

Localization in Classical Statistical Mechanics

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BIRS

Talk Outline

1. Model of Interest and Basic Question
2. Some Background
3. Kac Interactions and Our Results
4. Proof Outline
5. Research In Progress

The Classical XY model

- ▶ Let $\mathcal{G} = (\mathbb{Z}^d, \mathcal{E})$ with distance $\|x - y\|_2$ ($d \geq 2$).
- ▶ For $x \in \mathbb{Z}^d$, $\sigma_x \in \mathbb{S}^1$, $\sigma = (\sigma_x)_{x \in \mathbb{Z}^d} \in \Omega$
Equipped with i.i.d. uniform measure $\prod_{x \in \mathbb{Z}^d} d\nu(\sigma_x)$.
Define $-\mathcal{H}_L : \Omega \rightarrow \mathbb{R}$ by

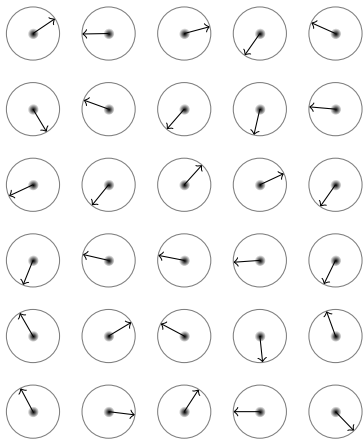
$$-\mathcal{H}_L(\sigma) = \sum_{\|x-y\|_2=1, x,y \in \Lambda_L} \sigma_x \cdot \sigma_y$$

- ▶ Gibbs Weight

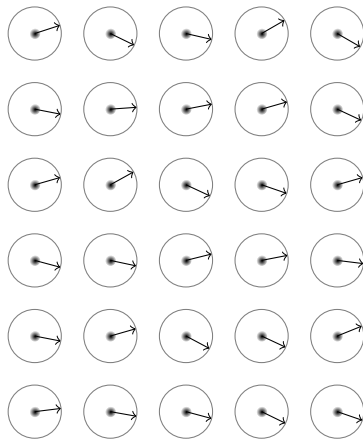
$$d\mu_L(\sigma) = Z^{-1} e^{-\beta \mathcal{H}_L(\sigma)} \prod_{x \in \Lambda_L} d\nu(\sigma_x)$$

- ▶ Averages denoted by

$$\langle \cdot \rangle.$$



High Temperature, β small



Low Temperature, β large

Ordering in Gibbs States

- ▶ Given $A \subset \mathbb{Z}^d$ finite, how does

$$m_L(A) = \left\langle \frac{1}{|A|} \sum_{x \in A} \sigma_x \right\rangle_L$$

behave as A varies. Are there preferred directions?

- ▶ Technically, need **Boundary Conditions**:

$$\begin{aligned} -\mathcal{H}(\sigma|\eta) = & \sum_{\|x-y\|_2=1, x,y \in \Lambda_L} \sigma_x \cdot \sigma_y \\ & + \sum_{\|x-y\|_2=1, x \in \Lambda_L, y \notin \Lambda_L} \sigma_x \cdot \eta_y \quad (1) \end{aligned}$$

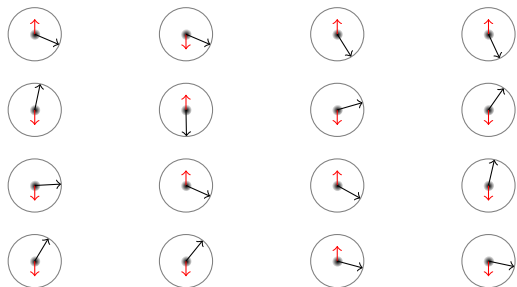
Natural choices: $\eta_x \equiv u$, $u \in \mathbb{S}^1$.

Local Fields

For $x \in \mathbb{Z}^d$, $\alpha_x(\omega)$ i.i.d. $\{\pm 1\}$ Bernoulli.

$$-\mathcal{H}^\omega(\sigma|\eta) = \sum_{\|x-y\|_2=1, x,y \in \Lambda_L} \sigma_x \cdot \sigma_y + \epsilon \sum_{x \in \Lambda_L} \alpha_x(\omega) e_2 \cdot \sigma_x + \sum_{\|x-y\|_2=1, x \in \Lambda_L, y \notin \Lambda_L} \sigma_x \cdot \eta_y \quad (2)$$

Gibbs State as before.



Basic Question

- ▶ Given ϵ , does this model order at β large (T low)?
- ▶ Given $A \subset \mathbb{Z}^d$ finite, how does

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- ▶ Believed that ordering occurs in \hat{e}_1 direction.
- ▶ Robust: Any symmetric distribution for α_x . Even may take bias, e.g. $\mathbb{P}(\alpha_x = 1) = p$.
- ▶ Similar Phenomena should exist in any spin model with continuous spin space.

Why is this Interesting? Two Related Models:

- ▶ From now on, we restrict attention to \mathbb{Z}^d , $d = 2$.

Model 1: As above with $\epsilon = 0$:

$$-\mathcal{H}(\sigma|\eta) = \sum_{\|x-y\|_2=1, x,y \in \Lambda_L} \sigma_x \cdot \sigma_y + \sum_{\|x-y\|_2=1, x \in \Lambda_L, y \notin \Lambda_L} \sigma_x \cdot \eta_y \quad (3)$$

- ▶ **Fact:** No matter the η you take,

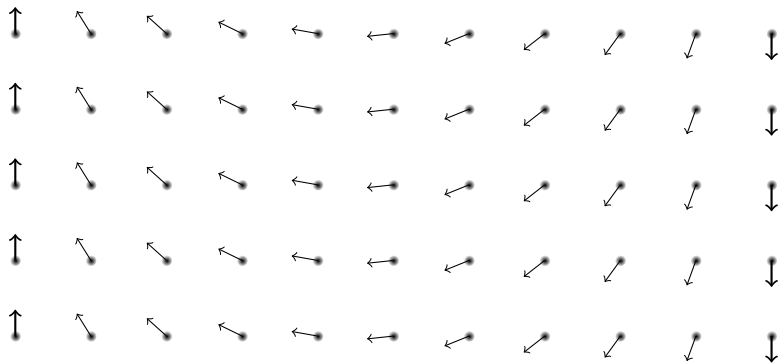
$$m_L(A) \rightarrow 0 \text{ as } L \rightarrow \infty$$

ANY $\beta > 0$.



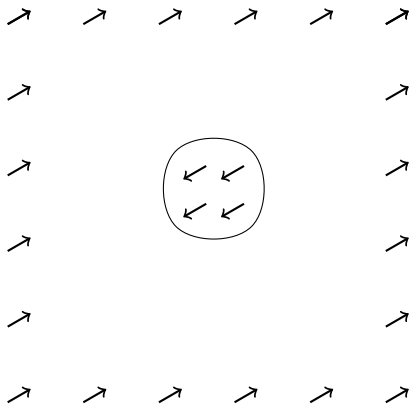
In $L \times L$ slab

$$-\inf_{\sigma} \mathcal{H}(\sigma | \uparrow_L \downarrow_R) + \inf_{\sigma} \mathcal{H}(\sigma | \uparrow_L \uparrow_R) \sim -L^2 \times (1 - \cos(\frac{\pi}{L})) \sim -\frac{\pi^2}{2}.$$



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Cost: $\frac{1}{\log L}$ for $d = 2$

L^{d-2} for $d \geq 3$

- ▶ Rigorous Statement known as Mermin-Wagner Theorem ['66], Dobrushin-Shlosman ['75] etc.

Theorem

In $d = 2$ for any short range spin system with continuous symmetry ALL Gibbs states are invariant.

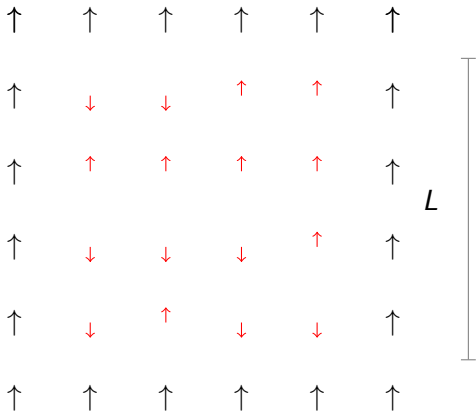
- **Model 2:** $\sigma_x \in \{\pm\hat{e}_2\}$ and $\nu(\sigma_x) = \frac{1}{2}\delta_{e_2}(\sigma_x) + \frac{1}{2}\delta_{-\hat{e}_2}(\sigma_x)$,
 ϵ field strength.

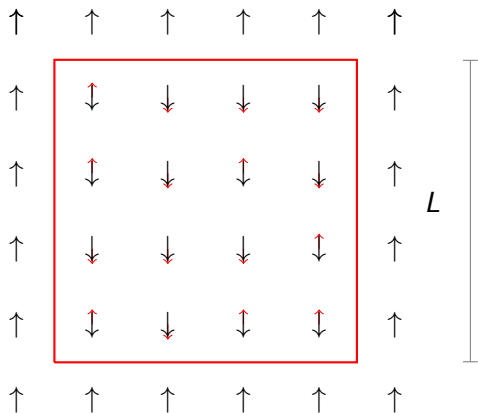
$$\begin{aligned}
 -\mathcal{H}^\omega(\sigma|\eta) = & \sum_{\|x-y\|_2=1, x,y \in \Lambda_L} \sigma_x \cdot \sigma_y + \epsilon \sum_{x \in \Lambda_L} \alpha_x(\omega) e_2 \cdot \sigma_x \\
 & + \sum_{\|x-y\|_2=1, x \in \Lambda_L, y \notin \Lambda_L} \sigma_x \cdot \eta_y \quad (4)
 \end{aligned}$$

Theorem (Imry-Ma '75 Aizenman-Wehr '89)

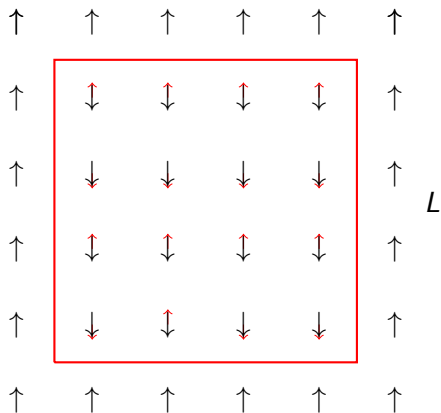
ANY $\beta > 0$, for all η and $\epsilon > 0$,

$$\mathbb{E}_\alpha [\langle \sigma_x \rangle_L] \rightarrow 0$$





Contour Energy: $\sim L^{d-1}$



L

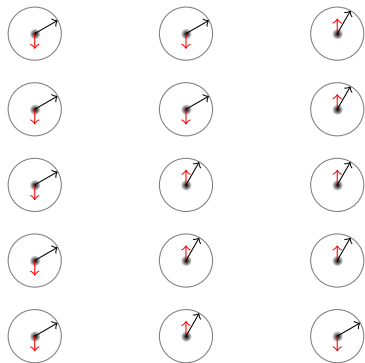
Contour Energy: $\sim L^{d-1}$

Local Field Energy Difference:

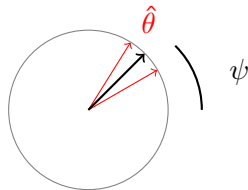
$$2\epsilon \sum_{x \in \Lambda_L} \alpha_x \sim \epsilon \mathcal{N}_L L^{\frac{d}{2}}$$

Back to our Question

► Intuitive picture at β large:

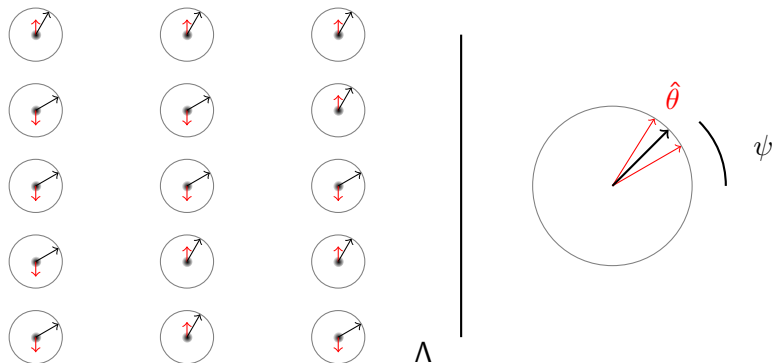


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Back to our Question

► Intuitive picture at β large:



Fix ψ ; write $\theta_x = \psi + \hat{\theta}_x$:

$$\begin{aligned} & -\mathcal{H}_\Lambda(\theta) \\ \asymp & -\frac{1}{2} \sum_{\langle x,y \rangle \in \Lambda} [\hat{\theta}_x - \hat{\theta}_y]^2 + \epsilon \sum_{x \in \Lambda} \frac{1}{\deg_\Lambda(x)} [\alpha_x - \alpha_\Lambda] [\cos(\psi) \hat{\theta}_x] + \epsilon |\Lambda| \alpha_\Lambda \sin(\psi) \end{aligned} \quad (5)$$

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- ▶ Optimizing over $\hat{\theta}_x$ gives

$$\epsilon^2 \cos^2(\psi) \sum_{\langle x,y \rangle} [g_x - g_y]^2 |\Lambda| + \alpha_\Lambda \sin(\psi)$$

where

$$-\Delta_N \cdot g_x = \alpha_x - \alpha_\Lambda$$

Our Contribution: Kac Interactions

Let $J : \mathbb{R}^d \rightarrow [0, \infty)$, $\int J = 1$ be smooth with compact support and $J_K(u) = K^{-d} J(\frac{u}{K})$.

$$\begin{aligned} -\mathcal{H}^\omega(\sigma|\eta) = & \frac{1}{2} \sum_{x,y \in \Lambda_L} J_K(x-y) \sigma_x \cdot \sigma_y + \epsilon \sum_{x \in \Lambda_L} \alpha_x(\omega) \mathbf{e}_2 \cdot \sigma_x \\ & + \frac{1}{2} \sum_{x \in \Lambda_L, y \notin \Lambda_L} J_K(x-y) \sigma_x \cdot \eta_y \quad (6) \end{aligned}$$

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Block Averages:

$$M_z = K^{-d(1+\delta)} \sum_{x \in B_{K^{1+\delta}}} \sigma_x.$$

Theorem (Main Theorem)

Let $d \geq 2$ be fixed. For ϵ, ξ small and $\beta > \beta(\epsilon, \xi)$, $K \geq K(\epsilon)$,

$$\|\langle M_z \rangle_L^{\omega, \rightarrow} - \rho(\beta, \epsilon) \mathbf{e}_1\|_2 \leq \xi$$

as $L \rightarrow \infty$ except for $z \in \mathcal{D}^\omega \subset \mathbb{Z}^d - \omega$ a.s.

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$$\rho(\beta, \epsilon) \rightarrow \sqrt{1 - \epsilon^2} \text{ as } \beta \rightarrow \infty.$$

Proof Outline

- ▶ Mean Field Analysis $\rightsquigarrow -\phi$.
- ▶ Coarse Graining and Large Deviation Estimates for Magnetization Profiles.
- ▶ Randomness and Contour Identification.
- ▶ Peierls Contour Counting Estimates.

Mean Field Theory

- ▶ For $\delta > 0$, let $k_{<} = K^{1-\delta}$. All $x \in B_{k_{<}}$ feel similar interaction.

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- ▶ Typically $\frac{|B_{k_{<}^\pm}|}{|B_{k_{<}}|} = \frac{1}{2} + O(|B_{k_{<}}|^{-\frac{1}{2}})$

- ▶ Random variables of interest:

$$M^\pm = \frac{1}{|B_{k<}^\pm|} \sum_{x \in B_{k<}^\pm} \sigma_x$$

- ▶ Large Deviation Principle for (M^+, M^-)

$$\langle \mathbf{1}_{\{M^\pm \sim m^\pm\}} \rangle \sim e^{-\beta |B_{k<}| (\phi(m^+, m^-) - \inf \phi)}.$$

$$\begin{aligned} \phi(m^+, m^-) = & -\frac{1}{8} \|m^+ + m^-\|_2^2 - \frac{\epsilon}{2} \hat{e}_2 \cdot (m^+ - m^-) \\ & - \frac{1}{2\beta} (S(m^+) + S(m^-)) \quad (7) \end{aligned}$$

and

$$S(m) = \inf_{h \in \mathbb{R}^2} \left(\log \int_{\mathbb{S}^1} d\nu(\sigma) e^{\sigma \cdot h} - m \cdot h \right)$$

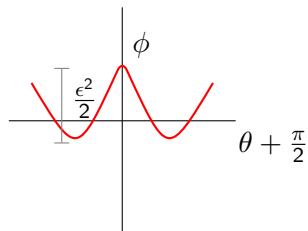
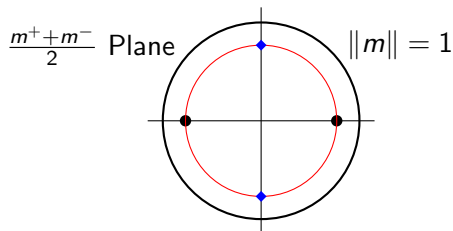
The Limits of this Method

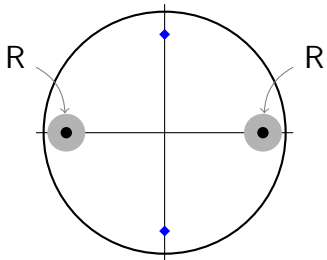
Want to bring LDP

$$\langle \mathbf{1}_{\{M^\pm \sim m^\pm\}} \rangle \sim e^{-\beta |B_{k<}| (\phi(m^+, m^-) - \inf \phi)}.$$

to \mathbb{Z}^d .

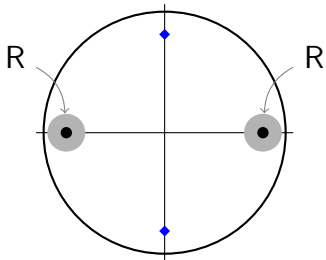
Optimizers of $\phi(m^+, m^-)$ characterized by $\frac{m^+ + m^-}{2}$.





Best we can hope for:

$$\langle \mathbf{1}_{\{M^\pm \text{ not in } R\}} \rangle \sim e^{-c\beta|B_{k < \epsilon^2}|}.$$

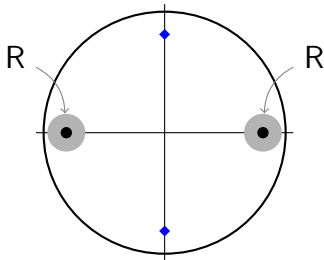


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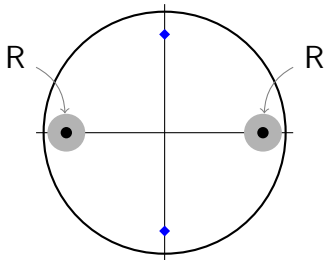
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Leads to correction:

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To suppress this need

$$|B_{k<}| \gg \epsilon^{-2}$$