

Notes on Fukaya category

Tianyu Yuan

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This note is a beginner version of the beginner's introduction to Fukaya category [1]. Some parts are also based on [2–4].

1 Lagrangian Floer Cohomology

In this note we assume that the background symplectic manifold (M, ω) is exact, i.e. $\omega = d\theta$. Roughly speaking, we want to study Lagrangian submanifolds in M and their intersection theory. Here we just consider exact Lagrangians L , i.e. $\theta|_L = df$ for some function $f \in C^\infty(L, \mathbb{R})$.

Example For a compact real analytic smooth manifold X , consider the cotangent bundle $M = T^*X$, then there is a canonical exact symplectic form $\omega = d\theta$, where θ is the canonical Liouville 1-form locally given by $\theta = \sum_i p_i dq_i$. Now for a submanifold $Y \subset X$, assign a defining function $m : X \rightarrow \mathbb{R}_{\geq 0}$ for $\partial Y \subset X$, then we have the standard Lagrangian:

$$L_{Y,f} = T_Y^*X + \Gamma_{df} \subset T^*X$$

where $f = \log m$ and Γ_{df} is the graph of df . Clearly $L_{Y,f}$ depends only on $f|_Y$. Since locally $\theta = \sum_i p_i dq_i$ and $p_i = \frac{\partial f}{\partial q_i}$, $\theta = dF$, where F is the pullback of f from Y to $L_{Y,f}$. Therefore $L_{Y,f}$ is exact.

1.1 Floer complex

Now we introduce the Floer complex. Roughly speaking, for each pair of Lagrangians (L_0, L_1) , we define a cochain complex $CF(L_0, L_1)$, which is freely

generated by intersection points of L_0 and L_1 over some field \mathbb{K} , together with a differential $d : CF(L_0, L_1) \rightarrow CF(L_0, L_1)[1]$.

Naturally there arise some issues: We want a well-defined Floer cohomology $HF(L_0, L_1) = \text{Ker } d / \text{Im } d$, so we need $d^2 = 0$; If L_0 and L_1 are not transversal, we have to perturb them by Hamiltonian isotopy so that the Floer complex is well defined; Also, we want to give $CF(L_0, L_1)$ a grading.

For now simply assume L_0, L_1 are compact and transversal, then define the Floer differential by counting J -holomorphic strips in M with boundary in L_0 and L_1 connecting p and q . Then the coefficient of q in dp is given by that number:

Definition 1 (*Floer differential*)

$$d(p) = \sum_{q \in \chi(L_0, L_1), \text{ind}([u])=1} (\#\mathcal{M}(p, q; [u], J))q, \quad (1)$$

More precisely, we are counting the space of maps $u : \mathbb{R} \times [0, 1] \rightarrow M$ which solves the Cauchy-Riemann equation $\bar{\partial}_J u = 0$:

$$\frac{\partial u}{\partial s} + J(u) \frac{\partial u}{\partial t} = 0, \quad (2)$$

$$\begin{cases} u(s, 0) \in L_0 \text{ and } u(s, 1) \in L_1, \\ \lim_{s \rightarrow +\infty} u(s, t) = p, \lim_{s \rightarrow -\infty} u(s, t) = q, \end{cases} \quad (3)$$

where J is any almost complex structure compatible with ω . Fix a homotopy class $[u] \in \pi_2(M, L_0 \cup L_1)$, then let $\#\hat{\mathcal{M}}(p, q; [u], J)$ be the moduli space of all solutions in $[u]$. Let $\#\mathcal{M}(p, q; [u], J)$ be $\#\hat{\mathcal{M}}(p, q; [u], J)$ quotiented by reparametrization of the domain, which in this case is the 1-dimensional translation of s .

The above boundary Cauchy-Riemann problem is linearized to be Fredholm. Denote the linearization of $\bar{\partial}_J$ at u by $D_{\bar{\partial}_J u}$, then the Fredholm index is defined by $\text{ind}([u]) = \dim \text{Ker } D_{\bar{\partial}_J u} / \text{Coker } D_{\bar{\partial}_J u}$. The dimension of moduli space $\#\hat{\mathcal{M}}(p, q; [u], J)$ is $\text{ind}([u])$ if $D_{\bar{\partial}_J u}$ is surjective everywhere on $\#\hat{\mathcal{M}}(p, q; [u], J)$. For simplicity we have to believe that for generic J this is true.

To count signed moduli spaces, orientation is necessary. Here either assume $\text{char}(\mathbb{K}) = 2$ or that L_0, L_1 are oriented and spin, then orientation is guaranteed.

1.2 Grading and Maslov index

Now it is natural to give $CF(L_0, L_1)$ gradings. To do this, try to assign to each Lagrangian L a phase function $\phi_L : L \rightarrow S^1$. If $2c_1(TM) = 0$, the bicanonical bundle $\Lambda_{\mathbb{C}}^{top} T^*M \otimes \Lambda_{\mathbb{C}}^{top} T^*M$ is trivial, then choose a nonzero section μ , define $\phi_L(p) = \frac{\mu(v_1 \wedge \dots \wedge v_n \otimes v_1 \wedge \dots \wedge v_n)}{\|\mu(v_1 \wedge \dots \wedge v_n \otimes v_1 \wedge \dots \wedge v_n)\|}$. It is easy to check the independence on basis of $T_p L$.

We still need to lift the phase function $\phi_L : L \rightarrow S^1$ to $\tilde{\phi}_L : L \rightarrow \mathbb{R}$, so require that the Maslov class $[\phi_L] \in [L, S^1]$ vanishes.

Assume $2c_1(TM) = 0$ and all Maslov classes vanish, then we can define:

Definition 2

$$\deg(p) = \tilde{\phi}_{L_0}(p) - \tilde{\phi}_{L_1}(p) + \lambda(p, L_0, L_1), \quad (4)$$

where $\lambda(p, L_0, L_1)$ accounts for the difference between $T_p L_0$ and $T_p L_1$, lifted to \mathbb{R} by a short path in $LGr(T_p M)$. Then the Maslov index is defines as:

$$\text{ind}(u) = \deg(q) - \deg(p), \quad (5)$$

where u is any J -holomorphic strip connecting p to q .

For standard Lagrangians $L_{Y,f} = T_Y^*X + \Gamma_{df} \subset T^*X$ in T^*X , grading is canonically well-defined:

Proposition 1 *The bicanonical bundle $\Lambda_{\mathbb{C}}^{top} T^*(T^*X) \otimes \Lambda_{\mathbb{C}}^{top} T^*(T^*X)$ is canonically trivial; The Maslov class $[\phi_{L_{Y,f}}] = 0$. Therefore, there is a canonical grading of $L_{Y,f}$.*

It turns out that Maslov index is the same as Fredholm index:

Proposition 2 *(A relative version of Riemann-Roch)*

$$\text{ind}(u) = \text{ind}_{\text{Fredholm}}(D_u), \quad (6)$$

1.3 Compactness

Compactness of $\mathcal{M}(p, q; [u], J)$ is given by Gromov compactness theorem. First assume a uniform energy bound. There are three possible limit behavior: strip breaking, disc bubbling and sphere bubbling. By assumption that M and L are both exact, the latter two possibilities are eliminated.

Now by some gluing statement, we can say broken strips correspond to the boundary of moduli space of strips with $ind([u]) = 2$:

$$\overline{\partial\mathcal{M}}(p, q; [u], J) = \coprod_{[u']+[u'']=[u], r \in \chi(L_0, L_1)} (\mathcal{M}(p, r; [u'], J) \times \mathcal{M}(r, q; [u''], J)), \quad (7)$$

from which it is easy to see $d^2 = 0$. Moreover, we have

Theorem 1 (Floer) *The Floer differential is well-defined: $d^2 = 0$; The Floer cohomology $HF(L_0, L_1)$ is independent of J and Hamiltonian isotopies of L_0 and L_1 .*

1.4 Transversality

In case L_0, L_1 do not intersect transversally (for example $L_0 = L_1$), try to perturb one of them by (time-dependent) Hamiltonian isotopy, i.e. choose a generic Hamiltonian $H \in C^\infty([0, 1] \times M, \mathbb{R})$ and let $CF(L_0, L_1)$ be generated by points in $L_0 \cap (\phi_H^1)^{-1}L_1$.

Equivalently, we may perturb the Cauchy-Riemann equation instead of Lagrangians:

$$\frac{\partial u}{\partial s} + J(t, u) \left(\frac{\partial u}{\partial t} - X_H(t, u) \right) = 0, \quad (8)$$

$$\left\{ \begin{array}{l} u(s, 0) \in L_0 \text{ and } u(s, 1) \in L_1, \\ \lim_{s \rightarrow +\infty} u(s, t) \text{ and } \lim_{s \rightarrow -\infty} u(s, t) \text{ are flow of } X_H \text{ from } L_0 \text{ to } L_1, \end{array} \right. \quad (9)$$

2 Fukaya category

2.1 Composition map

We start the discussion of Fukaya category by introducing higher composition maps:

Definition 3 (Higher operations)

$$\begin{aligned} \mu^k : CF(L_{k-1}, L_k) \otimes \dots \otimes CF(L_1, L_2) \otimes CF(L_0, L_1) &\rightarrow CF(L_0, L_k)[2 - k], \\ \mu^k(p_k, \dots, p_1) &= \sum_{ind([u])=2-k, q \in \chi(L_0, L_k)} (\#\mathcal{M}(p_1, \dots, p_k, q; [u], J))q, \end{aligned} \quad (10)$$

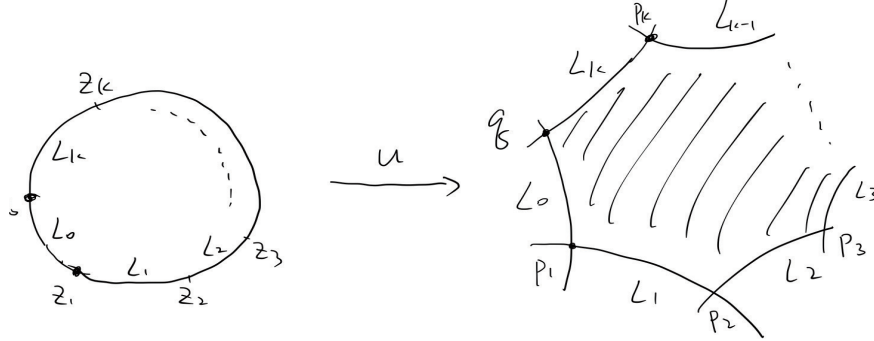


Figure 1: $\mathcal{M}(p_1, \dots, p_k, q; [u], J)$

where $\mathcal{M}(p_1, \dots, p_k, q; [u], J)$ is the moduli space of k -pointed J -holomorphic disk with boundary in L_0, \dots, L_k as shown in Figure 1, quotiented by reparametrization. The Maslov index in the sum comes from

$$\dim \mathcal{M}(p_1, \dots, p_k, q; [u], J) = k - 2 + \text{ind}([u]) = k - 2 + \text{deg}(q) - \sum_{i=1}^k \text{deg}(p_i), \quad (11)$$

Note that $\mu^1 = d$ and $(\mu^1)^2 = d^2 = 0$. This can be generalized to $k \geq 1$, which is called A_∞ -relation. Still we consider the limit curve (boundary) of every 1-dim moduli space $\mathcal{M}(p_1, \dots, p_k, q; [u], J)$. Similar as the case of $k = 1$, the limit curve cannot be disc bubbling or sphere bubbling, so the only possible behavior is nodal discs splitting the k -pointed disc. The signed sum of all these limit curves is zero, which implies

Proposition 3 (A_∞ -relations)

$$\sum_{l=1}^k \sum_{j=0}^{k-l} (-1)^* \mu^{k+1-l}(p_k, \dots, p_{j+l+1}, \mu^l(p_{j+l}, \dots, p_{j+1}), p_j, \dots, p_1) = 0, \quad (12)$$

where $*$ = $j + \text{deg}(p_1) + \dots + \text{deg}(p_j)$.

2.2 Definition of $Fuk(M, \omega)$

Definition 4 Let (M, ω) be exact with $2c_1(TM) = 0$.

$$Ob(Fuk(M, \omega)) = \{\text{exact Lagrangian } L \text{ with spin structure and graded lift}\}, \quad (13)$$

For each pair of objects L, L' which need not be transversal, choose time-dependent perturbation $H_{L,L'}$ and $J_{L,L'}$. Then for every L_0, \dots, L_k and J – holomorphic discs, choose H and L compatible with each pair's near each end. Such perturbation exists.

$$hom(L, L') = CF(L, L'; H_{L,L'}, J_{L,L'}). \quad (14)$$

Then $Fuk(M, \omega)$ is the A_∞ -category with composition maps given by Definition 3.

Usually a triangulated category is better so that we can talk about generators and mapping cones. Therefore we embed $Fuk(M, \omega)$ into a larger one $TwFuk(M, \omega)$, which contains twisted complexes of objects in $Fuk(M, \omega)$. Precisely,

Definition 5 $TwFuk(M, \omega)$

$$Ob(TwFuk(M, \omega)) = (E, \delta^E), \quad (15)$$

where $E = \bigoplus_{i=1}^N L_i[k_i]$. $\delta^E \in End^1(E)$ is lower triangular differential s.t. $\sum_{k \geq 1} \mu^k(\delta^E, \dots, \delta^E) = 0$.

A degree d morphism a in $hom(E, E')$ is just the direct sum of all $a_{ij} \in hom^{d+k'_j-k_i}(E_i, E'_j)$.

Composition maps are given by

$$\mu_{Tw}^k(a_k, \dots, a_1) = \sum_{j_0, \dots, j_k \geq 0} \mu^{k+j_0+\dots+j_k}(\delta^k, \dots, \delta^k, a_k, \dots, \delta^1, \dots, \delta^1, a_1, \delta^0, \dots, \delta^0). \quad (16)$$

Definition 6 (mapping cones) If $f \in hom^0(E, E')$ is closed, then

$$\begin{array}{ccc} (E, \delta) & \xrightarrow{f} & (E', \delta') \\ & \swarrow (1) & \downarrow \\ & & \left(E[1] \oplus E', \begin{pmatrix} \delta & 0 \\ f & \delta' \end{pmatrix} \right) \end{array} \quad (17)$$

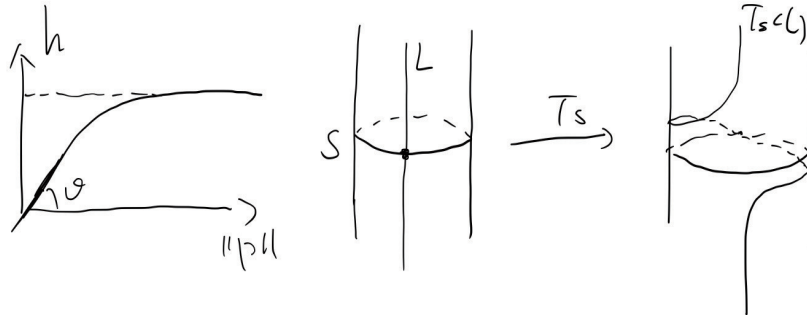


Figure 2: Dehn twist

Example (Dehn twists) Given a Lagrangian sphere S in (M, ω) , by Weinstein, a neighbourhood of S in M is just like T^*S . Now perform Dehn twists by a Hamiltonian flow $H(p, q) = h(\|p\|)$ in the complement of zero section as Figure 2.

Then $\tau_S(L)$ is the mapping cone:

Theorem 2 (Seidel)

$$\begin{array}{ccc}
 HF^*(S, L) \otimes S & \xrightarrow{ev} & L \\
 \swarrow (1) & & \downarrow \\
 & & \tau_S(L)
 \end{array} \tag{18}$$

2.3 Non-compact Lagrangians and perturbation

Here gives an example on dealing with non-compact Lagrangians by [3]. Let $U \subset X$ be open, fix a defining function m for $X \setminus U$, let $f = \log m$. Let L be standard Lagrangian given by the graph df . Fix stratification $\mathcal{S} = \{S_\alpha\}$ of X , and let $\Lambda_{\mathcal{S}} = \cup_\alpha T_{S_\alpha}^* X \subset T^*X$ be the corresponding conical Lagrangian.

If L and $\Lambda_{\mathcal{S}}$ intersect at infinity, then their intersections may be non-compact. To restrict all intersections inside a compact region, need to perturb L s.t. they are separated at infinity.

Proposition 4 (Nadler-Zaslow) *There exist $\eta > 0$ and $\delta > 0$ such that for all $\delta' \in (0, \delta]$, the normalized geodesic flow satisfies*

$$\gamma_{\delta'}(\bar{L}_{m \leq \eta}) \cap \bar{\Lambda}_{\mathcal{S}} = \emptyset, \quad (19)$$

where $\gamma_{\delta'}$ is the flow of Hamiltonian $H(x, \xi) = |\xi|$ on $T^*X \setminus X$.

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