On Ternary Quadratic Forms

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Dedicated to the memory of Arnold E. Ross

1 Introduction.

Let $q(x) = q(x_1, x_2, x_3)$ be a positive definite ternary quadratic form with integral coefficients. In 1946 Ross and Pall [RP] conjectured that every sufficiently large square-free integer that is represented by q modulo N for all N is in fact integrally represented. This conjecture was proven in 1988 [Du1] as an application of bounds for sums of Kloosterman sums given by Iwaniec [Iw] (see also [DS-P] and [Du2] for an exposition and references.)

In this paper I will quantify "sufficiently large" in this result in terms of the determinant D of q, which is defined by the 3×3 determinant

$$D = \det(\partial^2 q / \partial x_i \partial x_j)$$

and is known to be a positive even integer.

Theorem 1 There is an absolute constant c > 0 so that q(x) = n has an integral solution provided that

$$n > c D^{337}$$

is square-free and that the congruence $q(x) \equiv n \pmod{8D^3}$ has a solution. The constant c is ineffective.

2 A uniform asymptotic formula.

Theorem 1 follows from a uniform asymptotic formula for the number of representations

$$r_q(n) = \#\{x \in \mathbb{Z}^3 ; q(x) = n\}$$

This quantity does not change if q is replaced by a \mathbb{Z} -equivalent form, that is one obtained from q by an invertible integral change of variables. The genus

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G of q consists of those forms that are equivalent to q under invertible variable changes over the p-adic integers, for all p. The genus G is known to split into finitely many \mathbb{Z} -equivalence classes. The average number of representations of n by forms in G is

$$r(n,G) = M(G)^{-1} \sum_{\{q\} \in G} \omega_q^{-1} r_q(n)$$
(1)

the sum being over $\mathbbm{Z}\text{-inequivalent}\;q$ in G. Here ω_q is the number of automorphs of q and

$$M(G) = \sum_{\{q\} \in G} \omega_q^{-1} \tag{2}$$

is the mass (or weight) of the genus G. We shall establish the following uniform asymptotic formula.

Theorem 2 Fix $\epsilon > 0$. Then for n square-free we have

$$r_q(n) = r(n,G) + O(D^{\frac{11}{2}}n^{\frac{1}{2} - \frac{1}{28}}(Dn)^{\epsilon})$$

where the implied constant depends only on ϵ .

It follows from the fundamental work of Siegel [Si] that for n square-free we have the formula

$$r(n,G) = \kappa L(1,\chi)\mathcal{N}(n;q)\sqrt{n/D},$$
(3)

where $\chi(\cdot) = \left(\frac{-2Dn^*}{\cdot}\right)$ is the Kronecker symbol with n^* the discriminant of $\mathbb{Q}(\sqrt{n})$ and $L(s,\chi)$ is the Dirichlet *L*-function. Also

$$\mathcal{N}(n;q) = \#\{x \in (\mathbb{Z}/m\mathbb{Z})^3; q(x) \equiv n \pmod{m}\} m^{-2}$$
(4)

for $m = (2D)^3$ and $\kappa > 0$ is a constant. Theorem 1 is a consequence of the Theorem 2 together with the following lower bound, that is valid for all $\epsilon > 0$ under the condition that the congruence $q(x) \equiv n \pmod{8D^3}$ has a solution:

$$r(n,G) \gg_{\epsilon} n^{\frac{1}{2}} D^{-\frac{13}{2}} (nD)^{-\epsilon}.$$
 (5)

This follows from (3) and Siegel's lower bound for $L(1, \chi)$ and is where the ineffectivity of c in Theorem 1 arises.

As in [DS-P] one may refine Theorem 2 by replacing r(n, G) by the associated quantity for the spinor genus of q. This allows one to weaken the condition on n to be that the square part of n be prime to the level of q (or has a limited ged with the level). Further refinements as in [DS-P] should be possible as well.

3 Weight 3/2 modular forms.

In this section we will reduce the proof of Theorem 2 to an upper bound for the L^2 norm of a cusp form by applying a known estimate from [DS-P] for the Fourier coefficients of a cusp form of weight 3/2.

The level of q is the smallest $N \in \mathbb{Z}^+$ so that NA^{-1} is even integral, where $A = (\partial^2 q / \partial x_i \partial x_j)$. It is known that (see [Le] p. 404)

$$N \mid 2D \mid 4N^2 \,. \tag{6}$$

Let $a(n) = r(n, G) - r_q(n)$. Then by [Si] we have that

$$f(z) = \sum_{n=1}^{\infty} a(n)e(nz) \in S_{\frac{3}{2}}(\Gamma_0(N), \psi),$$

the space of cusp forms of weight $\frac{3}{2}$ for $\Gamma_0(N)$ with character ψ , a real character mod N. By the proof of Lemma 2 in [DS-P] we have that for n square-free

$$|a(n)| \ll_{\epsilon} ||f|| N^{\frac{1}{2} - \frac{1}{28} + \epsilon}$$
(7)

where

$$||f||^2 = \int_{\Gamma_0(N) \setminus H} |f|^2 y^{\frac{3}{2}} d\mu$$

with $d\mu = dx dx/y^2$. To proceed, we must estimate ||f||. Note that we are using $\Gamma_0(N)$ instead of $\Gamma_1(N)$ as is done in Lemma 2 of [DS-P]. A variant of the argument used in the next result occurred in [DFI].

Lemma 1 Let $f(z) = \sum_{n=1}^{\infty} a(n)e(nz)$ be a cusp form of weight $\frac{3}{2}$ in $S_{\frac{3}{2}}(\Gamma_0(N), \psi)$. Then

$$||f||^2 \ll \Gamma(\alpha) d(N) N^{2\alpha} \sum_{n=1}^{\infty} |a(n)|^2 n^{-\alpha},$$

where $\alpha > \frac{1}{2}$ is any number so that the series converges. Here $d(\cdot)$ is the divisor function and the implied constant is absolute.

Lemma 1 will be proven in the next section. Granting this, to finish the proof of Theorem 2 we need an upper bound for $a(n) = r(n, G) - r_q(n)$ that is uniform in D, hence in N. From (1) we have

$$|a(n)| \ll M(G)r(n,G) \tag{8}$$

since $\omega_q \leq 48$ ([Ne] p. 180). Now from (2)

$$M(G) \le h_3(D)$$

where $h_3(D)$ is the number of all classes of q with determinant D. By Minkowski reduction theory (see [Ca]) we have the crude bound

$$h_3(D) \ll \sum_{\substack{0 < a \le b \le c \\ abc \le 2D}} a^2 b \ll D^2.$$

Thus from (3) and (8)

$$|a(n)| \ll_{\epsilon} N^{\frac{1}{2}} D^{\frac{3}{2}} (nD)^{\epsilon}$$

since from (4)

$$\mathcal{N}(n;q) \ll D^{\epsilon}$$
.

Taking $\epsilon = 2 + 2\epsilon$ in the lemma gives

$$||f||^2 \ll N^4 D^3 (ND)^{4\epsilon}$$

and so Theorem 2 follows by (6) and (7).

4 A bound for the L^2 norm of a cusp form.

In this section we will prove Lemma 1 and thus complete the proof of Theorem 2. Let \mathcal{D} be the standard fundamental domain for the modular group $\Gamma = \Gamma_0(1)$

$$\mathcal{D} = \{ z = x + iy; |x| \le 1/2 \text{ and } |z| \ge 1 \}.$$

A fundamental domain for $\Gamma_0(N)$ is $\cup_{\sigma} \sigma \mathcal{D}$, where σ runs over coset representatives for $\Gamma_0(N)$ in Γ . Thus

$$||f||^2 = \sum_{\sigma} \int_{\mathcal{D}} |f(\sigma z)|^2 \operatorname{Im}(\sigma z)^{\frac{3}{2}} d\mu.$$
(9)

For each σ we have a Fourier expansion in the cusp $\sigma(\infty)$:

$$f(\sigma z) = \varepsilon_{\sigma} \sum_{n>0} a_{\sigma}(n) e\left(\frac{nz}{w_{\sigma}}\right)$$

where $|\varepsilon_{\sigma}| = 1$ and w_{σ} is the width of the cusp. Thus

$$\int_{\mathcal{D}} |f(\sigma z)|^{2} (\operatorname{Im} \sigma z)^{\frac{3}{2}} d\mu \leq \int_{\frac{\sqrt{3}}{2}}^{\infty} \int_{0}^{w_{\sigma}} |f(\sigma z)|^{2} (\operatorname{Im} \sigma z)^{\frac{3}{2}} d\mu$$
$$= w_{\sigma} \int_{\frac{\sqrt{3}}{2}}^{C_{\sigma}} \sum_{n} |a_{\sigma}(n)|^{2} e^{-4\pi n y/w_{\sigma}} y^{-1/2} dy$$
$$= 2 \int_{\frac{\sqrt{3}}{2}}^{C_{\sigma}} \int_{\frac{-w_{\sigma}}{2}}^{\frac{w_{\sigma}}{2}} |f(\sigma z)|^{2} (\operatorname{Im} \sigma z)^{\frac{3}{2}} d\mu \quad (10)$$

where $C_{\sigma} = \frac{\sqrt{3}}{2} + w_{\sigma}C$ and C > 0 is an absolute constant so that

$$\int_{t}^{\infty} e^{-y} y^{-\frac{1}{2}} dy \le 2 \int_{t}^{t+C} e^{-y} y^{-\frac{1}{2}} dy$$

for all t > 0. We may choose the coset representations

$$\sigma = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix} \text{ with } \gamma \mid N \text{ and } |\delta| \leq \frac{\gamma}{2}.$$

For such σ the cusp $\sigma(\infty) = \frac{\alpha}{\gamma}$ has width $w_{\sigma} = N/(\gamma^2, N) \le N/\gamma$, so that

$$C_{\sigma} = \frac{\sqrt{3}}{2} + w_{\sigma}C \le \frac{\sqrt{3}}{2} + \frac{NC}{\gamma}.$$

Now for z in range of (10) we have

$$\operatorname{Im}(\sigma z) = y((\gamma x + \delta)^2 + \gamma^2 y^2)^{-1} \ge 4y((N + \gamma)^2 + \gamma^2 y^2)^{-1} \gg C^{-1} N^{-2}$$

and

and

$$\left|\operatorname{Re}(\sigma z) - \frac{\alpha}{\gamma}\right| = \gamma^{-2} \left|x + \frac{\delta}{\gamma}\right| \left(\left(x + \frac{\delta}{\gamma}\right)^2 + y^2\right)^{-1} \qquad \ll \gamma^{-2} y^{-1} \le 1.$$

Thus changing variables z to $\sigma^{-1}z$ in (10) we get

$$\int_{D} |f(\sigma z)|^{2} (\operatorname{Im} \sigma z)^{\frac{3}{2}} d\mu \ll \int_{C^{-1}N^{-2}} \int_{0}^{1} |f(z)|^{2} y^{\frac{3}{2}} d\mu$$

and hence from (9)

$$||f||^2 \ll Nd(N) \int_{C^{-1}N^{-2}}^{\infty} \int_0^1 |f(z)|^2 \ y^{\frac{3}{2}} \ d\mu \tag{11}$$

since the index of $\Gamma_0(N)$ in Γ is

$$N\prod_{p|N}\left(1+\frac{1}{p}\right)\ll Nd(N)\,.$$

Putting the Fourier expansion of f(z) into (11) we get

$$\begin{split} \|f\|^2 \ll Nd(N) \sum_{n=1}^{\infty} |a(n)|^2 \int_{C^{-1}N^{-2}}^{\infty} e^{-4\pi ny} y^{\frac{1}{2}} \frac{dy}{y} \\ &= C^{-\frac{1}{2}} d(N) \sum_{n=1}^{\infty} |a(n)|^2 \int_{1}^{\infty} e^{-\frac{4\pi nt}{CN^2}} t^{\frac{1}{2}} \frac{dt}{t} \\ &= C^{-\frac{1}{2}} d(N) \sum_{n=1}^{\infty} |a(n)|^2 \int_{1}^{\infty} e^{-\frac{4\pi nt}{CN^2}} t^{\alpha} \frac{dt}{t} \end{split}$$

for any $\alpha > \frac{1}{2}$, and this is

$$\leq C^{-\frac{1}{2}} d(N) \sum_{n=1}^{\infty} |a(n)|^2 \int_0^{\infty} e^{-\frac{4\pi nt}{CN^2}} t^{\alpha} \frac{dt}{t} \\ \ll \Gamma(\alpha) N^{2\alpha} d(N) \sum_{n=1}^{\infty} |a(n)|^2 n^{-\alpha}$$

for any $\alpha > \frac{1}{2}$ so that the series converges.

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