac{dx}{x^2} = -\frac{1}{e^{2px}+1} + 1 = \frac{e^{px}}{e^{px}+e^{-px}} = \frac{1}{2} [1 + \tanh px]$$



As  $p \rightarrow \infty$ , this approaches a "step" function:



### Multidimensional $\delta$ -functions

$\delta(\vec{r} - \vec{r}_0)$  what are units of  $\delta$ -function?

$$= \delta(x - x_0) \delta(y - y_0) \delta(z - z_0)$$

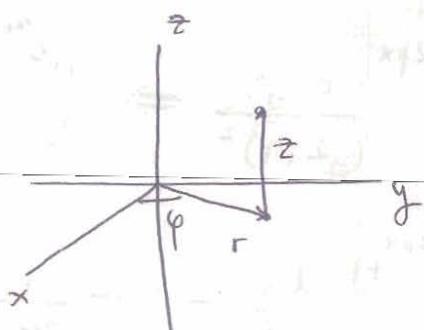
$$\int f(\vec{r}') \delta(\vec{r} - \vec{r}') d^3 r' = f(\vec{r})$$

in cylindrical coordinates

$$d\vec{r} = \hat{r} dr + \hat{\varphi} r d\varphi + \hat{z} dz$$

$$\nabla = \hat{r} \partial_r + \frac{1}{r} \hat{\varphi} \partial_\varphi + \hat{z} \partial_z$$

$$d^2 \vec{r} = r d\varphi dz$$

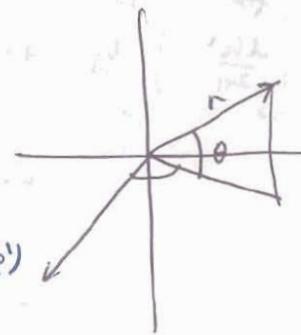


$$\int d^2 \vec{r}' \delta(\vec{r} - \vec{r}') = \int \delta(r - r') r' d\varphi' dz' = 1$$

$$\therefore \delta(\vec{r} - \vec{r}') = \frac{1}{4\pi r'} \delta(r - r') \delta(\varphi - \varphi') \delta(z - z')$$

## Spherical Coordinates

$$d^3\vec{r} = r^2 dr d\sigma = r^2 dr \sin\theta d\theta d\phi$$



$$\therefore \delta(\vec{r} - \vec{r}') = \frac{1}{|\vec{r}|^2} \frac{1}{\sin\theta} \delta(r-r') \delta(\theta-\theta') \delta(\phi-\phi')$$

## Properties of Fourier Transform

### ④ Linearity

$$\int_{-\infty}^{+\infty} [af(x) + bg(x)] e^{-ikx} dx = a F(k) + b G(k)$$

### ④ Scaling

$$\int_{-\infty}^{+\infty} f(ax) e^{-ikx} dx = \int_{-\infty}^{+\infty} f(y) e^{-i\frac{ky}{a}} \frac{dy}{a} = \frac{1}{a} F\left(\frac{k}{a}\right)$$

### ④ Multidimensional

$$\prod_{i=1}^N \int_{-\infty}^{+\infty} f(x_1, x_2, x_3, \dots, x_N) e^{-i(k_1 x_1 + k_2 x_2 + \dots + k_N x_N)} dx_i$$

### ④ Convolution

$$f * g(x) = \int_{-\infty}^{+\infty} g(y) f(x-y) dy = ?$$

$$g(y) = \int_{-\infty}^{+\infty} G(k) e^{iky} \frac{dk}{2\pi}$$

$$= \int_{-\infty}^{+\infty} F(k') e^{ik'(x-y)} \frac{dk'}{2\pi}$$

$$[f * g](x) = \int_{-\infty}^{+\infty} dy \left( \int_{-\infty}^{+\infty} G(k) e^{iky} \frac{dk}{2\pi} \right) \left( \int_{-\infty}^{+\infty} F(k') e^{ik'(x-y)} \frac{dk'}{2\pi} \right) =$$

$$= \int_{-\infty}^{+\infty} \frac{dk}{2\pi} \int_{-\infty}^{+\infty} \frac{dk'}{2\pi} \int_{-\infty}^{+\infty} dy G(k) F(k') e^{iy(k-k')} e^{ik'x}$$

$\int_{-\infty}^{+\infty} dy e^{iy(k-k')} \Rightarrow 2\pi \delta(k-k')$

do the  $dk'$  integral,

$$[f * g](x) = \int_{-\infty}^{+\infty} \frac{dk}{2\pi} G(k) F(k) e^{-ikx} \equiv \text{Fourier transform of } G(k) F(k).$$

Inverting,  $\int_{-\infty}^{+\infty} [f * g](x) e^{-ikx} dx = G(k) F(k)$

Fourier transform of convolution = product of the Fourier transforms of each function.

What is  $\int_{-\infty}^{+\infty} e^{ikx} F(k) G(k) H(k) \frac{dk}{2\pi}$  ?

Parseval's relation

$$\int_{-\infty}^{+\infty} g(y) f(-y) dy = \int_{-\infty}^{+\infty} G(k) F(k) \frac{dk}{2\pi}$$

### ④ Real functions

if  $f(x) \in \mathbb{R}$   $\int_{-\infty}^{+\infty} F(k) e^{ikx} \frac{dk}{2\pi}$  is real,

$$\int_{-\infty}^{+\infty} F(-k) e^{-ikx} \frac{dk}{2\pi} = \int_{-\infty}^{+\infty} F^*(k) e^{-ikx} \frac{dk}{2\pi}$$

∴ for real  $f(x)$   $F^*(k) = F(-k)$

## ④ Derivatives & Fourier transforms

$$\int_{-\infty}^{+\infty} \left( \frac{\partial f}{\partial x} \right) e^{-ikx} dx \xrightarrow{\text{integration by parts}} f(x) e^{-ikx} \Big|_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} -ik f(x) e^{-ikx} dx$$

$$= 0 + ik \int_{-\infty}^{+\infty} f(x) e^{-ikx} dx$$

$$= ik F(k)$$

In general

$$\int_{-\infty}^{+\infty} \left( \frac{\partial f}{\partial x_i} \right) e^{-i(\underbrace{k_1 x_1 + k_2 x_2 + \dots}_{\vec{k} \cdot \vec{x}}} dx = i k_i F(k_1, k_2, \dots, k_N)$$

higher derivatives:

$$\int_{-\infty}^{+\infty} \left[ \frac{d^n f(x)}{dx^n} \right] e^{-ikx} dx = (ik)^n F(k)$$

vector derivatives: gradient  $\vec{\nabla}$ :

$$\int_{-\infty}^{+\infty} (\vec{\nabla} f) e^{-i\vec{k} \cdot \vec{x}} dx_1 dx_2 \dots dx_N$$

$$= i(\hat{x}_1 k_1 + \hat{x}_2 k_2 + \dots + \hat{x}_N k_N) F(\vec{k}) = i \vec{k} F(\vec{k})$$

Laplacian  $\nabla^2$ :

$$\int_{-\infty}^{+\infty} (\nabla^2 f) e^{-i\vec{k} \cdot \vec{x}} d\vec{x} = -(\vec{k} \cdot \vec{k}) F(\vec{k})$$

FT turns derivatives to multiplication w/  $ik \Rightarrow$  solve differential eqns.

## Diffusion-convection equation

$$\frac{\partial c(\vec{x}, t)}{\partial t} = -\vec{v} \cdot \vec{\nabla} c + D \nabla^2 c(\vec{x}, t)$$

Fourier transform eqn:

$$\begin{aligned} & \int_{-\infty}^{+\infty} \left[ \frac{\partial c(\vec{x}, t)}{\partial t} \right] e^{-i\vec{k} \cdot \vec{x}} d\vec{x} \\ &= \frac{\partial \tilde{c}(\vec{k}, t)}{\partial t} = -\vec{v} \cdot (i\vec{k}) \tilde{c}(\vec{k}, t) - D k^2 \tilde{c}(\vec{k}, t) \\ &= -[i\vec{v} \cdot \vec{k} + D k^2] \tilde{c}(\vec{k}, t) \end{aligned}$$

Now solve the ODE in time:

$$\tilde{c}(\vec{k}, t) = \tilde{c}(\vec{k}, 0) e^{-[i\vec{v} \cdot \vec{k} + D k^2] t}$$

what is  $\tilde{c}(\vec{k}, 0)$ ?

$$\tilde{c}(\vec{k}, 0) = \int c(\vec{x}, 0) e^{-i\vec{k} \cdot \vec{x}} d\vec{x}$$

If  $c(\vec{x}, 0)$  starts as a point source,

$$c(\vec{x}, 0) = C_0 \delta(\vec{x}), \text{ and } \tilde{c}(\vec{k}, 0) = C_0$$

$$\tilde{c}(\vec{k}, t) = C_0 e^{-[i\vec{v} \cdot \vec{k} + D k^2] t}$$

find  $C(\vec{x}, t)$  by inverting

$$C(\vec{x}, t) = \int C_0 e^{-[i\vec{v} \cdot \vec{k} + Dh^2]t} \frac{dk}{(2\pi)^3} e^{i\vec{k} \cdot \vec{x}}$$

for 1D,  $C(x, t) = C_0 \int_{-\infty}^{+\infty} \frac{dk}{2\pi} e^{-ivkt - Dh^2 t} e^{ikx}$

complete the square:  $(k+a)^2 = k^2 + 2ak + a^2 =$   
 $k^2 - i \frac{(x-vt)}{Dt} k + a^2$   
 $a = -\frac{i}{2} \frac{(x-vt)}{Dt}$

$$C(x, t) = e^{-Dt a^2} C_0 \int_{-\infty}^{+\infty} \frac{dk'}{2\pi} e^{-Dt k'^2}$$

$$= C_0 e^{-\frac{Dt}{4} \frac{(x-vt)^2}{(Dt)^2}} \cdot \frac{1}{\sqrt{4\pi Dt}}$$

$$= \frac{C_0}{\sqrt{4\pi Dt}} e^{-\frac{(x-vt)^2}{4Dt}} \Rightarrow \text{compare w/ result from probability diffusion.}$$

for 3D  $C(\vec{x}, t) = C_0 \int \frac{d^3 k}{(2\pi)^3} e^{-i\vec{v} \cdot \vec{k} t - Dh^2 t} e^{i\vec{k} \cdot \vec{x}}$

$$= C_0 \int \frac{k^2 dk}{(2\pi)^3} \sin \theta d\phi \underbrace{\cos \theta}_{\text{reference axis } \hat{z}} e^{+i\vec{k} \cdot (\vec{x} - \vec{v} t) - Dh^2 t}$$

$$\int_{-1}^{+1} e^{ik|\vec{x} - \vec{v} t| \cos \theta} d(\cos \theta) = \frac{2 \sin(k|\vec{x} - \vec{v} t|)}{k|\vec{x} - \vec{v} t|}$$

$$\int d\phi \rightarrow 2\pi$$

$$= \frac{2C_0 \sqrt{\pi}}{(2\pi)^3 4(Dt)^{3/2}} \frac{e^{-\frac{|\vec{x} - \vec{v} t|^2}{4Dt}}}{\frac{1}{(4\pi Dt)^{3/2}} e^{-\frac{|\vec{x} - \vec{v} t|^2}{4Dt}}} =$$

$$= \frac{C_0}{k |\vec{x} - \vec{v} t|}$$

$$= \int_0^{2\pi} \frac{k^2 dk}{(2\pi)^3} e^{-Dt k^2} \int_0^{2\pi} \frac{\sin k |\vec{x} - \vec{v} t|}{(2\pi)^3} d\phi$$

## Wave Equation

$$\ddot{\phi} = c^2 \frac{\partial^2 \phi}{\partial x^2}$$

String under tension

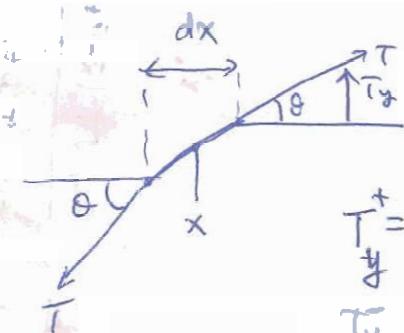
$\rho$  = mass/unit length

Local force balance:

$$dx \rho h(x,t) = \text{acceleration} \times \text{mass}$$

$$= \text{sum of forces}$$

normal



$$T_x^+ = T \sin \theta (x + \frac{dx}{2}) \approx T \theta (x + \frac{dx}{2})$$

$$T_y^+ = T \sin \theta (x - \frac{dx}{2}) \approx T \theta (x - \frac{dx}{2})$$

$$+T_x^+ - T_y^- = \text{sum of normal forces} = T (\theta (x + \frac{dx}{2}) - \theta (x - \frac{dx}{2}))$$

$$\approx +T dx \theta'(x) = +T \frac{d^2 h}{dx^2}$$

$$\therefore \ddot{\phi}(x,t) = +T \frac{d^2 h}{dx^2}$$

$$\ddot{h}(x,t) = \left(\frac{T}{\rho}\right) \frac{d^2 h}{dx^2} \quad \text{wave equation}$$

$$\frac{T}{\rho} = c^2 = \frac{\text{Force}}{\text{mass/length}} = \frac{\text{mass} \cdot \text{length}}{(\text{time})^2 \text{ mass}} \frac{\text{length}}{\text{mass}} = \frac{L^2}{T^2} = c^2$$

c = wave speed.