Magic Angles for Twisted Bilayer Graphene

# Microlocal Analysis and Mathematical Physics In honor of Anders Melin's 80th birthday 

Maciej Zworski

September 21, 2023




A project in the time of covid-19
2020: Simon Becker, Mark Embree, Jens Wittsten, MZ: BEWZ

2022: Simon Becker, Tristan Humbert, MZ: BHZ
2023: Michael Hitrik, MZ: HZ; Simon Becker, MZ: BZ


Motivation: bilayer graphene
graphite

graphene


## Motivation: bilayer graphene

## graphite



MacGyver in the physics (ab


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Geim-Novoselov '04

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Cao et al '18, Yankovitz et al '18: superconductivity at $\theta \simeq 1.08^{\circ}$

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Cao et al '18, Yankovitz et al '18: superconductivity at $\theta \simeq 1.08^{\circ}$ Predicted by Bistritzer-MacDonald '11

The chiral model of TBG

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PHYSICAL REVIEW LETTERS 122, 106405 (2019)

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Grigory Tarnopolsky, Alex Jura Kruchkov, ${ }^{*}$ and Ashvin Vishwanath Department of Physics, Harvard University, Cambridge, Massachusetts 02138, USA

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z=x_{1}+i x_{2}, \quad D_{\bar{z}}:=\frac{1}{2 i}\left(\partial_{x_{1}}+i \partial_{x_{2}}\right) \\
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Seeley 85: $P(\alpha)=e^{i x} D_{x}+\alpha e^{i x}, x \in \mathbb{S}^{1}, \operatorname{Spec}(P(\alpha))=\mathbb{C}, \alpha \in \mathbb{Z}$.

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Editors' Suggestion

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Theorem (BEWZ '20) There exists a discrete set $\mathcal{A} \subset \mathbb{C}$ such that

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\operatorname{Spec}_{L_{0}^{2}(\mathbb{C} / \Lambda)} D(\alpha)= \begin{cases}\Lambda^{*}+\{K,-K\} & \alpha \notin \mathcal{A} \\ \mathbb{C} & \alpha \in \mathcal{A}\end{cases}
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| 1 | 0.58566355838955 |  |
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Magic angles vs. Scattering resonances


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Magic $\alpha$ 's


Resonances for $B_{\mathbb{H}^{2}}(0,1)$

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(Known for obstacles in hyperbolic plane: Vodev, Borthwick...)

## Flat bands

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Theorem (BHZ '22; implicit in BEWZ '20)

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\exists \mathrm{k} \notin \Lambda^{*}+\{K,-K\} \quad E_{1}(\alpha, \mathrm{k})=0 \Longrightarrow \forall \mathrm{k} E_{1}(\alpha, \mathrm{k})=0 .
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Theorem (BHZ '22) For all $p>1$

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\sum_{\alpha \in \mathcal{A}} \alpha^{-2 p} \in \frac{\pi}{\sqrt{3}} \mathbb{Q}
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Remark: Luskin-Watson '21 showed $|\mathcal{A} \cap(0.583,0.589)| \geq 1$

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Example $Q(\alpha, k)=D(\alpha)+k, m\left(\alpha_{0}, k\right)=\mathbb{1}_{\{K,-K\}+\Lambda^{*}}$ (symmetry protected states!)

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Example $Q(\alpha, k)=\left(2 D_{\bar{z}}+k\right)^{2}-\alpha^{2} U(z) U(-z)$ : a scalar model in which $m(0, k)=2 \mathbb{1}_{\Lambda^{*}}(k)>\operatorname{dim} \operatorname{ker} Q(0, k)$.

Works for general potentials with $\mathbb{Z}_{3}^{2} \rtimes \mathbb{Z}_{3}$ symmetries

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U_{\theta}(z):=\sum_{k=0}^{2} \omega^{k}\left(\cos ^{2} \theta e^{\frac{1}{2}\left(\bar{z} \omega^{k}-z \bar{\omega}^{k}\right)}+\sin ^{2} \theta e^{\bar{z} \omega^{k}-z \bar{\omega}^{k}}\right)
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Structure of eigenfunctions at the flat band

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Theorem ( BHZ '22) If $\alpha \in \mathcal{A}$ is simple then the unique zero has to appear at the stacking point $z_{S}:=-z(K)=\sqrt{3} / i$.

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Dirac point at $k \Longleftrightarrow k \in \operatorname{Spec}_{L_{0}^{2}(\mathbb{C} / \Gamma)} D_{B}(\alpha)$
Theorem ( BZ '23) If $\underline{\alpha} \in \mathcal{A}$ is simple ( + one more condition) and $0<B_{0} \ll 1$ then there are no flat bands and for $\alpha \sim \underline{\alpha}$ Dirac points (eigenvalues of $D_{B}(\alpha)$ ) are close to the $\Gamma$ point.

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## Protected states



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\sigma=d \xi_{1} \wedge d x_{1}+d \xi_{2} \wedge d x_{2}=2 \operatorname{Re} d \zeta \wedge d z, \quad z=x_{1}+i x_{2}, \quad \zeta=\frac{1}{2}\left(\xi_{1}-i \xi_{2}\right)
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Passage to PDE with analytic coefficients has its complications

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## Application to the chiral model of TBG

Exponential decay of solutions near $x_{0}$ guaranteed by

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\begin{gathered}
q\left(x_{0}, \xi\right)=0 \Longrightarrow\{q, \bar{q}\}=0, \quad\{q,\{q, \bar{q}\}\} \neq 0 \\
q=(2 \bar{\zeta})^{2}-U(z) U(-z), \quad q=0 \Leftrightarrow 2 \bar{\zeta}= \pm \sqrt{U(z) U(-z)} \\
\{q, \bar{q}\}= \pm 8 i \operatorname{Im}\left((\overline{U(z) U(-z)})^{\frac{1}{2}} \partial_{z}(U(z) U(-z))\right)
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Precalculus $\Longrightarrow\{q, \bar{q}\}=0$ on $\pi^{-1}$ (hexagon) $\cap q^{-1}(0)$
hexagon $=$ spanned by $\pm z_{S}+\Lambda, z_{S}=i / \sqrt{3}, \omega z_{S} \equiv z_{S} \bmod \Lambda$

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$\{q,\{q, \bar{q}\}\}=\frac{128}{9} \pi^{2}(c-1)^{2}(2 c+1)(2 c-9), \quad c:=\cos (2 \pi \sqrt{3} t / 3)$

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Lower order terms matter as we see in the figure above!


Anonymous '23 (communicated by Simon Becker)
Proofs are for dinosaurs - why don't you just put it on a computer.


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