

# Geodesic Normal Coordinates on a Surface.

To help us in understanding the properties of geodesics on surface, we shall be making use of a special kind of coordinate system. To define this, we start with a surface  $S$ , a point  $p_0$  in  $S$  (and hence in  $\mathbb{R}^3$ ) and a parameterization  $S(u, v)$  of  $S$  near  $p_0$  with  $S(0, 0) = p_0$  and  $S_u|_{(0,0)}$  and  $S_v|_{(0,0)}$  an orthonormal basis for the tangent space of  $S$  at  $p_0$  (or, equivalently,  $E|_{(0,0)} = 1$ ,  $G|_{(0,0)} = 1$  and  $F|_{(0,0)} = 0$ ).

We can do this without loss of generality, since we can always arrange these conditions, starting with an arbitrary parameterization, by making a translation and linear transformation of the parameters. [Alternatively, one can think of "graphing the surface over its tangent plane" as we discussed earlier, that is writing it as  $S(u, v) = (u, v, f(u, v))$  after a translation and rotation). With  $S$  as described, we define a new parameterization  $\hat{S}$  by  $\hat{S}(a, b) = \gamma_{a,b}(1)$  where  $\gamma_{a,b}$  = the geodesic in  $S$  with  $\gamma_{a,b}(0) = p_0$  and  $\gamma'_{a,b}(0) = aS_u + bS_v$ . The basic Existence and Uniqueness Theorem from differential equations applied to the equations of geodesics gives that  $\hat{S}$  is defined on some neighborhood  $W$  of  $(0, 0)$  in  $\mathbb{R}^2$  and is  $C^\infty$  on  $W$ . We can also show that



$\hat{S}$  is nonsingular at  $(0,0)$  so that  $\hat{S}$  is a local parameterization of  $S$  near  $p_0$ . For this, we need to compute  $\hat{S}_a$  and  $\hat{S}_b$ .

Now note that  $\gamma_{a,b}(t) = \gamma_{a,b}(\alpha t)$  for each  $\alpha$  (so small that the maps are defined); this comes from Uniqueness together with the homogeneity in  $t$  of the geodesic equations:  $\gamma(t)$  is a geodesic if and only if  $\gamma(\alpha t)$  is a geodesic, ( $\alpha \neq 0$ ).

$$\begin{aligned} \text{So } \frac{d}{dt} \hat{S}(ta, tb) \Big|_{t=0} &\stackrel{\text{def.}}{=} \frac{d}{dt} \gamma_{ta, tb}(t) \Big|_{t=0} \\ &= \frac{d}{dt} \gamma_{a,b}(t) \Big|_{t=0} = aS_u + bS_v. \end{aligned}$$

In particular  $\hat{S}_a = S_u$  and  $\hat{S}_b = S_v$

since  $\hat{S}_a =$  (by definition)  $\frac{d}{dt} \hat{S}_{ta} \Big|_{t=0}$  and

similarly for  $\hat{S}_b$ . Thus  $\hat{S}_a \Big|_{(0,0)} \times \hat{S}_b \Big|_{(0,0)} =$

$$S_u \Big|_{(0,0)} \times S_v \Big|_{(0,0)} \neq 0.$$

The parameterization of  $S$  we have just defined is called a geodesic normal coordinate system at  $p_0 \in S$ . It is in fact unique up to linear orthogonal changes of  $(a, b)$ .



What is particularly interesting is the associated idea of geodesic polar coordinates obtained from polar coordinates on the  $(a, b)$  plane.

We take the  $a$ -axis as  $\theta=0$  and  $r = \sqrt{a^2+b^2}$  as usual. So we get geodesic polar coordinates (by definition) on  $S$  by sending  $(r, \theta) \rightarrow \hat{S}(r \cos \theta, r \sin \theta) \stackrel{\text{def}}{=} G(r, \theta)$

These are coordinates on  $S$  (near  $p_0$ ) in the same sense that polar coordinates are coordinates on  $\mathbb{R}^2$ : i.e.,  $(r, \theta) \rightarrow \hat{S}(r \cos \theta, r \sin \theta)$  is  $C^\infty$  but it is singular at  $r=0$  and many-to-one as far as  $\theta$  is concerned.

We are going to want to study  $\langle G_r, G_r \rangle$ ,  $\langle G_r, G_\theta \rangle$  and  $\langle G_\theta, G_\theta \rangle$

( $G$  for geodesic!) Note first that by definition a curve  $r \rightarrow G(r, \theta_0)$ ,  $\theta_0$  fixed,  $r > 0$ , is a geodesic. (Look at the definition of  $\hat{S}$ ). In fact it is a geodesic with arclength parameter. So  $\langle G_r, G_r \rangle \equiv 1$  ( $r > 0$ ).

Our next main goal is to show that  $\langle G_r, G_\theta \rangle = 0$ . But before we do that, let us look at some examples.



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(1)  $\mathbb{R}^2 \subset \mathbb{R}^3$   $S(u, v) = (u, v, 0), p_0 = (0, 0, 0)$   
 Geodesic polar coordinates are just polar coordinates on  $\mathbb{R}^2$ , that is  
 $G(r, \theta) = (r \cos \theta, r \sin \theta, 0)$

(2)  $S^2$  (unit 2-sphere =  $\{(x, y, z) : x^2 + y^2 + z^2 = 1\}$ )  
 in  $\mathbb{R}^3$ ,  $p_0 = (0, 0, 1)$   
 $S(u, v) = (u, v, \sqrt{1 - u^2 - v^2})$ .

$G(r, \theta) =$  endpoint of geodesic starting at  $(0, 0, 1)$ , arc length parameter, traversed up to distance  $r$ . Since geodesics are great circles

$$G(r, \theta) = (\sin r \cos \theta, \sin r \sin \theta, \cos r)$$

$$G_r = (\cos r \cos \theta, \cos r \sin \theta, -\sin r)$$

and

$$\begin{aligned} \langle G_r, G_r \rangle &= \cos^2 r (\cos^2 \theta + \sin^2 \theta) + (-\sin r)^2 = \\ &= \cos^2 r + \sin^2 r = 1 \quad \text{as} \end{aligned}$$

expected.

$$G_\theta = (-\sin r \sin \theta, \sin r \cos \theta, 0)$$

and

$$\langle G_r, G_\theta \rangle = 0 \quad \text{as we shall prove in general.}$$

Finally

$$\begin{aligned} \langle G_\theta, G_\theta \rangle &= (-\sin r \sin \theta)^2 + (\sin r \cos \theta)^2 \\ &= \sin^2 r. \end{aligned}$$

Think about how this gives the expected answer for the  $r =$  positive constant,  $\theta$  goes from 0 to  $2\pi$  (a curve which is a circle, radius  $|\sin r|$ )



We now turn to our goal:  
Gauss's Lemma:  $\langle G_r, G_\theta \rangle \equiv 0$ .

Proof: We compute (for  $r > 0$ )

$$\frac{d}{dr} \langle G_r, G_\theta \rangle = \langle G_{rr}, G_\theta \rangle + \langle G_r, G_{r\theta} \rangle$$

$$= \langle T(G_{rr}), G_\theta \rangle + \langle G_r, G_{r\theta} \rangle$$

0 "since  $r$ -curve is geodesic!

$$= 0 + \frac{1}{2} \frac{\partial}{\partial \theta} \langle G_r, G_r \rangle = 0$$

since  $\langle G_r, G_r \rangle \equiv 1$ . So  $\langle G_r, G_\theta \rangle|_{\theta \text{ fixed}, r > 0}$

is constant. But  $\lim_{r \rightarrow 0^+} \langle G_r, G_\theta \rangle|_{\theta \text{ fixed}} = 0$

(since the  $\Delta$  change in  $S$  associated to change  $\Delta \theta$  goes to 0 as  $r$  goes to 0: write everything in  $(a, b)$  terms to see this). So  $\langle G_r, G_\theta \rangle \equiv 0$  ( $\theta$  fixed) any  $r$ , as required.  $\square$

This Lemma actually implies that geodesics are locally minimal connections.

Here is the idea: A curve  $\sigma(t) = (r(t), \theta(t))$  has  $\sigma'(t) = \left(\frac{dr}{dt}, \frac{d\theta}{dt}\right)$  in  $r, \theta$  terms, that is its  $\mathbb{R}^3$  tangent vector is

$$\frac{dr}{dt} G_r + \frac{d\theta}{dt} G_\theta$$



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Using what we know ( $\langle G_r, G_r \rangle = 1$ ,  $\langle G_r, G_\theta \rangle = 0$ ) we get

$$\|\sigma'(t)\|^2 = \left(\frac{dr}{dt}\right)^2 + \langle G_\theta, G_\theta \rangle \left(\frac{d\theta}{dt}\right)^2$$

Now suppose we want to go from  $p_0$  ( $r=0$ ) to some  $q = (r_0, \theta_0)$ .

Then

$$\begin{aligned} \text{length}(\sigma) &= \int \sqrt{\left(\frac{dr}{dt}\right)^2 + \langle G_\theta, G_\theta \rangle \left(\frac{d\theta}{dt}\right)^2} \\ &\geq \int \left|\frac{dr}{dt}\right| \geq \text{change in } r = r_0. \end{aligned}$$

So  $\text{length}(\sigma) \geq r_0$ . On the other hand, the geodesic  $t \rightarrow (t, \theta_0)$  in  $(r, \theta)$  coordinates has length  $r_0$ . So no other curve <sup>(with same endpoints)</sup> is shorter. Also if  $\text{length}(\sigma) = r_0$ , then (since  $\langle G_\theta, G_\theta \rangle > 0$ ) it must be that  $\frac{d\theta}{dt} = 0$  and

since  $\int \left|\frac{dr}{dt}\right| = \text{change in } r$ , it must also be that  $r(t)$  is monotone.

So  $\sigma$  is the geodesic indicated with some (monotone) parameterization.

This of course all only works if  $q$  is "close" to  $p_0$ , namely  $q \in \hat{S}$  (some disc around  $(0,0)$  in  $\mathbb{R}^2$  on which  $\hat{S}$  is 1-1, onto its image nonsingularly).