

Isomer selective infrared spectroscopy of neutral metal clusters

André Fielicke,¹ Christian Ratsch,^{1,2} Gert von Helden,¹ and Gerard Meijer¹

¹*Fritz-Haber-Institut der Max-Planck-Gesellschaft,
Faradayweg 4-6, D-14195 Berlin, Germany*

²*Department of Mathematics, UCLA and California NanoSystems Institute,
UCLA, Los Angeles, CA 90095-1555, USA*

(Dated: January 18, 2005)

Abstract

We report experimental infrared spectra of neutral metal clusters in the gas phase. Multiple photon dissociation of the argon complexes of niobium clusters is used to obtain vibrational spectra in the 80–400 cm^{-1} region. The observed spectra for Nb_9Ar_n ($n = 1 - 4$) are different for different values of n . This is explained by the presence of two isomers of Nb_9 that have different affinities towards Ar and the isomer specific infrared spectra are obtained. The structures of the isomers are determined by comparing the observed spectra with the outcome of density-functional theory calculations.

PACS numbers: 36.40.Mr, 61.46.+w

The knowledge on the geometric and electronic structure of a cluster forms the basis for the understanding of its physical and chemical properties. Geometric structure determination of metal clusters in the gas phase is difficult and for systems where isomers are present, the structure determination is even more challenging. For metal clusters, the coexistence of different isomeric structures has been experimentally revealed using mainly two methods: i) investigation of the kinetics with reactant molecules, where bi- or multiple-exponential reaction kinetics are signatures for the presence of isomers with different reactivities [1–5], or ii) (ion) mobility measurements, where isomers with different collision cross sections can be separated and detected *via* their different arrival time distributions [6–8]. Further evidence for metal cluster isomers is obtained from the measurements of magnetic moments [9], from unimolecular photodissociation dynamics [10], and from photoelectron spectroscopy [11]. Very recently, optical spectroscopy on matrix deposited silver clusters allowed for the identification of the presence of silver cluster isomers [12].

In principle, vibrational spectroscopy can be used to obtain structural information on metal clusters, but until recently, this approach was limited to selected small clusters that have mainly been investigated *via* matrix isolation spectroscopy [13] or to clusters where vibrational structure in electronic excitation spectra has been found [14–18]. For cationic metal clusters, we have recently demonstrated that far-infrared (vibrational) resonance enhanced multiple photon dissociation (FIR-MPD) spectroscopy of their complexes with rare gas atoms is a suitable method to obtain their vibrational spectra, and to thus get direct information on their geometric structures [19]. Here we report on FIR-MPD spectroscopic studies of neutral metal clusters. We are focussing in particular on Nb₉, to demonstrate that isomer specific vibrational spectra can be obtained. The structures of the isomers are determined by comparison of the observed FIR-MPD spectra to calculated IR absorption spectra as obtained from density-functional theory calculations.

The experimental set-up has been described elsewhere [19] and only a brief description is given here. Argon complexes of niobium clusters are produced in a pulsed laser ablation cluster source that is partly cooled to ≈ 80 K. As carrier gas, a mixture of 20–40% Argon in Helium is used. The binding of the rare gas atoms to neutral metal clusters is much weaker than the binding to cationic metal clusters and complex formation is, therefore, by far more sensitive to source parameters such as timings, partial pressures, and fluence of the vaporization laser. The cluster complexes are detected in a time-of-flight mass spectrometer

after ionization with either an ArF excimer laser (6.42 eV/photon) or a frequency-doubled dye laser. As the ionization potentials (IPs) for niobium clusters are typically in the 4.6–5.5 eV range [20, 21] and as the IPs of their Ar complexes will be slightly lower, the 6.42 eV photon energy allows for one photon ionization of all clusters and their Ar complexes. The fluence of the ArF laser is kept sufficiently low to prevent absorption of multiple UV photons that would lead to fragmentation of the weakly bound complexes. With this method, niobium cluster complexes containing up to four Ar atoms have been detected. Alternatively, the frequency doubled output of a dye laser is used for the more selective ionization near the threshold (see below). The molecular beam is overlapped with a counter-propagating IR laser beam emitted by the Free Electron Laser for Infrared eXperiments (FELIX) [22]. When the IR light is resonant with a transition in the cluster complex, the complex can absorb several photons and subsequently evaporate off one or more rare gas atoms. The resulting abundance changes of the rare gas complexes are recorded as a function of the IR frequency. Details on the normalization of the spectra to the laser intensity and on the conversion of the depletion spectra into absorption spectra are described elsewhere [19].

Figure 1 shows parts of the typical mass spectra in the range of Nb_9Ar_n ($n = 0 - 4$). In the spectrum in fig. 1(a) the intensities of the rare gas complexes are usually a factor of ≈ 50 lower than those of the bare metal clusters. Additional peaks can be assigned to oxide and carbide contaminations. For $n = 3$ and 4 the peaks of the complexes overlap with those of the contaminants. However, this does not lead to interferences in the IR spectra of Nb_9Ar_n since only the rare gas complexes can be dissociated.

Far-infrared absorption spectra of free neutral niobium clusters are measured *via* multiple photon dissociation of the Nb_mAr_n complexes in the 80–400 cm^{-1} region. Using ArF ionization for detection, spectra are recorded for clusters containing five to more than 20 Nb atoms and different numbers of Ar atoms. These results will be reported elsewhere, but important for the discussion here is that for almost all of the clusters the observed FIR-MPD spectra are invariant with the number of attached Ar atoms. Only small shifts of the absorption peaks (on the order of 1–3 cm^{-1}) and slight broadenings are observed for a given Nb_m cluster with an increasing number of attached Ar atoms.

A special case are the Nb_9Ar_n ($n = 1 - 4$) complexes. Their FIR-MPD spectra are shown in fig. 2. For Nb_9Ar_n ($n = 1 - 4$) absorption features are only found in the 180–300 cm^{-1} region. In the FIR-MPD spectrum of Nb_9Ar four main bands are found at $\approx 195, 202, 252$

(P3), and at 280 cm^{-1} (P4). In addition, a minor peak can be identified around 215 cm^{-1} (P2) and there is a shoulder on the low frequency side of the peak P3. With increasing number of Ar atoms in the complex, the double peak around 200 cm^{-1} merges into a single peak (P1), the peak P2 as well as the shoulder on P3 gain in intensity whereas the originally dominating peaks decrease in intensity.

The strong dependence of the FIR-MPD spectra of Nb_9Ar_n on the number of Ar atoms can be explained by the presence of (at least) two isomers of Nb_9 that have different affinities towards the Ar atoms. To be more precise, the relative affinity of these Nb_9 isomers towards Ar must gradually change with the number n of Ar atoms. This then leads to a change in the relative contribution of the Nb_9 isomers to the spectra of the Nb_9Ar_n complexes as a function of n , and can explain the observations. The presence of (at least) two isomers of Nb_9 has been concluded earlier from the observed biexponential reaction kinetics with hydrogen or nitrogen [1–3] and the different IPs for the isomers have been experimentally determined to be 4.92 eV and 5.20 eV [20]. Optical absorption spectra for Nb_9 have been recorded using photodissociation of their rare gas cluster complexes [23, 24] and the photoionization threshold behavior of these complexes has been investigated as well [25]. Also in these studies, differences in the propensities to bind Ar atoms have been found for the isomers of Nb_9 .

Alternatively, one might explain the observed changes in the FIR-MPD spectra of Nb_9Ar_n as a function of n by assuming that there is only one Nb_9 isomer but that the oscillator strengths of the different vibrational modes are differently affected by complexing the metal cluster with an increasing number of Ar atoms. This explanation seems highly unlikely, however, as we have not observed any significant influence of the rare gas atoms on the FIR-MPD spectra for other clusters.

The difference in the IPs of the two Nb_9 isomers enables the selective ionization of the isomer with the lowest IP using near threshold ionization. The onset of ionization for Nb_9Ar is shown in the inset of fig. 1(b). An ionization threshold of $4.845 \pm 0.002\text{ eV}$ (under field-free conditions) is found, in agreement with previous findings [25]. Figure 1(b) shows a mass spectrum that has been obtained using near threshold ionization with 250 nm radiation. Under the assumption that the Ar induced shift of the IPs is approximately the same for both isomers ($\approx 0.07\text{ eV}$), then the Ar complex of the second isomer should have an IP of $\approx 5.13\text{ eV}$, well above the photon energy at 250 nm.

The FIR-MPD spectrum of Nb_9Ar measured with near threshold ionization is shown in fig. 2(b). Only the Ar complex of one Nb_9 isomer can be ionized and therefore this spectrum can be assigned to the isomer with the lowest IP (isomer **1**). This FIR-MPD spectrum is quite similar to the FIR-MPD spectrum of Nb_9Ar measured with ArF ionization. Therefore, we conclude that isomer **1** has a higher propensity for binding a single Ar atom than isomer **2**. The observed changes in the FIR-MPD spectra indicate that with increasing number of Ar atoms in the complex the abundance of isomer **1** decreases relative to the abundance of isomer **2**.

The spectra of the individual isomers can be extracted from the spectra of the isomeric mixture by iterative subtractive deconvolution. As a starting point, it is assumed that the spectrum of Nb_9Ar is dominated by isomer **1**, whereas isomer **2** dominates the spectrum of Nb_9Ar_4 . The deconvolution has been applied after smoothing the data with a three point fast Fourier transform filter. No information from the FIR-MPD spectrum as obtained after selective ionization has been used in this approach. The thus obtained isomer selective FIR-MPD spectra are shown in fig. 3(a). Combinations of these spectra in the appropriate intensity ratio are shown overlayed with the observed spectra in fig. 2(a). The intensity ratio changes from 1:0.3 to 1:1.9 when n increases from one to four, but these numbers only indicate the relative abundance ratio of the two isomers as a function of n since the absolute IR absorption intensities and the ionization efficiencies are not known. The FIR-MPD spectrum of isomer **1** as extracted from the spectra of the isomeric mixture is very similar to the spectrum that is obtained using selective ionization, shown once more in fig. 3(b).

Density-functional theory (DFT) calculations have been performed to find the energetically most preferred structures for Nb_9 , and their vibrational spectra have been calculated. We have used the DMOL³ code [26] and the generalized gradient approximation (GGA) [27] for the exchange correlation functional. More details of the computational procedure will be given elsewhere [28]. As a check, for Nb_2 this approach yields a calculated vibrational frequency of 445 cm^{-1} , close to the experimental value of 425 cm^{-1} [31]. We tested a large number of different geometric structures for Nb_9 , including some that have been suggested previously in the literature [29, 30]. Moreover, we tested a number of geometric distortions of these structures, and calculated the energy and the IR absorption spectra for different spin states. The vibrational frequencies presented here are not scaled.

In Figure 3(c) the results of the calculations for three different structures of Nb₉ are shown. The lowest energy structure **A** can be described as a distorted triply capped trigonal prism [32]. Structure **B** is only 0.15 eV higher in energy and is a non-planar hexagon with a dimer above and an atom below the hexagon. The next stable structure **C**, a distorted triply capped octahedron, is noticeably higher in energy (0.82 eV). All these structures correspond to a doublet electronic state and structures with higher multiplicity are found to be significantly higher in energy. The IR spectrum of structure **B** nicely agrees with the experimental spectrum of isomer **1**; the band positions as well as the IR intensities are reproduced rather well. There is considerably less agreement between the observed spectrum of isomer **2** and the calculated IR spectra of either structure **A** or structure **C**. Nevertheless, the comparison of the spectra suggests that isomer **2** may have structure **A**; it is the energetically lowest structure, and there is some agreement between the spectra. The two calculated bands at 189 cm⁻¹ and at 211 cm⁻¹ can be related to features in the experimental spectrum. Below 200 cm⁻¹ the experimental spectrum is noisy and not well-resolved, preventing a firm conclusion on the presence or absence of the band calculated at 168 cm⁻¹ in the experimental spectrum. The calculations further predict IR bands at 241, 253, 278, and 279 cm⁻¹. This is in agreement with the band positions in the experimental spectrum, but the intensities do not match. One has to keep in mind, however, that the experimental spectra are obtained by multiple photon absorption followed by dissociation, and that intensities in FIR-MPD spectra can differ considerably from those in linear FIR absorption spectra. Further evidence for the assignments proposed here is derived from the calculation of the (adiabatic) ionization energies. We find values of 5.13 eV for structure **A** and 4.98 eV for structure **B** that are in good agreement with the experimental values. The assignment of isomer **1**, the isomer with the lowest IP, to structure **B** seems quite certain. The assignment for isomer **2** is less certain, and we cannot exclude that this isomer has yet another structure, that has not been found in the calculations thus far.

Finally, the question remains to be answered why these two Nb₉ isomers show different Ar binding properties. The calculations show that structure **B** has a dipole moment of 0.70 D, considerably higher than the 0.12 D dipole moment of structure **A**. This can explain the observed preferred binding of the first (two) Ar atom(s) to isomer **1**. However, the large dipole moment of structure **B** is mainly caused by the presence of the single Nb atom below the hexagon unit that carries a significant positive partial charge. It is likely, therefore, that

after the first (two) Ar atoms(s) are attached to this side of the Nb₉ isomer, the next Ar atoms either have to attach to a less favorable side of the complex or they attach to the same side of the complex where they then experience a strongly shielded dipole. In either case, this can well explain the observed (relative) decrease in affinity for Ar atoms with increasing number of Ar atoms for isomer **1**.

In conclusion, we have shown that vibrational spectra of neutral metal clusters can be measured via multiple photon dissociation spectroscopy of their rare gas complexes using tunable radiation in the far-infrared. In special cases even isomer specific spectra can be obtained. Together with theoretical studies, these spectra enable the structure of metal cluster isomers to be unravelled.

Acknowledgments

We acknowledge stimulating discussions with D.M. Rayner and M. Scheffler. This work is part of the research program of the "Stichting voor Fundamenteel Onderzoek der Materie" (FOM), which is financially supported by the "Nederlandse organisatie voor Wetenschappelijk Onderzoek" (NWO). Financial support from the EU IHP Research Training Network (Delayed Ionization and Competing Cooling Mechanisms in Atomic Clusters) is gratefully acknowledged. We also thank Accelrys for providing the DMOL³ software to the California NanoSystems Institute.

-
- [1] M. R. Zakin, R. O. Brickman, D. M. Cox, and A. Kaldor, *J. Chem. Phys.* **88**, 3555 (1988).
 - [2] Y. Hamrick *et al.*, *J. Chem. Phys.* **88**, 4095 (1988); *J. Phys. Chem.* **93**, 6494 (1989).
 - [3] A. Bérces *et al.*, *J. Chem. Phys.* **108**, 5476 (1998).
 - [4] J. Ho, L. Zhu, E. K. Parks, and S. J. Riley, *J. Chem. Phys.* **99**, 140 (1993).
 - [5] C. Berg *et al.*, *Chem. Phys.* **262**, 143 (2000).
 - [6] G. von Helden, M.-T. Hsu, P. R. Kemper, M. T. Bowers, *J. Chem. Phys.* **95**, 3835 (1991).
 - [7] M. F. Jarrold and J. E. Bower, *J. Chem. Phys.* **98**, 2399 (1993); A. A. Shvartsburg and M. F. Jarrold, *Phys. Rev. A* **60**, 1235 (1999).
 - [8] P. Weis, T. Bierweiler, E. Vollmer, and M. M. Kappes, *J. Chem. Phys.* **117**, 9293 (2002).

- [9] J. P. Bucher, D. C. Douglass, and L. A. Bloomfield, *Phys. Rev. Lett.* **66**, 3052 (1991).
- [10] M. Vogel, K. Hansen, and L. Schweikhard, *Int. J. Mass Spectrom.* **233**, 117 (2004).
- [11] H. Kietzmann *et al.*, *Phys. Rev. Lett.* **77**, 4528 (1996); *J. Chem. Phys.* **109**, 2275 (1998).
- [12] C. Sieber *et al.*, *Phys. Rev. A* **70**, 041201 (2004).
- [13] J. R. Lombardi and B. Davis, *Chem. Rev.* **102**, 2431 (2002).
- [14] M. F. Jarrold and K. M. Creegan, *Chem. Phys. Lett.* **166**, 116 (1990).
- [15] K. J. Taylor *et al.*, *J. Chem. Phys.* **93**, 7515 (1990).
- [16] G. F. Ganteför, D. M. Cox, and A. Kaldor, *J. Chem. Phys.* **96**, 4102 (1992); H. Handschuh, G. Ganteför, and W. Eberhardt, *Rev. Sci. Instrum.* **66**, 3838 (1995).
- [17] D.S. Yang, A.M. James, D.M. Rayner, and P.A. Hackett, *Chem. Phys. Lett.* **231**, 177 (1994).
- [18] T. P. Marcy and D. G. Leopold, *Int. J. Mass Spectrom.* **196**, 653 (2000).
- [19] A. Fielicke *et al.*, *Phys. Rev. Lett.* **93**, 023401 (2004).
- [20] M. B. Knickelbein and S. Yang, *J. Chem. Phys.* **93**, 1476 (1990); **93**, 5760 (1990).
- [21] K. Athanassenas, *et al.*, *Chem. Phys. Lett.* **213**, 105 (1993).
- [22] D. Oepts, A. F. G. van der Meer, and P. W. van Amersfoort, *Infrared Phys. Technol.* **36**, 297 (1995).
- [23] M. B. Knickelbein and W. J. C. Menezes, *Phys. Rev. Lett.* **69**, 1046 (1992); W. J. C. Menezes and M. B. Knickelbein, *J. Chem. Phys.* **98**, 1856 (1993).
- [24] B. A. Collings, K. Athanassenas, D. M. Rayner, and P. A. Hackett, *Z. Phys. D* **26**, 36 (1993).
- [25] M. B. Knickelbein and W. J. C. Menezes, *Chem. Phys. Lett.* **184**, 433 (1991); *J. Phys. Chem.* **96**, 6611 (1992).
- [26] B. Delley, *J. Chem. Phys.* **92**, 508 (1990).
- [27] J. P. Perdew, K. Burke, and M. Ernzerhof, *Phys. Rev. Lett.* **77**, 3865 (1996).
- [28] C. Ratsch *et al.*, *J. Chem. Phys.* (in press).
- [29] H. Grönbeck, A. Rosén, and W. Andreoni, *Phys. Rev. A* **58**, 4630 (1998).
- [30] V. Kumar and Y. Kawazoe, *Phys. Rev. B* **65**, 125403 (2002).
- [31] A. M. James, P. Kowalczyk, R. Fournier, B. Simard, *J. Chem. Phys.* **99**, 8504 (1993).
- [32] We have marked in fig. 3(c) the core parts of the structures for better recognizability of the structural elements. Structural descriptions in terms of other elements are equally possible.

Figure captions

FIG 1: Mass spectra in the mass range of Nb_9Ar_n ($n = 0 - 4$) recorded using (a) 6.42 eV photons, and (b) 4.96 eV photons for the ionization. The conditions under which spectrum (b) is recorded are optimized for Nb_9Ar . At 4.96 eV (250 nm) the Ar complexes of only one isomer of Nb_9 can be ionized. A measurement of the ionization threshold is shown in the inset.

FIG 2: (a) Far-IR multiple photon dissociation (FIR-MPD) spectra of Nb_9Ar_n ($n = 1-4$). The neutral clusters are ionized with 6.42 eV photons. The major features in the spectra are indicated with **P1-P4**. Reconstructed spectra of mixtures of the two isomers in the given ratios (intensities of isomer **1** to isomer **2**) are shown overlayed with the experimental spectra. (b) FIR-MPD spectrum of Nb_9Ar as obtained by ionization using 250 nm radiation.

FIG 3: Isomer selective FIR-MPD spectra. Panel a) shows the isomer selective FIR-MPD spectra as extracted out of the spectra of Nb_9Ar_n (shown in fig. 2) for each of the two Nb_9 isomers **1** and **2**. Panel b) shows the FIR-MPD spectrum of isomer **1** as obtained by selective ionization of this isomer with 250 nm radiation. Panel c) shows the calculated IR absorption spectra for three low energy isomers of Nb_9 labelled **A**, **B**, and **C**. The calculated spectra are folded with a Gaussian line shape function with a width (FWHM) of 5 cm^{-1} .

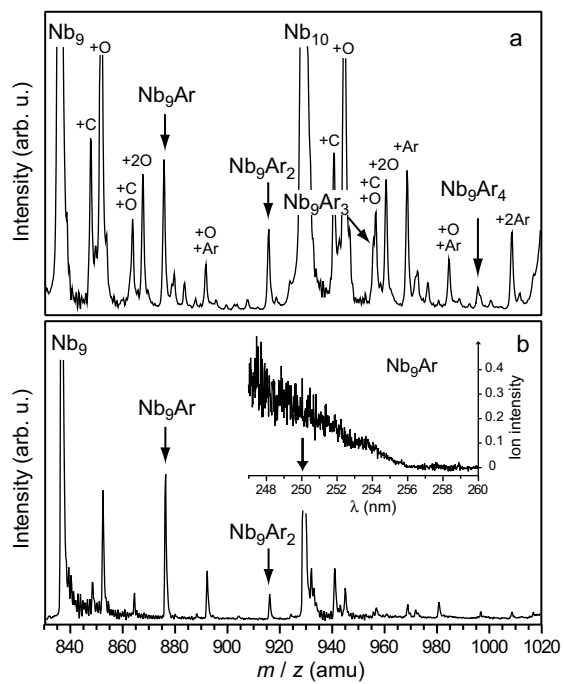


FIG. 1:

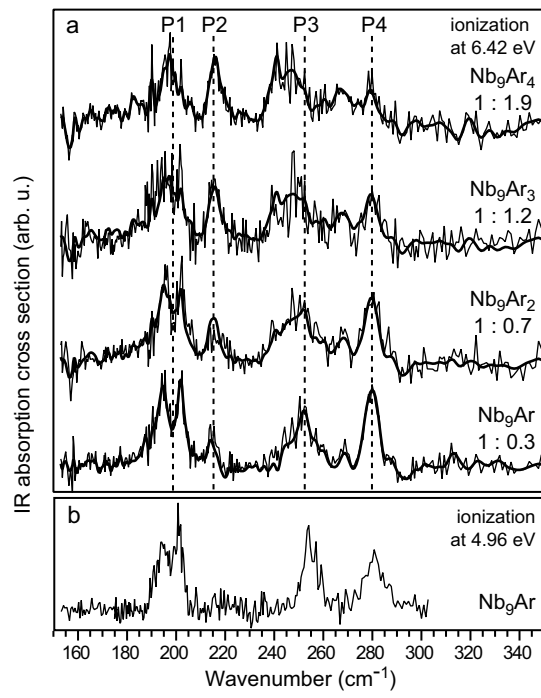


FIG. 2:

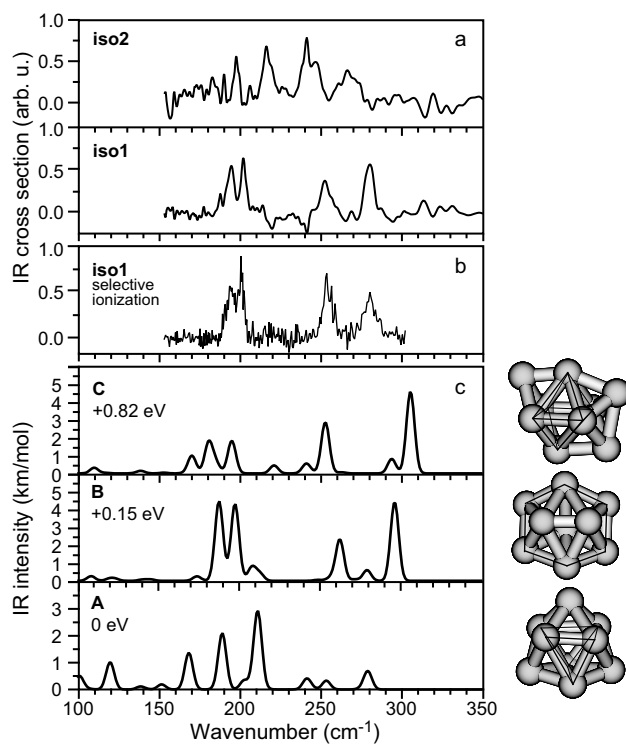


FIG. 3: