# LOCAL PROPERTIES OF SELF-DUAL HARMONIC 2-FORMS ON A 4-MANIFOLD

## KO HONDA

ABSTRACT. We prove a Moser-type theorem for self-dual harmonic 2-forms on closed 4manifolds, and use it to classify local forms on neighborhoods of singular circles on which the 2-form vanishes. Removing neighborhoods of the circles, we obtain a symplectic manifold with contact boundary – we show that the contact form on each  $S^1 \times S^2$ , after a slight modification, must be one of two possibilities.

## 1. INTRODUCTION

This paper is a study of generic self-dual (SD) harmonic 2-forms near their zero sets. Let  $M^4$  be a closed, oriented 4-manifold with  $b_2^+(M) > 0$ . Then one can show that, for a pair  $(\omega, g)$  consisting of a generic metric g and a generic self-dual harmonic 2-form  $\omega$  with respect to g,  $\omega$  is a transverse section of  $\bigwedge_g^+ \to M$  (i.e.,  $\omega$  is transverse to the zero section). Here  $\bigwedge_g^+$  is the self-dual subbundle of  $\bigwedge^2 T^*M \to M$  whose fiber over a point  $p \in M$  is  $\bigwedge_g^+(p) = \{\omega \in \bigwedge_p^2 T^*M | *_g \omega = \omega\}$ . For a generic  $(\omega, g)$ , the zero set C of  $\omega$  is therefore a disjoint union of embedded circles. Now, since  $\omega \wedge \omega = \omega \wedge *\omega$ ,  $\omega$  is nondegenerate at p if and only if  $\omega(p) \neq 0$ . That is,  $\omega$  is symplectic away from C – a union of circles – and is identically 0 on C. For more details, consult [3] or [2].

The interest in self-dual harmonic 2-forms on closed 4-manifolds comes, to a large extent, from our attempt to understand which closed 4-manifolds have symplectic structures. We therefore view the zero set C of  $\omega$  as an obstruction to the existence of a symplectic structure on M, and will sometimes refer to the self-dual harmonic forms as *singular symplectic forms*.

We briefly outline the contents of the paper. In Section 2, we introduce an almost complex structure J which is naturally associated to our singular symplectic form  $\omega$  and metric g. Section 3 is devoted to a discussion of a version of Moser's theorem (Theorem 2) which applies to our singular symplectic forms. In Section 4, we use the Moser-type theorem to classify local normal forms for the singular symplectic forms near an  $S^1$ , with an eye towards global results, and in the last section we discuss the induced contact structures on the boundaries of  $N(S^1)$ .

It turns out that most of the theorems described in paper were already known to various researchers, but never published. I hope this manuscript fills a gap in the literature, especially for readers of Taubes' papers [6, 7, 8, 9], which contain an analysis of *J*-holomorphic

Date: Revised January 13, 2004.

<sup>1991</sup> Mathematics Subject Classification. Primary 57R17; Secondary 57R57.

Key words and phrases. harmonic form, symplectic, 4-manifold.

curves with boundary on C, and the relationship to the nonvanishing of the Seiberg-Witten invariants.

## 2. Almost complex structures

Observe first that we can define an almost complex structure J on M - C, where C is the zero set of  $\omega$ .

**Proposition 1.** If  $\omega$  is a self-dual harmonic 2-form which is nondegenerate on M-C, then there exists a unique almost complex structure J compatible with  $\omega$  and  $\tilde{g}$  on M-C, where  $\tilde{g}$  is conformally equivalent to g.

*Proof.* Given a metric g, any 2-form  $\omega$  can be written pointwise as

$$\omega = c_1 e_1 e_2 + c_2 e_3 e_4,$$

where  $(e_1, ..., e_4)$  is an oriented orthonormal basis for  $T_p^*M$  at a point  $p \in M$ . Since  $\omega$  is self-dual,  $c_1 = c_2$ . Hence,

$$\omega = c(e_1e_2 + e_3e_4).$$

We claim this c is well-defined and smooth on M - C, up to sign. Since  $\frac{1}{2}\omega \wedge \omega = c^2 e_1 \dots e_4 = c^2 dv_g$ , with  $dv_g$  the volume form for g, we find that  $c^2$  is determined by  $\omega$  and g. Taking advantage of M - C being connected, we may fix c on all of M - C so that c > 0.

We then set  $J: e_1 \mapsto e_2, e_2 \mapsto -e_1, e_3 \mapsto e_4, e_4 \mapsto -e_3$ . This definition is equivalent to the following: Let  $\tilde{g} = cg$ , and define J such that  $\tilde{g}(x, y) = \omega(x, Jy)$ . Hence we see that if there is a J compatible with  $\omega$  and  $\tilde{g}$ , it must be unique. Thus J is compatible with  $\omega$  and  $\tilde{g} = cg$  on M - C.

Observe that  $\omega$  is defined on all of M and is zero on C,  $\tilde{g}$  can be defined on all of M and is zero on C, but is not smooth on C, while J is defined only on M - C.

## 3. Moser argument for self-dual harmonic 2-forms

Let  $\{\omega_t\}, t \in [0, 1]$ , be a family of self-dual harmonic 2-forms on M, where  $\omega_t$  is a transverse section of the corresponding  $\bigwedge_{a_t}^+$  for each t, and the following hold:

- (i)  $[\omega_t] \in H^2(M; \mathbf{R})$  is constant.
- (ii) The sets  $C_t = \{x \in M | \omega_t(x) = 0\}$  are all  $S^1$ 's; hence via an isotopy, we may assume that  $C = C_t$  is a fixed  $S^1$ .
- (iii)  $[\omega_t] \in H^2(M, C; \mathbf{R})$  does not vary with t.

For simplicity, we have assumed in (ii) that the zero set C is a single circle. The result also holds if there are more circles, provided their number remains constant with t. If C is contractible, then (iii) is equivalent to the following:

(iii') Let  $\Omega$  be an oriented surface with  $\partial \Omega = C$ . Then  $\int_{\Omega} \omega_t$  does not vary with t.

Then we have the following:

**Theorem 2.** There exists a 1-parameter family of  $C^0$ -homeomorphisms of M, which is smooth away from C and takes  $(M - C, \omega_0) \xrightarrow{\sim} (M - C, \omega_1)$  symplectically.

This generalizes the classical Moser's theorem:

**Theorem 3** (Moser). Let  $\{\omega_t\}$  be a family of symplectic forms on a closed manifold M. Provided  $[\omega_t] \in H^2(M; \mathbf{R})$  is fixed, there is a 1-parameter family of diffeomorphisms  $\phi_t$  such that  $\phi_t^* \omega_t = \omega_0$ .

Proof of Theorem 3. Let  $\eta_t$  be a 1-parameter family of 1-forms such that  $\frac{d\omega_t}{dt} = d\eta_t$ . (The nontrivial part of this proof is to construct a smooth family  $\eta_t$  using Hodge theory.) Thus, if we define  $X_t$  such that  $i_{X_t}\omega_t = -\eta_t$ , then  $\mathcal{L}_{X_t}\omega_t = (i_{X_t} \circ d + d \circ i_{X_t})\omega = -d\eta_t$ , which, integrated, gives a 1-parameter family  $\phi_t$  such that  $\phi_t^*\omega_t = \omega_0$ .

Proof of Theorem 2. The point here is to find a suitable  $\eta_t$  satisfying  $\frac{d\omega_t}{dt} = d\eta_t$ , and  $\eta_t|_C = 0$ "up to first order" near C.

**Claim.** There exists a 1-form  $\tilde{\eta}_t$  satisfying (1)  $\frac{d\omega_t}{dt} = d\tilde{\eta}_t$  and (2)  $i^*\tilde{\eta}_t$  is exact, where  $i: C \to M$  is the inclusion.

*Proof of Claim.* This follows from condition (iii) and the relative cohomology sequence:

 $H^1(M; \mathbf{R}) \to H^1(C; \mathbf{R}) \xrightarrow{\delta} H^2(M, C; \mathbf{R}) \to H^2(M; \mathbf{R}).$ 

In the de Rham setting,  $\delta$  is given as follows: Given a class  $[\alpha] \in H^1(C; \mathbf{R})$ , represented by a closed 1-form  $\alpha$  on C, we extend  $\alpha$  to (a not necessarily closed) 1-form  $\tilde{\alpha}$  on M. Then let  $\delta[\alpha] = [d\tilde{\alpha}] \in H^2(M, C; \mathbf{R}).$ 

Now choose any smooth family  $\tilde{\eta}_t$  satisfying  $\frac{d\omega_t}{dt} = d\tilde{\eta}_t$ . Since  $\left[\frac{d\omega_t}{dt}\right] = 0 \in H^2(M, C; \mathbf{R})$ , if  $[i^*\tilde{\eta}_t] \neq 0 \in H^1(C; \mathbf{R})$ , we can kill this class by adding to  $\tilde{\eta}_t$  a closed 1-form on M which represents a class in  $H^1(M; \mathbf{R})$  and which maps to  $[i^*\tilde{\eta}_t] \in H^1(C; \mathbf{R})$ . This modification can be performed smoothly with respect to t, using Hodge theory.

Now fix an  $\tilde{\eta}_t$  as in the Claim. Then there exists a family  $f_t$  of functions on C such that  $i^* \tilde{\eta}_t = df_t$ . Our goal is to extend  $f_t$  to a function on M such that  $\tilde{\eta}_t = df_t$  "up to first order" near C.

In order to extend  $f_t$  to a neighborhood N(C) of C, first observe that there is only one orientable rank 3 bundle over  $S^1$  ( $\pi_1(BSO(3)) = 0$  implies  $S^1 \to BSO(3)$  is homotopically trivial) and hence  $N(C) \simeq C \times D^3$ . Choose coordinates ( $\theta, x_1, x_2, x_3$ ) such that  $d\theta, dx_1, dx_2, dx_3$  at ( $\theta, 0$ ) are orthonormal.

Setting

$$f_t(\theta, x_1, x_2, x_3) = f_t(\theta, 0) + \sum_i \tilde{\eta}_i(\theta, 0) x_i + \frac{1}{2} \sum_{i,j} \frac{\partial \tilde{\eta}_i}{\partial x_j}(\theta, 0) x_i x_j$$

on N(C), where  $\tilde{\eta}_t = \tilde{\eta}_{\theta} d\theta + \sum_i \tilde{\eta}_i dx_i$ , we have

$$df_t(\theta, x_1, x_2, x_3) = \frac{\partial f_t}{\partial \theta}(\theta, 0)d\theta + \sum_i \frac{\partial \tilde{\eta}_i}{\partial \theta}(\theta, 0)x_i d\theta \\ + \sum_i \tilde{\eta}_i(\theta, 0)dx_i + \frac{1}{2}\sum_{i,j} \frac{\partial \tilde{\eta}_i}{\partial x_j}(\theta, 0)(x_i dx_j + x_j dx_i)$$

up to first order in the  $x_i$ 's. Now observing that

(1) 
$$\frac{\partial f}{\partial \theta}(\theta, 0) = \tilde{\eta}_{\theta}(\theta, 0),$$

(2) 
$$d\tilde{\eta}_t(\theta, 0) = 0,$$

and that Equation 2 gives

$$\frac{\partial \tilde{\eta}_{\theta}}{\partial x_{i}}(\theta, 0) = \frac{\partial \tilde{\eta}_{i}}{\partial \theta}(\theta, 0),$$
$$\frac{\partial \tilde{\eta}_{i}}{\partial x_{j}}(\theta, 0) = \frac{\partial \tilde{\eta}_{j}}{\partial x_{i}}(\theta, 0),$$

we obtain

$$df_t(\theta, x) = \left( \tilde{\eta}_{\theta}(\theta, 0) + \sum_i \frac{\partial \tilde{\eta}_{\theta}}{\partial x_i}(\theta, 0) x_i \right) d\theta + \sum_i \left( \tilde{\eta}_i(\theta, 0) + \sum_j \frac{\partial \tilde{\eta}_i}{\partial x_j}(\theta, 0) x_j \right) dx_i = \tilde{\eta}_{\theta}(\theta, x) d\theta + \sum_i \tilde{\eta}_i(\theta, x) dx_i$$

up to first order in x.

Damping  $f_t$  out to 0 outside N(C), we arrive at  $\eta_t = \tilde{\eta}_t - df_t$ . Finally, we obtain the vector field  $X_t$  such that  $i_{X_t}\omega_t = -\eta_t$ .  $X_t$  will then give rise to a 1-parameter family of symplectomorphisms, away from C, once we establish that  $X_t \to 0$  rapidly enough as  $p \in M$  approaches C.

On N(C),

(3)  

$$\omega_t = L_1(\theta, x)(d\theta dx_1 + dx_2 dx_3) + L_2(\theta, x)(d\theta dx_2 + dx_3 dx_1) + L_3(\theta, x)(d\theta dx_3 + dx_1 dx_2) + Q,$$

where  $L_i(\theta, x) = \sum_j L_{ij}(\theta) x_j$  and Q consists of forms in  $d\theta$  and  $dx_i$ , whose coefficients are quadratic or higher in the  $x_i$ . In terms of matrices,  $\omega_t$  corresponds to

$$A = \begin{pmatrix} 0 & L_1 & L_2 & L_3 \\ -L_1 & 0 & L_3 & -L_2 \\ -L_2 & -L_3 & 0 & L_1 \\ -L_3 & L_2 & -L_1 & 0 \end{pmatrix} + \widetilde{Q},$$

where  $\widetilde{Q}$  has quadratic or higher terms in the  $x_i$  and the matrix is with respect to the basis  $\{d\theta, dx_1, dx_2, dx_3\}$ .  $i_{X_t}\omega_t = -\eta_t$  then becomes

$$(a_{\theta} \ a_1 \ a_2 \ a_3)A = -(\eta_{\theta} \ \eta_1 \ \eta_2 \ \eta_3)$$

with  $X_t = a_{\theta} d\theta + \sum_i a_i dx_i$ . Thus,  $(a_{\theta} \ a_1 \ a_2 \ a_3) = -$ 

$$\begin{array}{rcl} a_{\theta} \ a_{1} \ a_{2} \ a_{3} \end{array} &=& -(\eta_{\theta} \ \eta_{1} \ \eta_{2} \ \eta_{3}) A^{-1} \\ \\ &=& \frac{(\eta_{\theta} \ \eta_{1} \ \eta_{2} \ \eta_{3})}{L_{1}^{2} + L_{2}^{2} + L_{3}^{2}} \begin{pmatrix} 0 & L_{1} & L_{2} & L_{3} \\ -L_{1} & 0 & L_{3} & -L_{2} \\ -L_{2} & -L_{3} & 0 & L_{1} \\ -L_{3} & L_{2} & -L_{1} & 0 \end{pmatrix}$$

up to first order in x. This means that  $|X_t| < k|x|$  near C; hence, as  $x \to 0$ ,  $|\phi_1(\theta, x) - \phi_0(\theta, x)| \to 0$ , where  $\phi_t$  is the flow such that  $\frac{d\phi_t}{dt} = X_t$ . This concludes our proof.  $\Box$ 

## 4. Local Normal Forms

On a neighborhood  $N(C) = C \times D^3$  of C,  $\omega$  can be written as in Equation 3. If  $\omega$  is generic, then it is transverse to the zero section of  $\bigwedge_g^+$ , and  $(L_{ij}(\theta))$  is nondegenerate for all  $\theta$ .

**Lemma 4.**  $(L_{ij}(\theta))$  is symmetric and traceless.

*Proof.* By comparing zeroth order terms in the  $x_i$ ,  $d\omega = 0$  implies

$$\frac{\partial L_1}{\partial x_1} + \frac{\partial L_2}{\partial x_2} + \frac{\partial L_3}{\partial x_3} = 0,$$
  
$$\frac{\partial L_2}{\partial x_3} - \frac{\partial L_3}{\partial x_2} = 0, \frac{\partial L_3}{\partial x_1} - \frac{\partial L_1}{\partial x_3} = 0, \frac{\partial L_1}{\partial x_2} - \frac{\partial L_2}{\partial x_1} = 0.$$

The traceless symmetric matrix  $(L_{ij}(\theta))$  thus has a basis  $\{v_1(\theta), v_2(\theta), v_3(\theta)\}$  of eigenvectors for each  $\theta$  (though the  $v_i$  are not necessarily continuous in  $\theta$ ). Since  $(L_{ij}(\theta))$  is traceless, either two of the eigenvalues are positive and the remaining is negative for all  $\theta$ , or vice versa. Hence,  $(L_{ij}(\theta))$  gives rise to a splitting of  $\mathbf{R}^3 \times S^1 \to S^1$  into a real line bundle over  $S^1$  and a rank 2 vector bundle over  $S^1$ . Such splittings are classified by homotopy classes of maps from  $S^1$  to  $\mathbf{RP}^2$ , and  $\pi_1(\mathbf{RP}^2) = \mathbf{Z}/2\mathbf{Z}$ . Hence, we have the well-known:

**Proposition 5.** There exist two homotopy classes of splittings of  $\mathbb{R}^3 \times S^1 \to S^1$ , the oriented one and the unoriented one.

What is rather remarkable is the following:

**Theorem 6.** For either of the two splitting types, there exists a SD harmonic 2-form for a flat metric on  $S^1 \times D^3$  whose zero set is  $C = S^1 \times \{0\}$  and which has the given splitting type.

*Proof.* We give representatives of both types.

(A) Corresponding to the oriented splitting, we have:

$$\omega_A = x_1(d\theta dx_1 + dx_2 dx_3) + x_2(d\theta dx_2 + dx_3 dx_1) - 2x_3(d\theta dx_3 + dx_1 dx_2) = *_3\mu + d\theta \wedge \mu,$$

where  $\mu = d(\frac{1}{2}(x_1^2 + x_2^2) - x_3^2)$ , and  $*_3$  is the \*-operator for the flat metric on  $D^3$ . Here,  $(L_{ij}(\theta)) = \text{diag}(1, 1, -2)$ , with fixed positive and negative eigenspaces. Note that  $\omega_A$  is  $S^1$ -invariant.

(B) We construct  $\omega_B$ , corresponding to the unoriented splitting, as follows. Starting with

$$\Omega = x_1(d\theta dx_1 + dx_2 dx_3) + x_2(d\theta dx_2 + dx_3 dx_1) - 2x_3(d\theta dx_3 + dx_1 dx_2)$$
  
on  $[0, 2\pi] \times D^3$ , we glue  $\phi : \{2\pi\} \times D^3 \to \{0\} \times D^3$  by sending  
 $\theta \mapsto \theta - 2\pi x_1 \mapsto x_1 x_2 \mapsto -x_2 x_3 \mapsto -x_3.$ 

One easily verifies that  $\phi^*\Omega = \Omega$ .

**Theorem 7.** Given a generic SD harmonic 2-form  $\omega$  on M, there exists a 1-parameter family of closed 2-forms  $\omega_t$ ,  $t \in [0, 1]$ , on M which satisfy the following:

- (1)  $\omega_0 = \omega$ .
- (2) The  $\omega_t$  are symplectic away from their common zero set C.
- (3)  $\omega_t = \omega$  except on a neighborhood N(C) of C.
- (4) On each connected component of N(C),  $\omega_1$  is (up to sign) one of the two local forms  $\omega_A$  or  $\omega_B$  of Theorem 6.
- (5)  $[\omega_t] \in H^2(M; \mathbf{R})$  is independent of t.

*Proof.* Suppose for simplicity that the zero set C of the SD 2-form  $\omega$  is a single circle. Assume that  $\omega$  on N(C) gives rise to an oriented splitting, i.e., we are in case (A). (Case (B) is identical.) After an orthonormal change of frame, we may write

$$\omega = (L_{11}(\theta)x_1 + L_{12}(\theta)x_2)(d\theta dx_1 + dx_2 dx_3) + (L_{21}(\theta)x_1 + L_{22}(\theta)x_2)(d\theta dx_2 + dx_3 dx_1) + \lambda_3(\theta)x_3(d\theta dx_3 + dx_1 dx_2) + Q,$$

with, say,  $(L_{ij}(\theta))_{1 \le i,j \le 2}$  positive definite and  $\lambda_3(\theta) < 0$ . Here, the  $L_{ij}(\theta)$  and  $\lambda_3(\theta)$  are differentiable in  $\theta$ .

Now, take a 1-parameter family  $\beta_t = (1 - t)\omega + t\omega_A$  on N(C). After shrinking N(C) if necessary,  $\beta_t$  is symplectic on N(C) - C. Using a local version of our Moser argument (see the proof of Theorem 2), we see that there exist homeomorphisms

$$\phi_t : (N_0(C), \beta_0 = \omega) \xrightarrow{\sim} (N_t(C), \beta_t)$$

where  $N_t(C)$  are small neighborhoods of C whose boundary depend smoothly on t,  $\phi_t = id$ on C, and  $\phi_t$  is smooth away from C. The  $\phi_t$  allow us to remove  $(N_0(C), \omega)$  and graft on  $(N_t(C), \beta_t)$ . Hence there exists a global family  $\omega_t$  on M with  $\omega_0 = \omega$  and  $\omega_1|_{N(C)} = \omega_A$ ,

after further shrinking N(C). Moreover, the perturbation can be performed in an arbitrarily small neighborhood of C and without altering the cohomology class.

In essence, Theorem 7 tells us that, in studying the singular circles of  $\omega$  from a global perspective, we may assume that the zeros are either (A) or (B).

## 5. Contact structures on the boundaries

In this section we investigate the boundary properties of  $\omega_A$  and  $\omega_B$ . More precisely, we have:

**Theorem 8.** There exist contact forms  $\lambda_A$  and  $\lambda_B$  on  $\partial N(C) = S^1 \times S^2$  such that  $d\lambda_A = i^* \omega_A$ and  $d\lambda_B = i^* \omega_B$ , where *i* is the inclusion  $S^1 \times S^2 \hookrightarrow N(C) = S^1 \times D^3$ .

*Proof.* (A) Consider the following  $S^1$ -invariant 1-form

$$\lambda = -\frac{1}{2}(x_1^2 + x_2^2 - 2x_3^2)d\theta + x_2x_3dx_1 - x_1x_3dx_2$$

on N(C). We then compute that  $d\lambda = \omega_A$  on N(C) and

$$\sum_{i} x_{i} dx_{i} \wedge \lambda \wedge d\lambda = \left(\frac{1}{2}(x_{1}^{2} + x_{2}^{2})(x_{1}^{2} + x_{2}^{2} + 2x_{3}^{2}) + 2x_{3}^{4}\right) d\theta dx_{1} dx_{2} dx_{3}$$

Since  $S^1 \times S^2 = \{\sum_i x_i^2 = 1\}$  is an integral submanifold of  $\sum_i x_i dx_i = 0$ ,  $i^*(\lambda \wedge d\lambda) \neq 0$ on  $S^1 \times S^2$  if and only if  $\lambda \wedge d\lambda \wedge \sum_i x_i dx_i \neq 0$  in a neighborhood of  $S^1 \times S^2$ . Noting that  $\lambda \wedge d\lambda \wedge \sum_i x_i dx_i = 0$  if and only if  $x_1 = x_2 = x_3 = 0$ , we conclude that  $\lambda_A = i^*\lambda$  is a contact 1-form on  $S^1 \times S^2$  and  $(M - N(S^1), \omega_A = d\lambda)$  is a symplectic manifold with contact-type boundary.

(B) Consider the 1-form

$$\lambda = -\frac{1}{2}(x_1^2 + x_2^2 - 2x_3^2)d\theta + x_2x_3dx_1 - x_1x_3dx_2$$

on  $[0, 2\pi] \times D^3$ . Using the notation from the proof of Theorem 6,  $d\lambda = \Omega$  and  $\phi^*\lambda = \lambda$ , so we glue together (after pulling back via *i*) a contact 1-form  $\lambda_B$  such that  $d\lambda_B = i^*\omega_B$ . The rest is the same as (A).

Let us now describe the orbits of the Reeb vector fields.

(A)  $\omega_A$  is compatible with a metric  $\tilde{g} = cg$ , where g is the standard product metric on  $S^1 \times D^3$ . We can then write the compatible J satisfying  $\tilde{g}(x, y) = \omega(x, Jy)$  as  $J = -\frac{1}{c}A$ , where

$$A = \begin{pmatrix} 0 & x_1 & x_2 & -2x_3 \\ -x_1 & 0 & -2x_3 & -x_2 \\ -x_2 & 2x_3 & 0 & x_1 \\ 2x_3 & x_2 & -x_1 & 0 \end{pmatrix}$$

represents  $\omega$  with respect to  $\{\theta, x_1, x_2, x_3\}$ . Now, the Reeb vector field X for  $\lambda_A$  is given, up to multiple, by

$$J\left(\sum_{i} x_i \frac{\partial}{\partial x_i}\right) = \frac{-1}{\sqrt{x_1^2 + x_2^2 + 4x_3^2}} \left( (x_1^2 + x_2^2 - 2x_3^2) \frac{\partial}{\partial \theta} - 3x_2 x_3 \frac{\partial}{\partial x_1} + 3x_1 x_3 \frac{\partial}{\partial x_2} \right)$$

Finally,  $\lambda_A(X) = 1$  implies that

$$X = \frac{1}{f} \left[ (x_1^2 + x_2^2 - 2x_3^2) \frac{\partial}{\partial \theta} - 3x_2 x_3 \frac{\partial}{\partial x_1} + 3x_1 x_3 \frac{\partial}{\partial x_2} \right],$$

with

$$f = -\frac{1}{2} \left[ (x_1^2 + x_2^2)(x_1^2 + x_2^2 + 2x_3^2) + 4x_3^4 \right].$$

Solving for the orbits,  $x_1^2 + x_2^2$  and  $x_3^2$  are fixed for each orbit, and hence,

$$x_1 = \sqrt{1 - r^2} \cos R_1(r) t$$
  

$$x_2 = \sqrt{1 - r^2} \sin R_1(r) t$$
  

$$x_3 = r$$
  

$$\theta = R_2(r) t + c,$$

where r is a constant, and  $R_1$  and  $R_2$  are functions of r.

In particular, the noteworthy closed orbits are  $S^1 \times \{(0,0,1)\}, S^1 \times \{(0,0,-1)\}$ , and  $S^1 \times \{(x_1, x_2, 0)\}$ , with  $x_1^2 + x_2^2 = 1$  and  $x_1, x_2$  fixed. These correspond to the stable and unstable gradient directions in the Morse theory of  $\frac{1}{2}(x_1^2 + x_2^2 - 2x_3^2)$  near (0,0,0). Moreover, the orbit  $S^1 \times \{(0,0,1)\}$  is nondegenerate, and so is the family  $S^1 \times \{(x_1, x_2, 0)\}$ . There are other closed orbits, but these do not seem to have any Morse-theoretic significance. The relevance of such periodic orbits is manifest in Taubes' paper [6].

(B) We apply the previous considerations and work on  $[0, 2\pi] \times S^2 / \sim$ . There is one orbit  $S^1 \times \{(0, 0, \pm 1)\}$ , which is a double of the orbits for (A). Since  $\phi$  identifies  $(2\pi, (x_1, x_2, 0)) \sim (0, (x_1, -x_2, 0))$ , we also have the doubled closed orbits  $S^1 \times \{(x_1, \pm x_2, 0)\}$ , with  $x_2 \neq 0$ , and the single closed orbits  $S^1 \times \{(1, 0, 0)\}, S^1 \times \{(-1, 0, 0)\}$ .

**Remark.** There is an example of a singularity of type (B) bounding a disk, which can be made to vanish.

Our final point of investigation is to determine whether the two contact forms represent the same *contact structure* on  $S^1 \times S^2$ . Let  $\xi_A$ ,  $\xi_B$  be the contact structures (i.e., 2-plane distributions) corresponding to the 1-forms  $\lambda_A$ ,  $\lambda_B$ . Then we remark the following

**Proposition 9.**  $\xi_A$  and  $\xi_B$  are both overtwisted.

*Proof.* A perturbation of the upper hemisphere  $\{x_3 \ge 0\} \cap (\{\text{pt.}\} \times S^2) \subset S^1 \times S^2$  is an overtwisted disk.

Since overtwisted contact structures on a 3-manifold N are classified by the homotopy type of the contact 2-plane distribution on TN (c.f. [1]), it suffices to determine whether  $\xi_A$ and  $\xi_B$  are homotopic as 2-plane distributions inside  $T(S^1 \times S^2)$ . We have the following: **Theorem 10.** The contact structures  $\xi_A$  and  $\xi_B$  are overtwisted contact structures for distinct homotopy classes of 2-plane fields on  $S^1 \times S^2$ .

*Proof.* First let us describe the homotopy classes of 2-plane fields of *degree* 1 on  $S^1 \times S^2$ . A trivialization of  $T(S^1 \times S^2)$  and a choice of nonsingular transverse vector field (in our case the Reeb vector field) gives rise to a Gauß map  $\psi : S^1 \times S^2 \to S^2$ , and the homotopy classes of 2-plane fields are in 1-1 correspondence with homotopy classes  $[S^1 \times S^2, S^2]$ . We are interested in maps where the induced map  $H_2(S^1 \times S^2; \mathbb{Z}) \to H_2(S^2; \mathbb{Z})$  has degree one. The set  $[S^1 \times S^2, S^1]_1$  of homotopy classes of maps of degree one is computed as follows:

$$S^{1} \times S^{2}, S^{2}]_{1} \simeq [S^{1}, \operatorname{Map}_{1}(S^{2}, S^{2})]$$
  

$$\simeq [S^{1}, \operatorname{Diff}^{+}(S^{2})]$$
  

$$\simeq [S^{1}, SO(3)]$$
  

$$\simeq \pi_{1}(SO(3)) \simeq \mathbb{Z}/2\mathbb{Z}.$$

Here,  $Map_1(S^2, S^2)$  is the set of maps of degree 1, and  $Diff^+(S^2)$  is the space of orientationpreserving diffeomorphisms of  $S^2$ . The representatives of the two classes are given by

(4) 
$$\Psi_n : (\theta, x) \mapsto R_{n\theta}(x),$$

with n = 0, 1, where  $R_{\alpha}$  is rotation by  $\alpha$  about the  $x_1$ -axis. (Note we could have chosen representatives to rotate about any axis – however, it is important to remember that  $R_{\alpha}$  is a rotation about the  $x_1$ -axis.)

Next trivialize  $T(S^1 \times S^2) \simeq \mathbf{R}^3 \times (S^1 \times S^2)$  as follows: First view  $S^2$  as the unit sphere  $\sum_i x_i^2 = 1$  inside  $\mathbf{R}^3$ . At the point  $(\theta, x_1, x_2, x_3) \in S^1 \times S^2$ , the tangent vector  $\frac{\partial}{\partial \theta}$  is mapped to the unit vector  $(x_1, x_2, x_3) \in \mathbf{R}^3$  and  $v \in T_{(x_1, x_2, x_3)}S^2$  to the corresponding tangent vector viewed inside  $\mathbf{R}^3$ .

Let us consider Case (A). After rescaling, the Reeb vector field is

$$X(\theta, x_1, x_2, x_3) = (-x_1^2 - x_2^2 + 2x_3^2, 3x_2x_3, -3x_1x_3, 0).$$

Let  $\psi: S^1 \times S^2 \to S^2$  be the corresponding Gauß map. Since  $\psi(\theta, x)$  is independent of t, we write  $u(x) = \psi(\theta, x)$ . Then u maps  $(0, 0, \pm 1) \mapsto (0, 0, \pm 1), (x_1, x_2, 0) \mapsto (-x_1, -x_2, 0)$ , and sends  $x_3 = const$  to  $x_3 = \tau(const)$ . Here  $\tau: [-1, 1] \to [-1, 1]$  sends -1, 0, 1 to itself and satisfies  $\tau(-x_3) = -\tau(x_3)$ . Since  $\tau$  is homotopic to the identity via a homotopy  $\tau_t$  which fixes -1, 0, 1 and satisfies  $\tau_t(-x_3) = -\tau_t(x_3), \psi$  is homotopic to:

$$(\theta, x) \mapsto S_{\pi + \sigma(x_3)}(x),$$

where  $S_{\alpha}$  is the rotation by  $\alpha$  along the  $x_3$ -axis, and  $\sigma(-x_3) = -\sigma(x_3)$ . We can then homotop  $\sigma = \sigma_0$  to  $\sigma_1(x_3) = 0$ , via  $\sigma_t$  which satisfies  $\sigma_t(-x_3) = -\sigma_t(x_3)$ . Therefore,  $\psi$  is homotopic to  $(\theta, x) \mapsto S_{\pi}(x)$ , which in turn is homotopic to  $\Psi_0$  in Equation 4.

Let us now consider Case (B). Instead of gluing  $\{2\pi\} \times S^2$  and  $\{0\} \times S^2$  via a twist, we use the diffeomorphism

$$\Phi : [0, 2\pi] \times S^2 \xrightarrow{\sim} [0, 2\pi] \times S^2,$$
  
$$(\theta, x) \mapsto (\theta, R_{\frac{\theta}{2}} x).$$

Hence, if we push forward via  $\Phi$ , we may glue  $\{2\pi\} \times S^2 \xrightarrow{\sim} \{0\} \times S^2$  identically. Our Reeb field for (B) is  $\Phi_*X$ , where  $X = (-x_1^2 - x_2^2 + 2x_3^2, 3x_2x_3, -3x_1x_3, 0)$  is defined over  $[0, 2\pi] \times S^2$ .

The corresponding Gauß map  $\psi: S^1 \times S^2 \to S^2$  is (homotopic to one) given on  $[0, 2\pi] \times S^2$  by

$$(\theta, x) \mapsto R_{\theta/2} u R_{\theta/2}^{-1}(x),$$

where  $u: S^2 \to S^2$  was described in the previous paragraph. Observe that  $R_{\pi}uR_{\pi}^{-1}(x) = u(x)$ , which allows us to glue  $\theta = 0$  and  $\theta = \pi$ . The homotopy  $u_t$  from u to  $S_{\pi}$  which was applied in the analysis of Case (A) also can be applied here, by observing that  $R_{\pi}u_tR_{\pi}^{-1} = u_t$ . Now,

$$R_{\theta/2}S_{\pi}R_{\theta/2}^{-1}(x) = R_{\theta/2}R_{\theta/2}S_{\pi}(x) = R_{\theta}S_{\pi}(x),$$

and hence  $\psi$  is homotopic to  $\Psi_1$ .

This shows that  $\xi_A$  and  $\xi_B$  are not homotopic as 2-plane fields.

We close with the following question:

**Question.** Can the contact homology theory developed by Hofer and Eliashberg (among others), or variations of it, shed any light on the singular symplectic forms?

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DEPARTMENT OF MATHEMATICS, UNIVERSITY OF SOUTHERN CALIFORNIA, LOS ANGELES, CA 90089-1113

*E-mail address*: khonda@math.usc.edu *URL*: http://math.usc.edu/~khonda